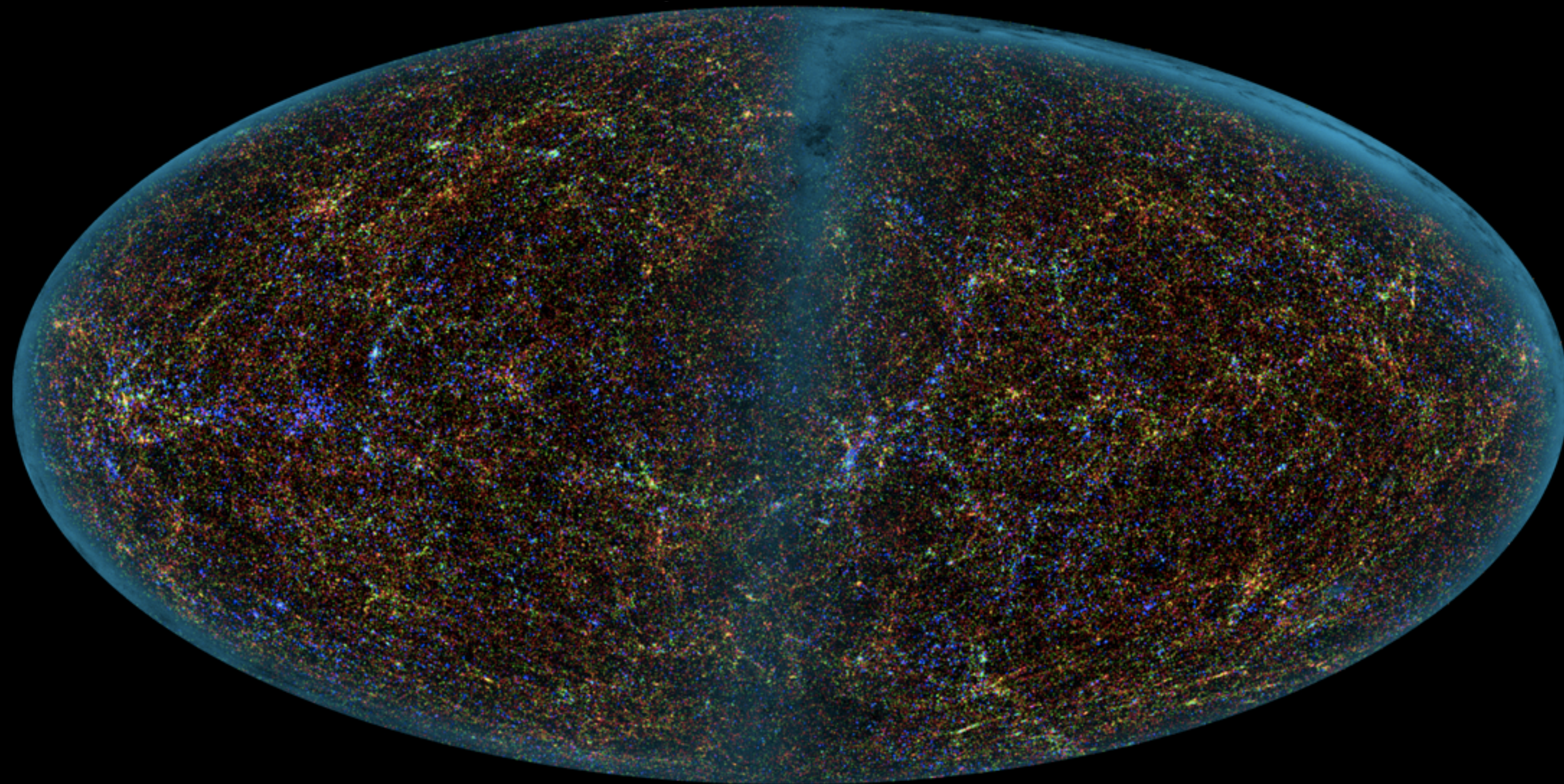


Cosmology

and Large Scale Structure



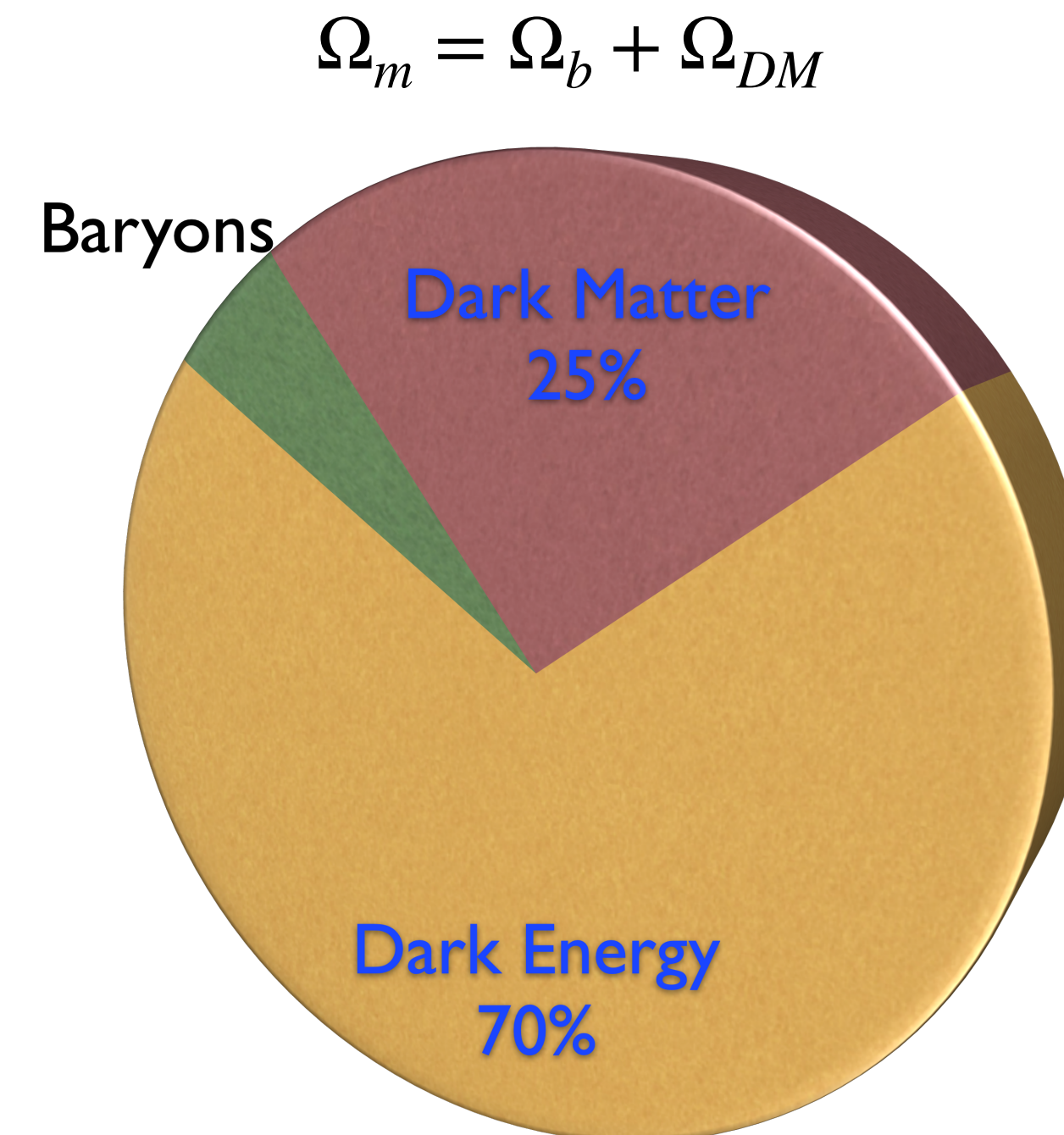
Today
Large Scale Structure
The Power Spectrum
Measurements of Ω_m

