

Basic parameters

Expansion rate H_0 Mass density Ω_m (relative to critical $\rho_c = \frac{3H_0^2}{8\pi G}$)including Ω_b baryon density Ω_{CDM} non-baryonic CDMCosmological constant Ω_Λ

Power spectrum

 σ_8 normalization of amplitude n tilt

Fundamental Observations

Distance Scale

Traditional Distance scale ladder

- Solar System
- Trigonometric parallax of nearby stars
- Main sequence fitting to star clusters
- Bright star standard candles (e.g. Cepheids)
- Secondary distance indicators
e.g., Type Ia SN, Tully-Fisher

Absolute methods (independent of the ladder above)

- e.g.
- water masers
 - gravitational lens time delay
 - SZ effect

Fundamental Observations

Age Scale

- Oldest stars $t_{GC} = 13.3 \pm 0.1$ Gyr (recent)
consistently in the range
- Globular clusters $t_{GC} = 12 - 15$ Gyr
calibration of stellar evolution depends on distance scale

- White dwarf

maximum ages along cooling function

measure drop off in luminosity function
(no WDs fainter than $M \approx 16$)

Account for preceding stellar evolution

$$t_{WD} = 12.5^{+1.4}_{-3.5} \text{ Gyr}$$

- Radio isotope chronometers eg. Th/Eu

$$t_{Th} = 12.8 \pm 3 \text{ Gyr} \quad \text{oldest star}$$

measure abundances of r-process radioactive elements
in ancient, metal poor stars relative to stable elements

- Interstellar dust grains

measure abundances of radio isotopes in meteoritic grains
account for evolution of parent star, evolution preceding it,
and age of solar system (each 4 or 5 Gyr):

$$t_{\text{grain}} \approx 13.7 \pm 1.3 \text{ Gyr}$$

Fundamental Observations

Mass Density

There are many ways to estimate the gravitational mass density of the universe - the over-all $\Omega_m = \frac{\rho_m}{\rho_{crit}}$ ($\rho_{crit} = \frac{3H_0^2}{8\pi G}$)

To give just a few examples:

- Cluster M/L

measure mass of cluster (e.g., virial $M \approx \frac{1.5 R_e \sigma^2}{G}$)
measure luminosity, combine to get M/L (assume to be global)
measure average luminosity density of the universe j
by integrating luminosity function of galaxies
combine to get density: $\rho_m = \frac{M}{L} j$

- weak lensing

measure lensing shear over large scales
 $\Omega_m \approx 0.2$ by this method so far

- peculiar velocity field

masses like clusters distort the Hubble flow

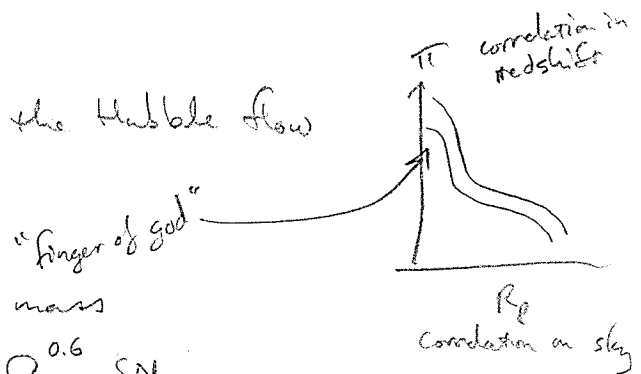
amount of distortion depends on mass

$$\frac{\delta V}{V} \approx \frac{d \ln H}{d \ln r} \frac{S_p}{f} \approx -\frac{1}{3} \frac{\Omega_m^{0.6}}{b} \frac{SN}{N}$$

b is bias factor relating mass ρ and galaxy number density N

$\Omega_m \approx 0.25$ by this method

- LOTS of others



Fundamental Observations

cosmological constant

Need Λ for many reasons, most especially

- Age problem

The universe is too young for its contents (t_{GC})
without Λ for observed (H_0, Ω_m)

- Cosmic acceleration

$q_0 < 0$ as seen in the Hubble diagram
for Type Ia SN

- Geometry as seen in the CMB

The location of the first peak of the
acoustic power spectrum tells us the geometry is flat.

$$\Omega_{tot} = 1 \quad \text{with} \quad \Omega_k \approx 0$$

so $\Omega_m + \Omega_k = 1$ but independently we measure

$\Omega_m < 1$. Need Λ to make up the deficit.