Empirical Pillars of the Hot Big Bang

1. Hubble Expansion
2. Big Bang Nucleosynthesis Ω_b
3. Cosmic Microwave Background

Auxiliary Hypotheses

- Dark matter Ω_{DM}
- Dark Energy Ω_Λ



Current mass-energy content of the universe

“Vanilla LCDM”

| | | | |
|---------------------------------------|------------------------------|--------------------|---|
| mass density | Ω_{m_0} | 0.30 | give or take a bit |
| normal matter | Ω_b | 0.05 | baryons - from BBN |
| mass that is <i>not</i> normal matter | Ω_{CDM} | 0.25 | cold dark matter |
| cosmic background radiation | Ω_r | 5×10^{-5} | photons |
| neutrinos | $0.001 < \Omega_\nu < 0.002$ | | for 3 neutrino flavors with $0.06 < \sum_{i=1}^3 m_{\nu_i} < 0.12 \text{ eV}$ upper limit from cosmic structure formation lower limit from neutrino oscillations |
| dark energy | Ω_Λ | 0.70 | energy density of vacuum |

$$\Omega_x = \frac{\rho_x}{\rho_{crit}}$$

$$\rho_{crit} = \frac{3H_0^2}{8\pi G}$$

e.g. $\Omega_\nu = \frac{\sum m_\nu}{93 \text{ eV}}$

since $n_\nu = \frac{9}{11} n_\gamma$

Measurements of the gravitating mass density

- Cluster M/L
 - measure M/L of a cluster, combine with measured luminosity density of universe.
- Weak lensing
 - measure shear over large scales
- Peculiar Velocity Field
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Measurements of the gravitating mass density

- Cluster M/L

- measure M/L of a cluster, combine with measured luminosity density of universe.
- j from integrating the luminosity function of galaxies:

$$\rho_m = \left(\frac{M}{L} \right)^j$$

- Also, cluster baryon fractions:

$$f_b = \frac{M_b}{M_{tot}} \quad \longrightarrow \quad \Omega_m = \frac{\Omega_b}{f_b}$$

- both assume clusters are representative of the whole.

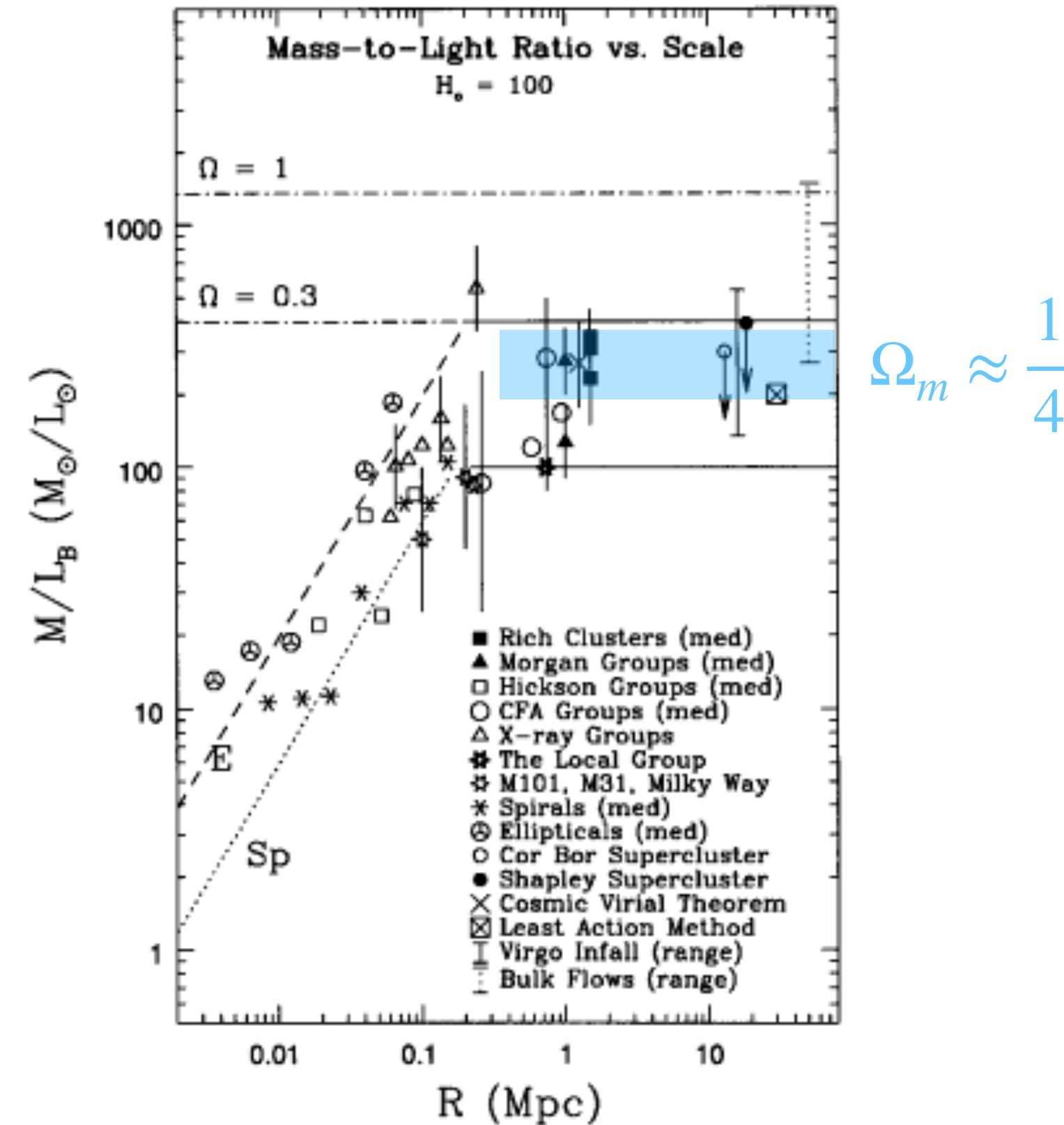
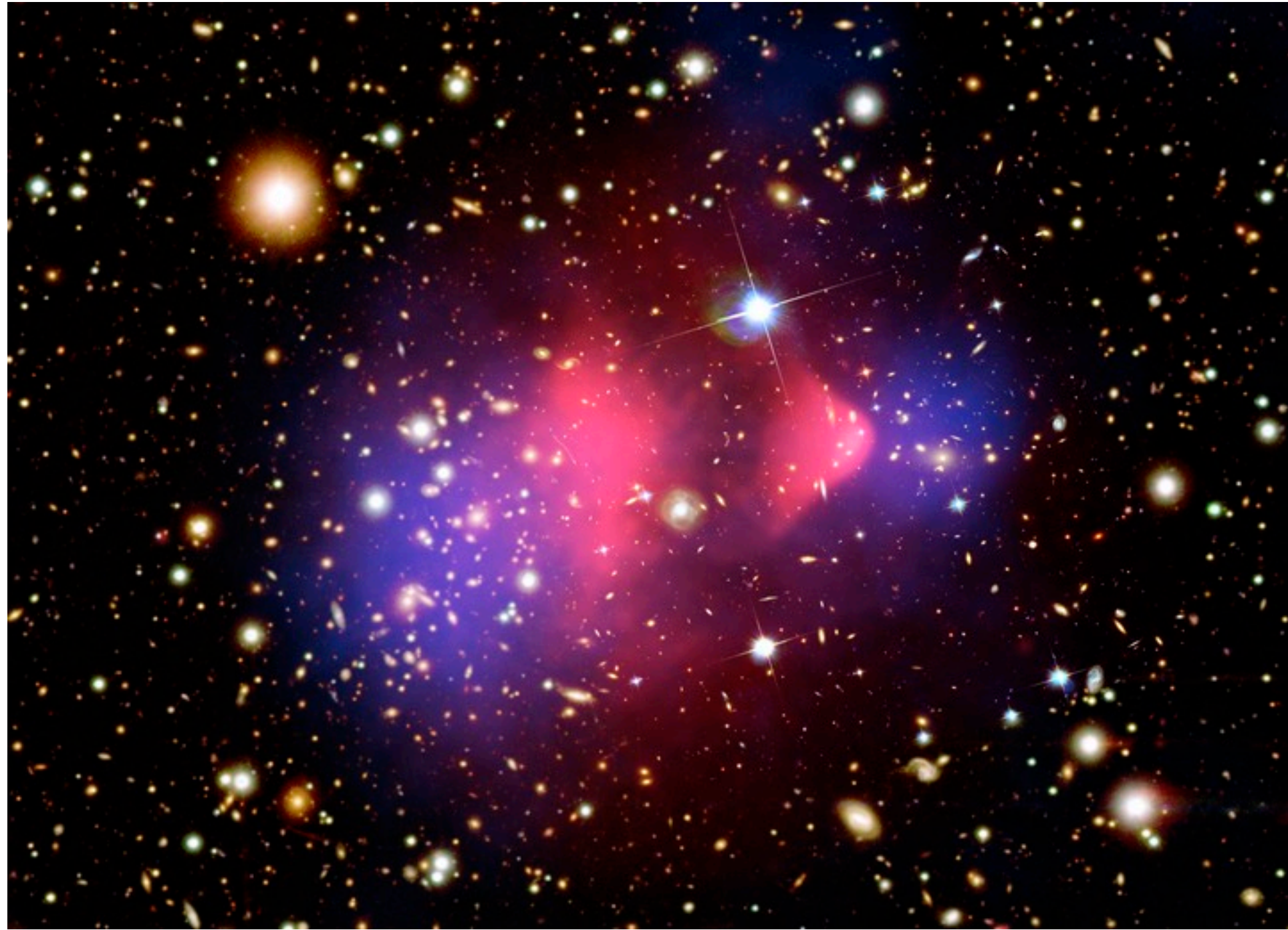


FIG. 2.—Composite mass-to-light ratio of different systems—galaxies, groups, clusters, and superclusters—as a function of scale. The best-fit $M/L_B \propto R$ lines for spirals and ellipticals (from Fig. 1) are shown. We present median values at different scales for the large samples of galaxies, groups and clusters, as well as specific values for some individual galaxies, X-ray groups, and superclusters. Typical 1σ uncertainties and 1σ scatter around median values are shown. Also presented, for comparison, are the M/L_B (or equivalently Ω) determinations from the cosmic virial theorem, the least action method, and the *range* of various reported results from the Virgocentric infall and large-scale bulk flows (assuming mass traces light). The M/L_B expected for $\Omega = 1$ and $\Omega = 0.3$ are indicated.

Bullet cluster

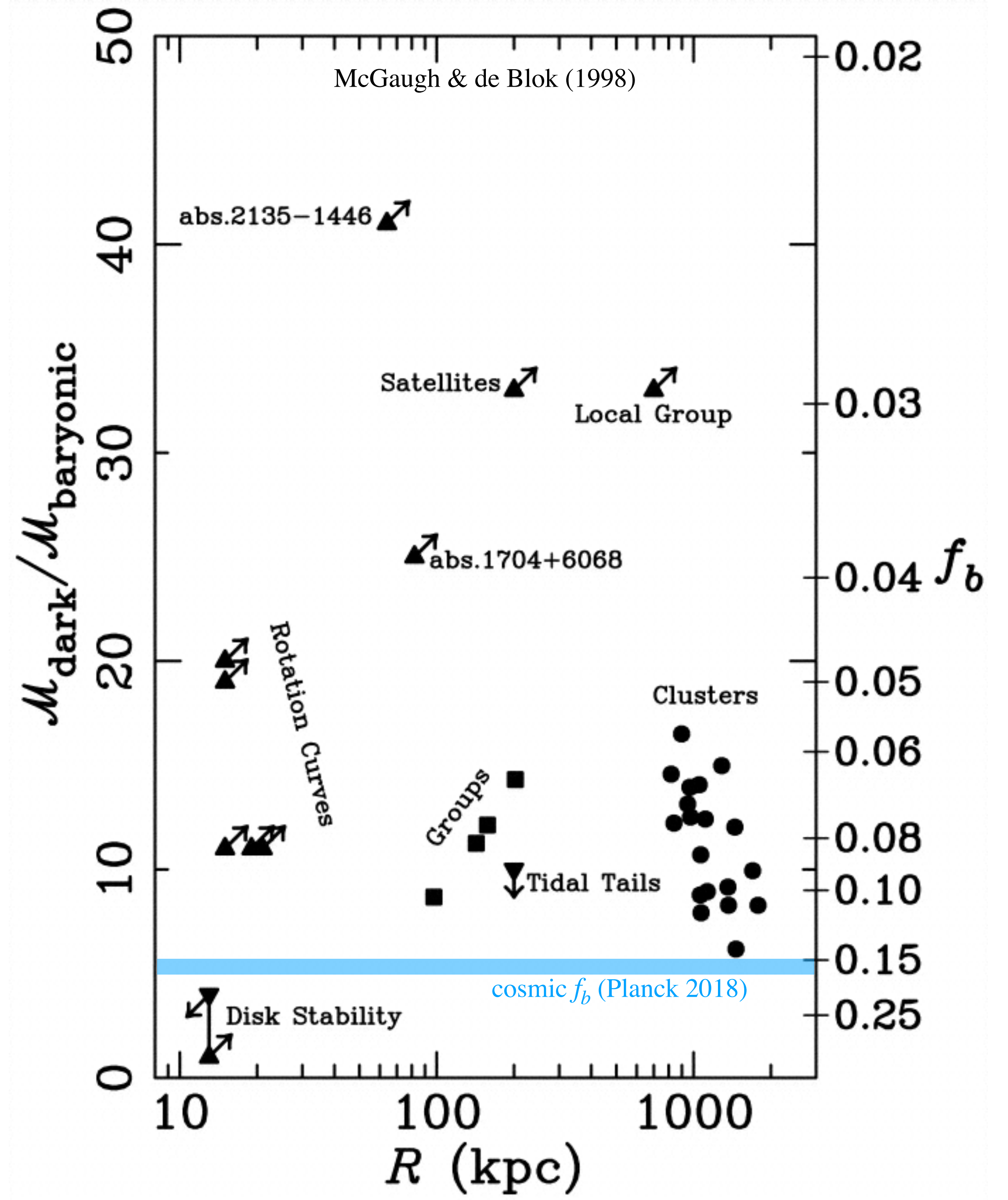


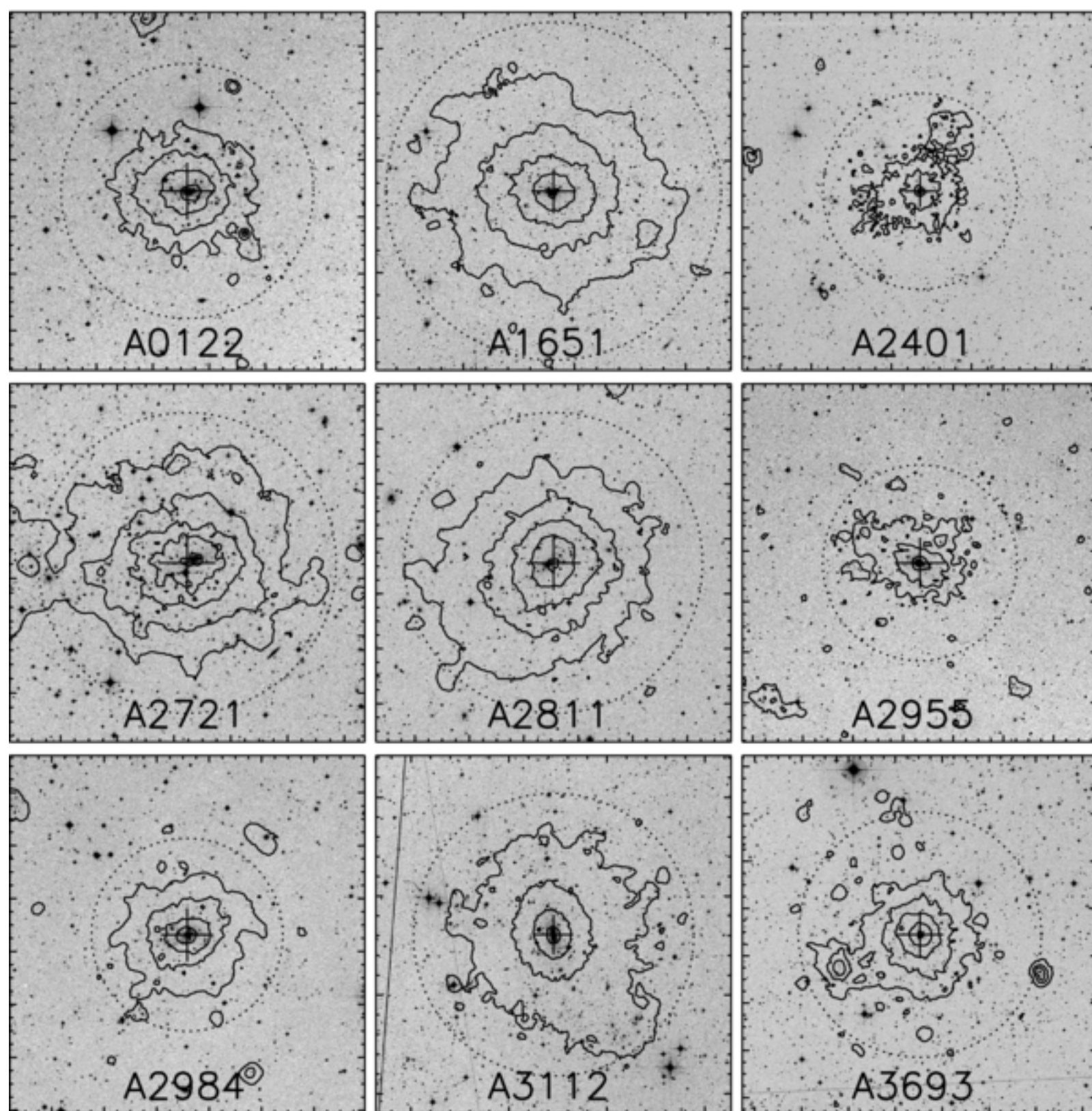
– cluster baryon fractions

$$f_b = \frac{M_b}{M_{tot}} \quad \circ \longrightarrow \quad \Omega_m = \frac{\Omega_b}{f_b}$$