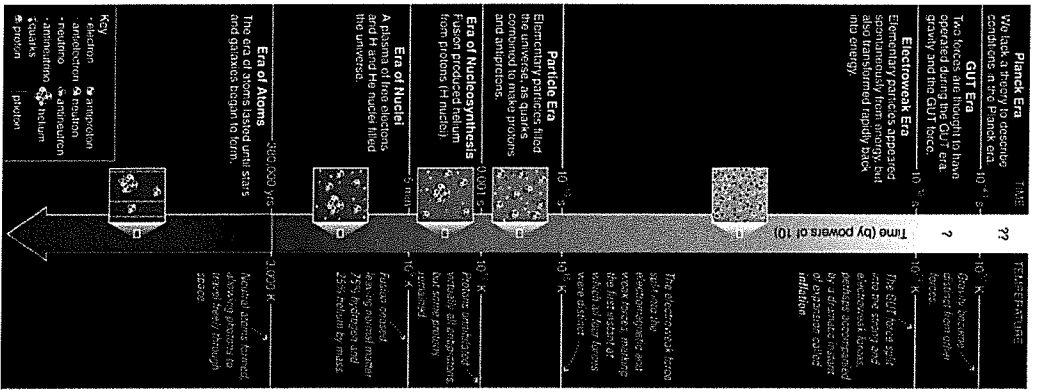


Figure 17.4 This timeline summarizes conditions and transitions that marked the early eras of the universe.

time Temp

Cosmic Timeline

Time	Event
$t \sim 10^{-43}$ s	Planck scale (<i>speculative</i>)
$t \sim 10^{-38}$ s	GUT scale (<i>speculative</i>)
$t \sim 10^{-35}$ s	Inflation (<i>speculative</i>)
$t \sim 10^{-12}$ s	Standard Model forces emerge
$t \sim 10^{-8}$ s	WIMPs decouple (<i>speculative</i>)
$t \sim 10^{-5}$ s	quarks condense into baryons (<i>baryogenesis</i>)
$t \sim 10^{-4}$ s	proton-antiproton annihilation ends
$t \sim 1$ s	neutrinos decouple
$t \sim 4$ s	electron-positron annihilation ends
$t \sim 10^2$ s	Big Bang Nucleosynthesis
$t \sim 10^5$ yr	Matter-radiation equality
$t \sim 4 \times 10^5$ yr	Atoms form, CMB emerges
$t \sim 5 \times 10^6$ yr	Gas temperature decouples from radiation
$t \sim 10^7$ yr	Dark Ages
$t \sim 5 \times 10^8$ yr	Cosmic dawn (first stars)
$t \sim 10^9$ yr	Galaxies form
$t \sim 4 \times 10^9$ yr	Peak star formation
$t \sim 9 \times 10^9$ yr	Sun forms
$t \sim 13 \times 10^9$ yr	Life on earth

Early U radiation dominated

$$t \lesssim 10^5 \text{ yr}$$

$$a \sim t^{1/2}$$

$$T \sim a^{-1}$$

$$Tt^2 \sim \text{constant}$$

Early Universe

Aside from the hypothetical period of exponential expansion during Inflation, the early universe is radiation dominated with

$$a(t) \sim t^{1/2} \quad T \sim a^{-1} \sim (1+z) \quad \epsilon_r = \alpha T^4$$

$\alpha = \text{radiation constant}$

Can substitute T for a in Friedmann equation

to find that $T^2 t = \text{constant}$

where the constant depends only on physical constants.