Measure cluster baryonic mass M_b from luminosity of X-ray gas (pink) plus stars in galaxies (yellow)

Measure cluster dynamical mass M_{tot} from X-ray temperature (pink) or weak lensing (blue) or velocity dispersion





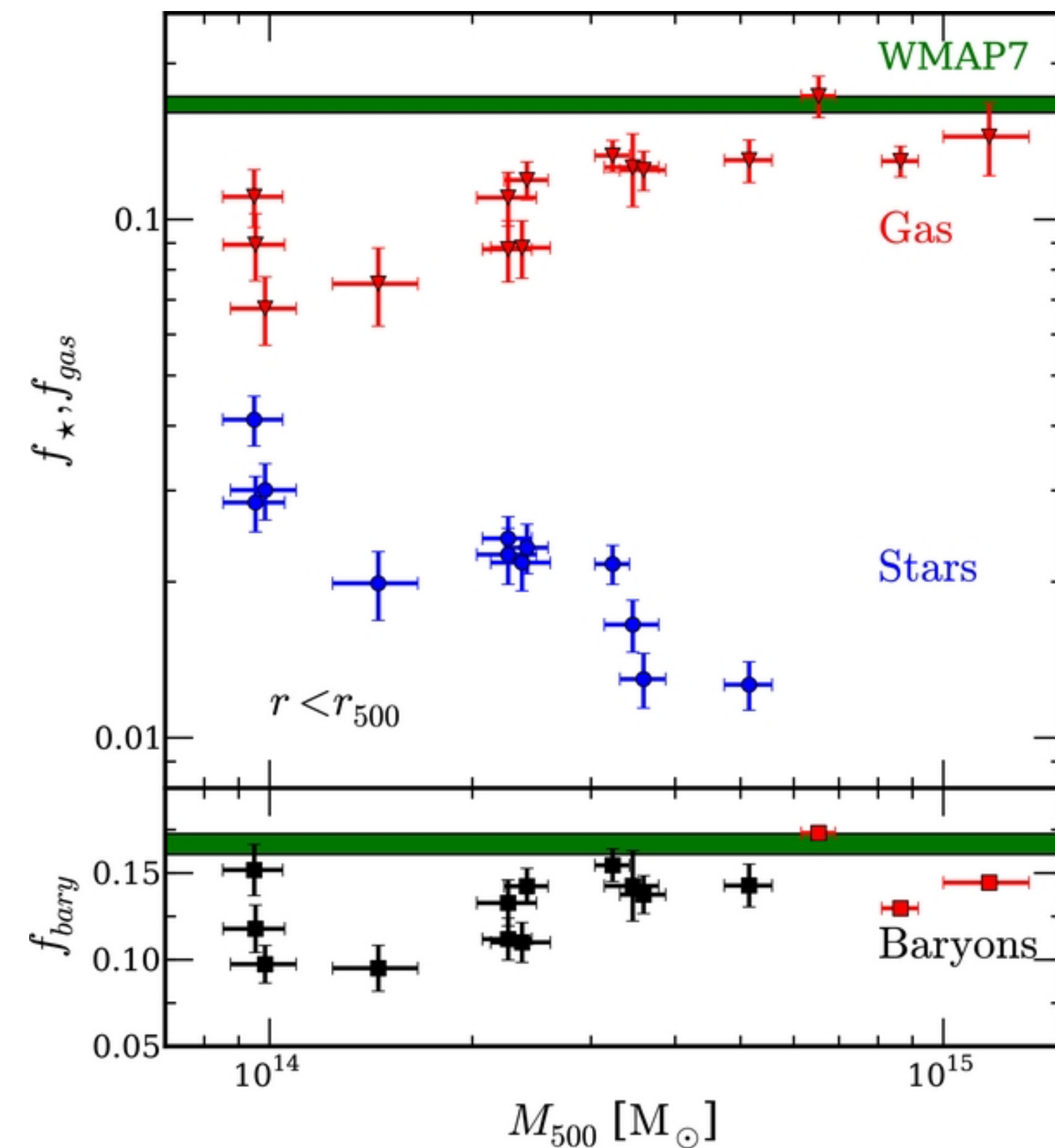
– cluster baryon fractions

$$f_b = \frac{M_b}{M_{tot}} \quad \circ \longrightarrow \quad \Omega_m = \frac{\Omega_b}{f_b}$$

Measure cluster baryonic mass M_b from luminosity of X-ray gas (contours)
plus stars in galaxies (black)

Measure cluster dynamical mass M_{tot} from X-ray temperature (contours)
or weak lensing
or velocity dispersion

Gonzalez et al. (2013)



Most of the baryonic mass in rich clusters is in the hot, X-ray emitting gas of the ICM (intracluster medium). Only the most massive clusters approach the cosmic fraction found in fits to the acoustic power spectrum of the CMB. Lower mass clusters suffer a missing baryon problem.

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Measurements of the gravitating mass density

- Weak lensing
 - measure shear over large scales

Dark Energy Survey
arxiv:2002.11124

$$\Omega_m \approx 0.18 \pm 0.04$$

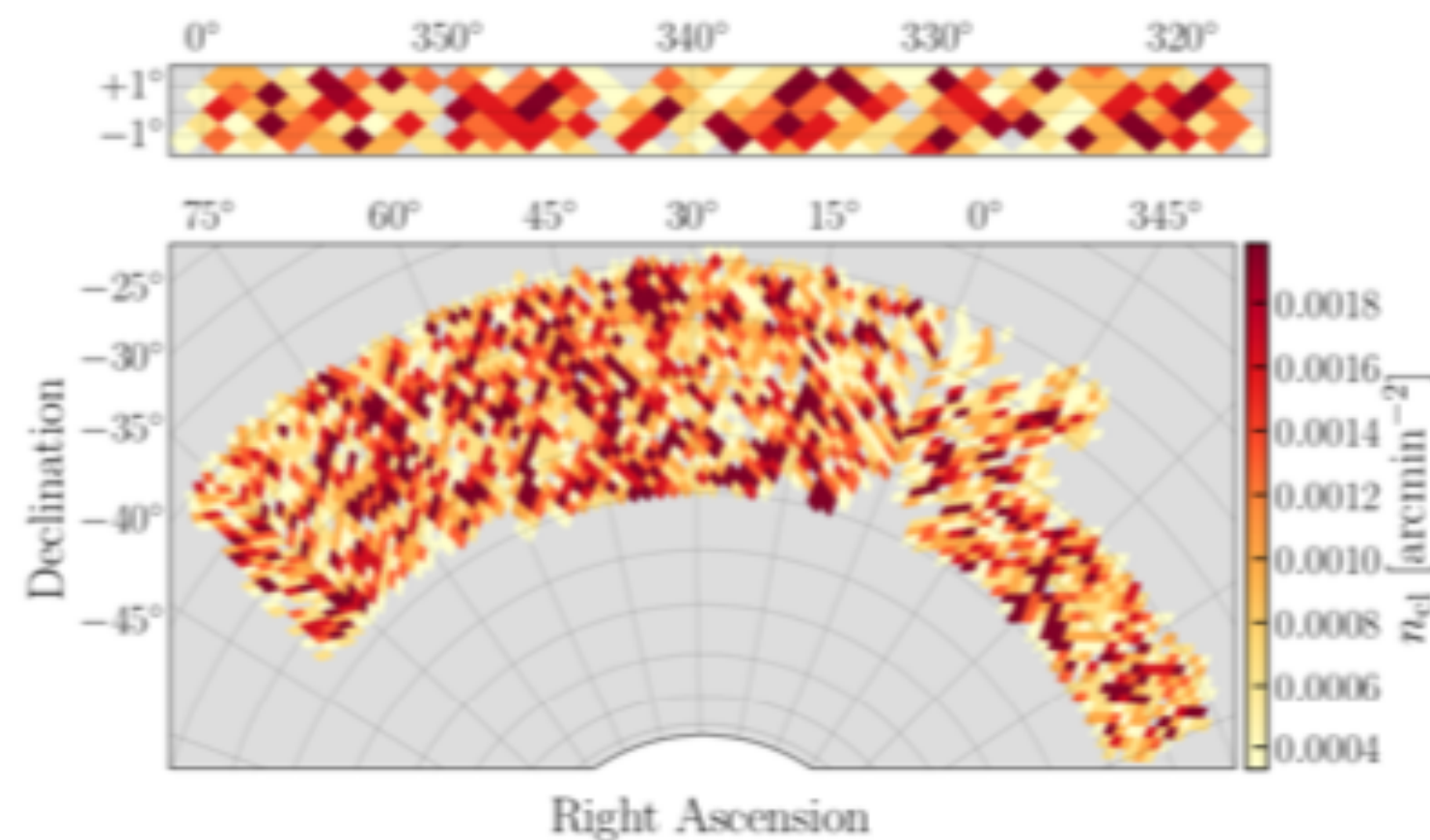
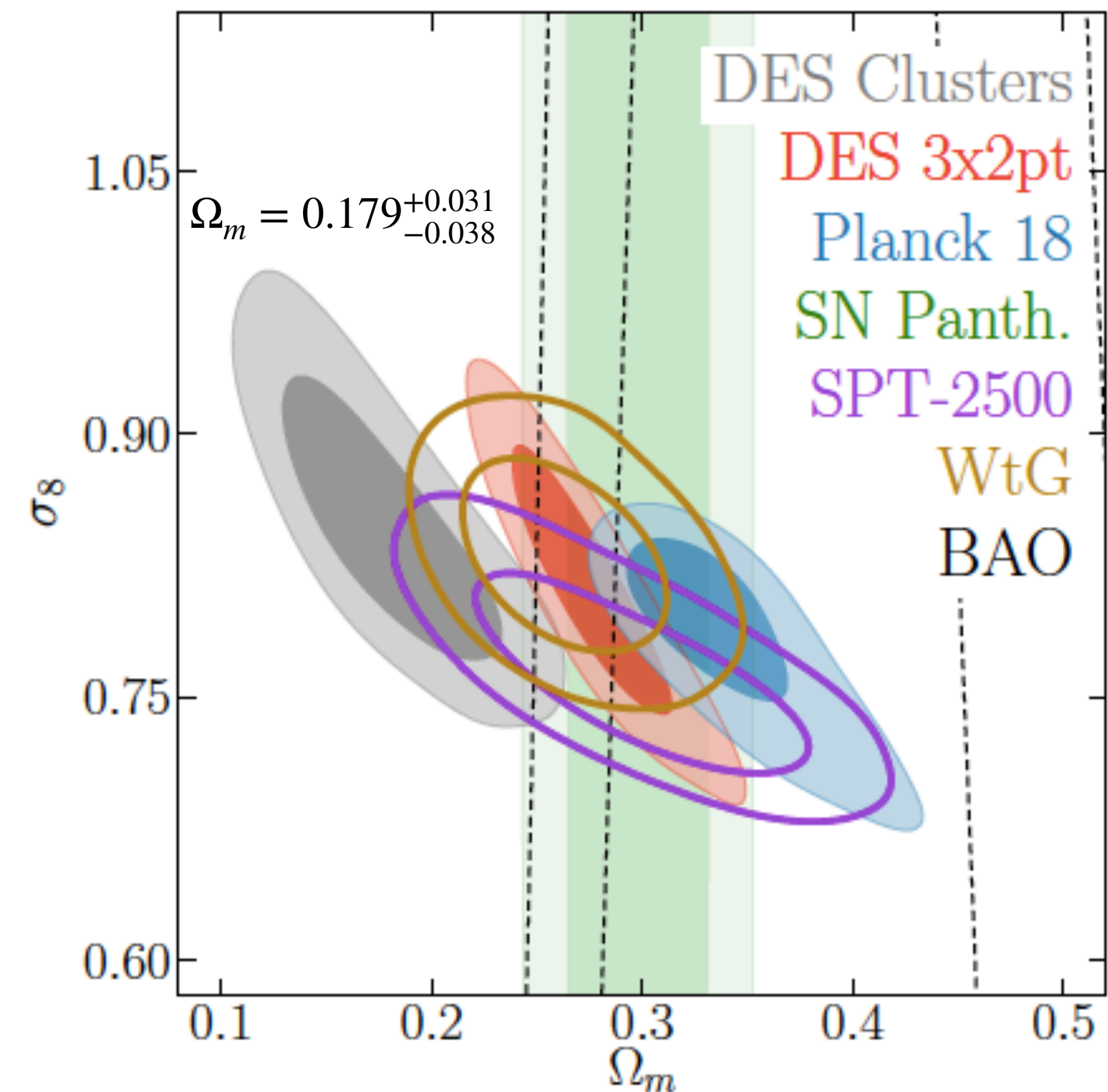


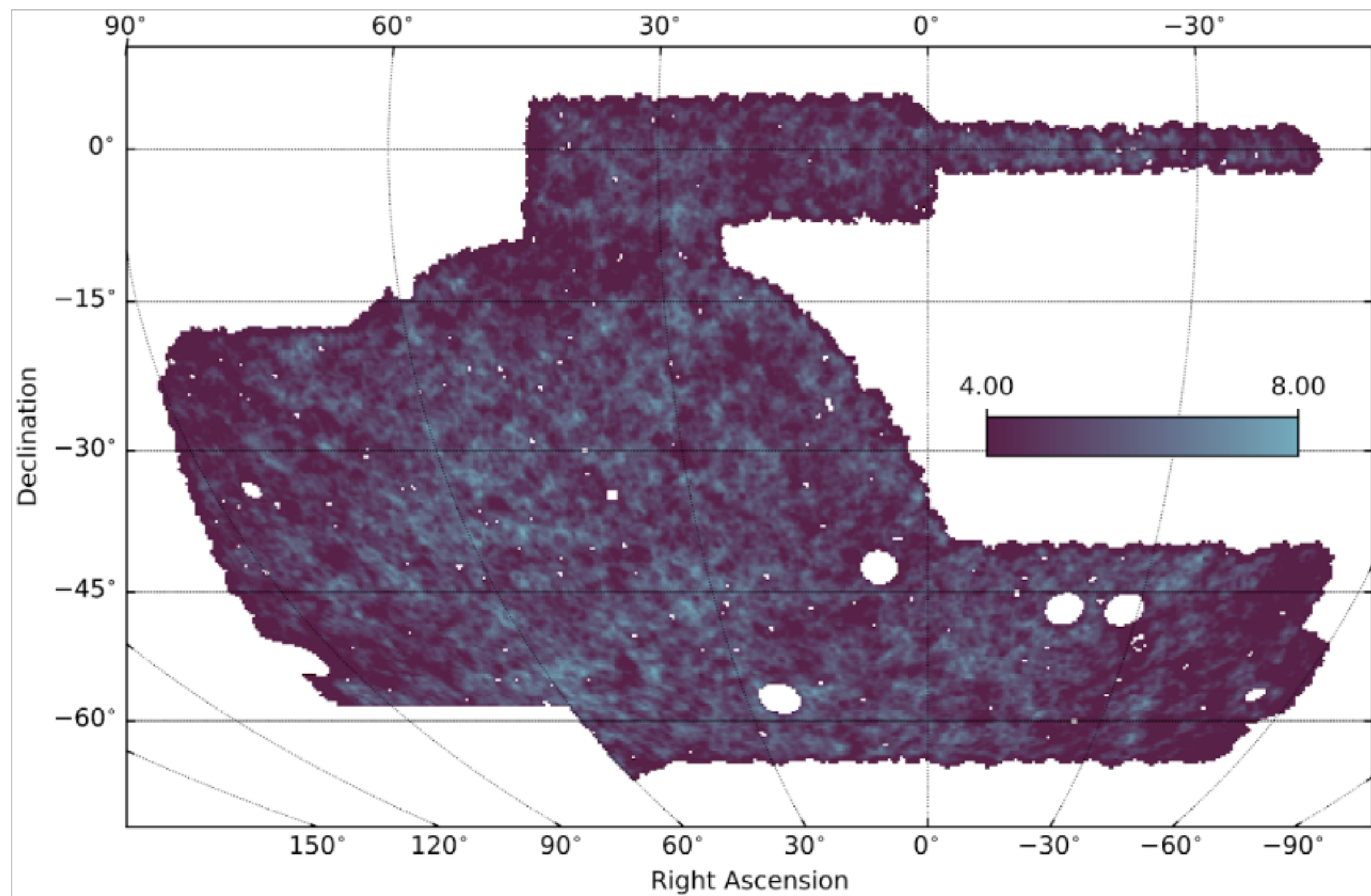
FIG. 1. The DES Y1 redMaPPer cluster density over the two non-contiguous regions of the Y1 footprint: the Stripe 82 region (116 deg²; *upper* panel) and the SPT region (1321 deg²; *lower* panel).