Particle - antiparticle pairs exist in equilibrium with the radiation field until they "freeze out" at a time when the temperature T_2 drops below the energy required for spontaneous pair production

electron - positron freeze-out occurs shortly after the neutrino background decouples from the radiation background, so the $e^- - e^+$ entropy contributes to T_2 but not T_ν .

Consequently, these relativistic backgrounds differ in temperature by

$$T_2 = \left(\frac{4}{11}\right)^{1/3} T_\nu$$

where the $\frac{4}{11}$ comes from the difference in degrees of freedom for bosons (photons) and fermions (neutrinos)

$$I_{\pm} = \int_0^{\infty} \frac{x^3 dx}{e^x \pm 1} \quad \begin{array}{l} + \text{ bosons} \\ - \text{ fermions} \end{array} \rightarrow \begin{array}{l} \epsilon_r = \alpha T_2^4 \\ \epsilon_r = \frac{7}{8} \alpha T_\nu^4 \end{array}$$

Cosmic Microwave Background (CMB)

A fundamental tenet of the Hot Big Bang cosmology is that the universe should be full of a relic radiation field. This was discovered in the '60s and is now known as the CMB.

It has a perfect thermal spectrum with $T_{\text{pho}} = 2.75 \text{ K}$.

The surface of last scattering at $z = 1090$ marks the transition from an opaque plasma to a transparent atomic gas after recombination forms free electrons and protons into hydrogen.

The CMB temperature is very uniform on the sky - it is very nearly the same in every direction we look: the early universe obeys the cosmological principle, being both homogeneous and isotropic when $t_2 = 3.8 \times 10^5 \text{ yr}$.

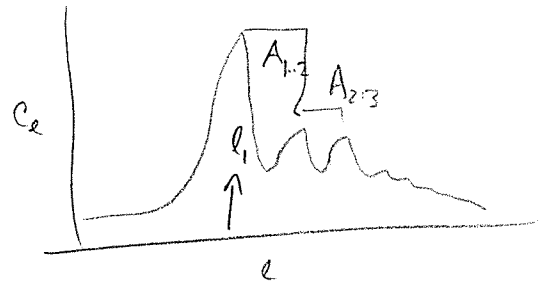
There is a dipole moment on the sky due to our motion wrt the CMB of $\sim 600 \text{ km s}^{-1}$ (discovered in '70s)

There are fluctuations on smaller scales at the level $\frac{\Delta T}{T} \approx 10^{-5}$ (discovered in '92)

These fluctuations represent early seeds for the formation of large scale structure, which were predicted to have a larger amplitude ($\frac{\Delta T}{T} \sim 10^{-2}$). That this is not the case is one line of evidence for non-baryonic cold dark matter, which is needed to grow structure in the allotted time.

(The other primary line of evidence for CDM is $\Omega_m > \Omega_b$.)

The pattern of fluctuations in the CMB, quantified as the acoustic power spectrum, is a sensitive probe of cosmic parameters



- The location of the first peak, l_1 , is primarily sensitive to the geometry of the universe. Its observed location indicates that the universe is flat $|\Omega_k| < 0.005$.
- The amplitudes of the peaks and their ratios $A_{1,2}, A_{2,3}, \dots$ are sensitive to the matter content of the universe. Alternate compression and rarefaction peaks are a competition between baryonic drag and the net driving term provided by non-baryonic dark matter.

Large Scale Structure

LSS starts from the small fluctuations seen in the CMB: $\delta = \frac{1}{3} \frac{\Delta}{T}$.
 From this initial condition, structure grows as $\delta \sim a(t)$.
 The dark matter forms structure first. The baryons subsequently fall into the holes formed by CDM after they fully decouple from the photons (not till $z \sim 200$).