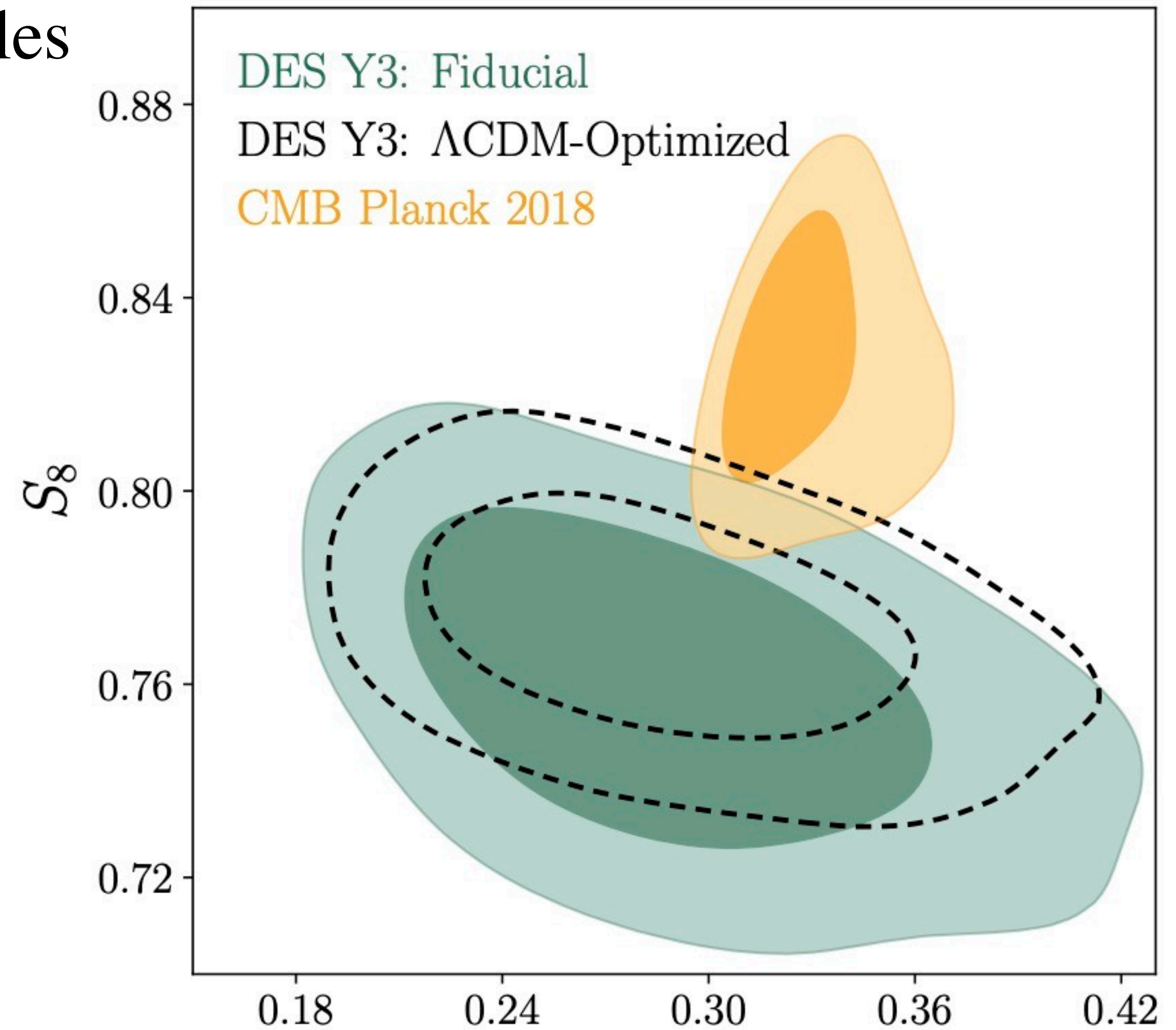
Measurements of the gravitating mass density

- Weak lensing
 - measure shear over large scales

Dark Energy Survey
Amon et al. (2022, PRD, 105, 023514)



$$S_8 = \sigma_8 \left(\frac{\Omega_m}{0.3} \right)^{1/2} = 0.759^{+0.025}_{-0.023}$$



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Virgo-centric infall

The Virgo cluster is the largest nearby over-density.
Its gravity distorts the Hubble flow.
We fall towards it so it appears to recede less than
it should by an amount that depends on its mass

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TONRY AND DAVIS (1981)

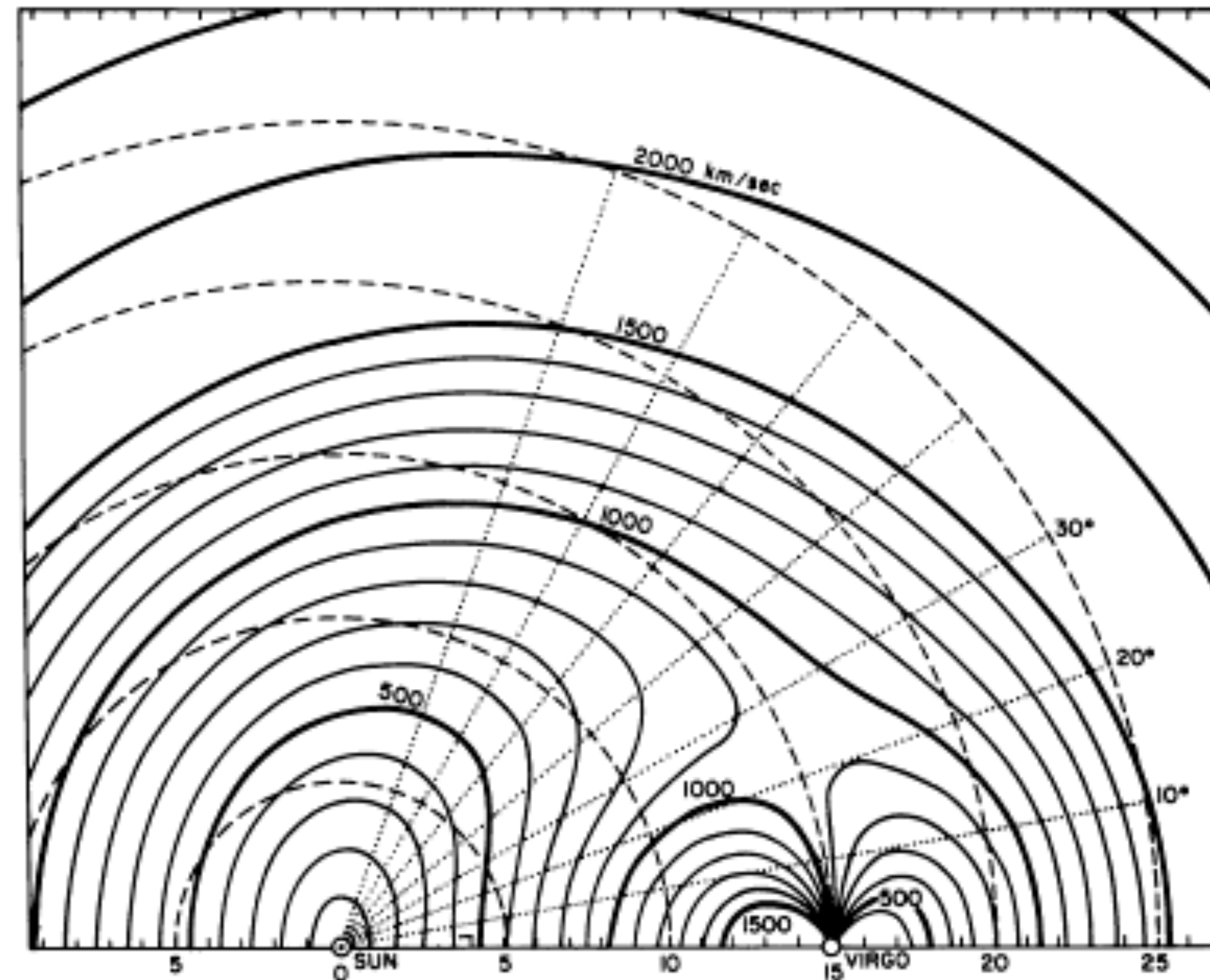


FIG. 1.—On a two-dimensional grid with the Earth and the Virgo cluster on the x axis, redshift contours are plotted for a Hubble flow perturbed by a Virgo-centric flow. An infall velocity of 400 km s^{-1} at our position is assumed. A pure Hubble flow would be concentric circles.

Measurements of the gravitating mass density

- Peculiar Velocity Field
 - measure deviations from Hubble flow

in linear regime $\frac{\delta\rho}{\rho} \ll 1$

$$\frac{\delta v}{v} \approx \frac{d \ln H}{d \ln \rho} \frac{\delta \rho}{\rho} \approx - \frac{1}{3} \frac{\Omega_m^{0.6}}{b} \frac{\delta \rho_g}{\rho_g}$$

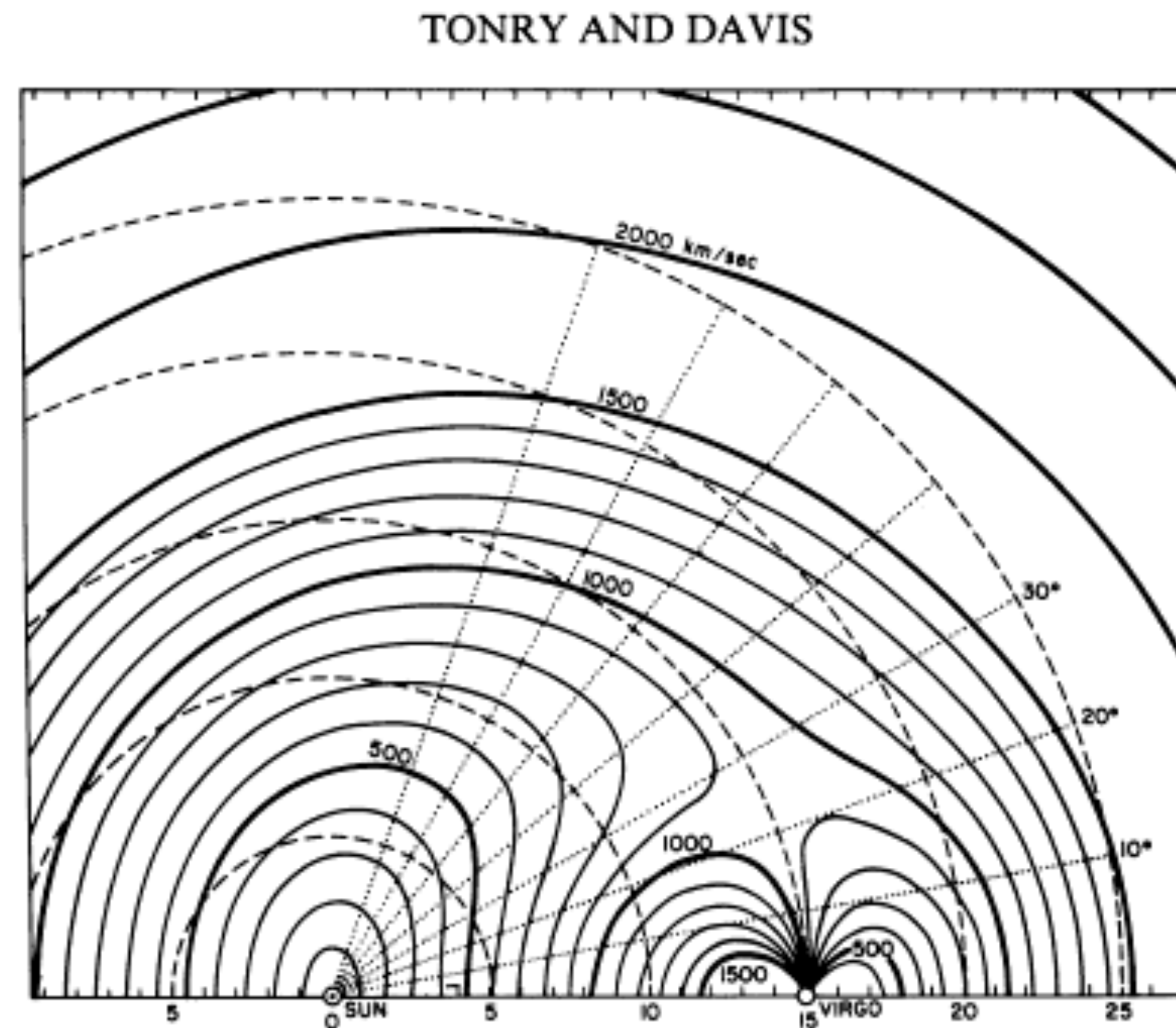
peculiar velocity

distortion in Hubble flow induced by

mass over-density

BIAS b relates galaxy over-densities to mass over-densities

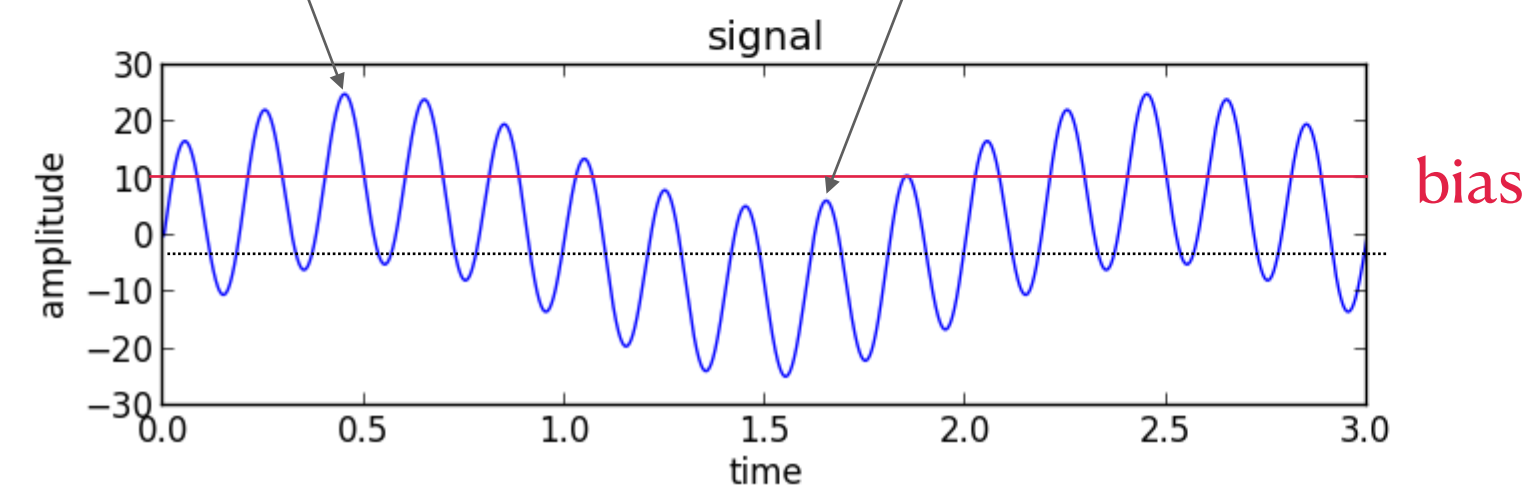
$$\Omega_m = 0.25 \pm 0.05$$



BIAS b
relates
galaxy over-densities
to mass over-densities

Peaks above line make a galaxy

Peaks below line do not



$$\frac{\delta\rho_g}{\rho_g} = b \frac{\delta\rho_m}{\rho_m}$$

$$\sigma_8 = \frac{1}{b} \quad \text{in a sphere of radius 8 Mpc}$$

Davis et al. (1980) found

$$\Omega_m = 0.4 \pm 0.1$$

with a modern distance scale this becomes

$$\Omega_m = 0.25 \pm 0.05$$

basically unchanged for over 40 years

Lines are lines of constant Ω_m

Ω_m

$\frac{\delta\rho}{\rho}$

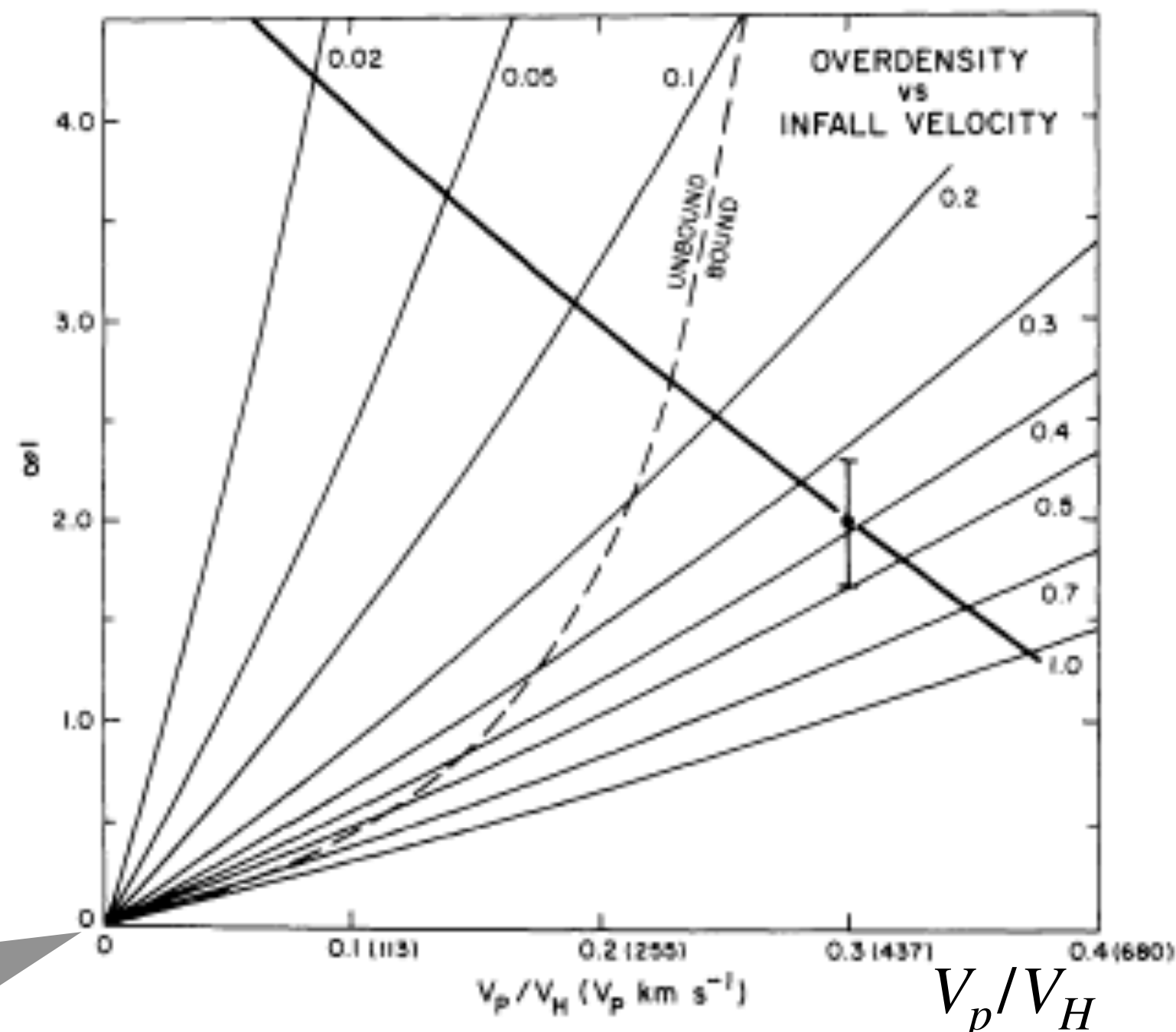


FIG. 1.—The mean overdensity of Virgo vs. v_p/v_H for various values of Ω . The x -axis is also labeled with v_p , using a recessional velocity to Virgo of 1020 km s^{-1} . The measured overdensity is prescribed by the heavy line, and is marked at the favored position as given by the anisotropy of the Hubble flow and microwave background radiation. The error bar is an estimate of the 90% confidence limit of our determination of $\bar{\delta}$. Models to the right of the dotted line are bound to Virgo.

ESTIMATES OF v_p

| Velocity | Source |
|----------------|-----------------------------------|
| 380 ± 75 | Smoot and Lubin 1979 |
| 480 ± 75 | Aaronson <i>et al.</i> 1980 |
| 350 ± 50 | de Vaucouleurs and Bollinger 1979 |
| $290 \pm 30^*$ | Yahil 1980 |
| 190 ± 130 | Schechter 1968 |

* Calculated with respect to the centroid at the local group as defined by Yahil *et al.* 1977.

mated, roughly by density of galaxies as r^{-2} ; if the mass reduced peculiar velocity this model will apply which have not yet Virgo core and with km s^{-1} are assumed scale peculiar motion effect on the comp $100 \text{ km s}^{-1} \text{ Mpc}^{-1}$, 300 km s^{-1} galactic Aaronson, and Hu sional velocity of Table 2 lists the co dom sample when b assumed here v_p/v_H The mean overd not at all sensitive

VIRGO

Distance^a (Mpc)

| | |
|-------|-------|
| 0-5 | |
| 5-10 | |
| 10-15 | |
| 15-20 | |
| 20-25 | |
| 25-30 | |
| 30-35 | |

$\bar{\delta}$

* Assumes distance and H

^b 600 obje

^c Within l

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$z = 0$

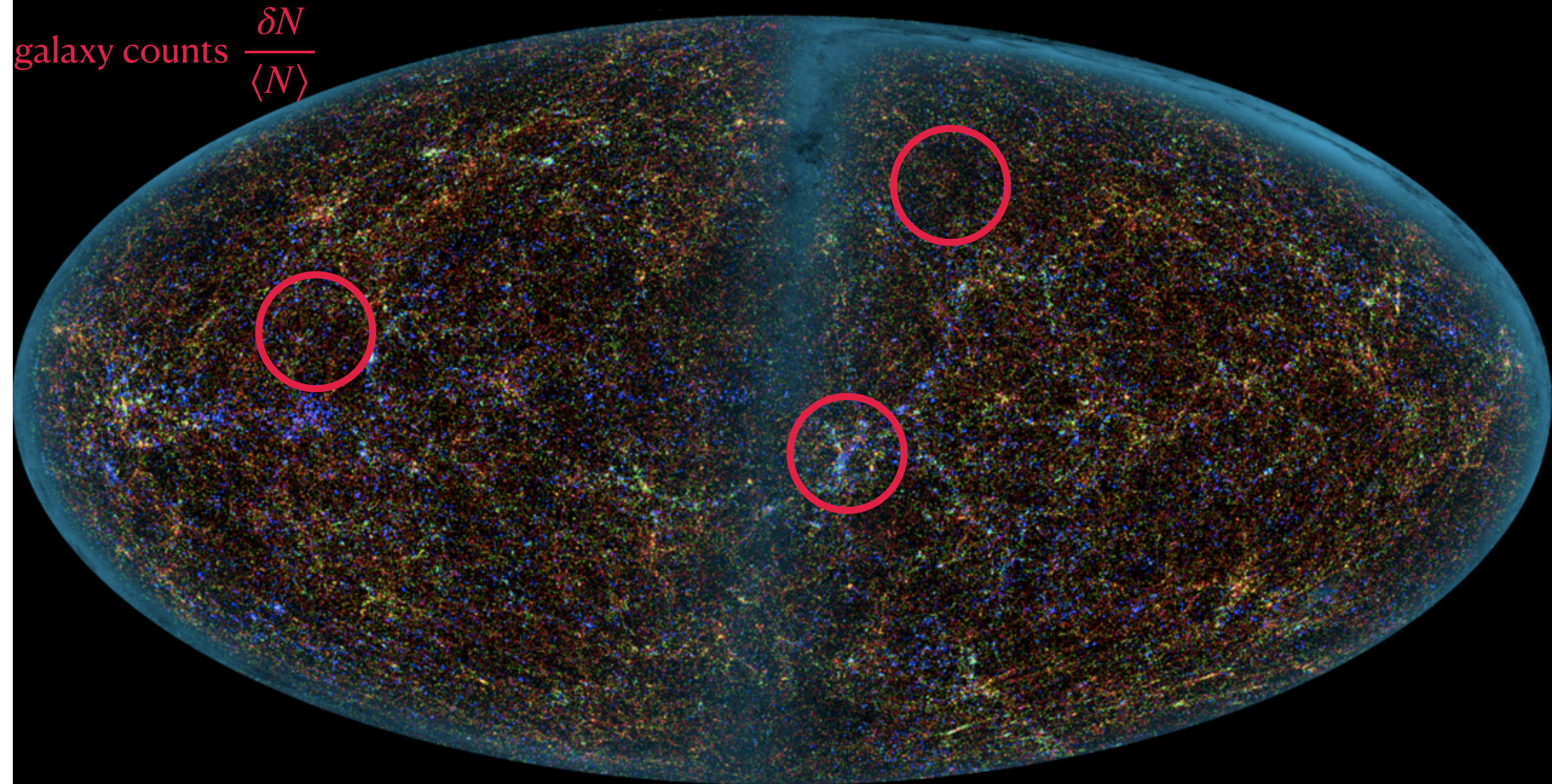
$z = 1090$

Structure formation basics:

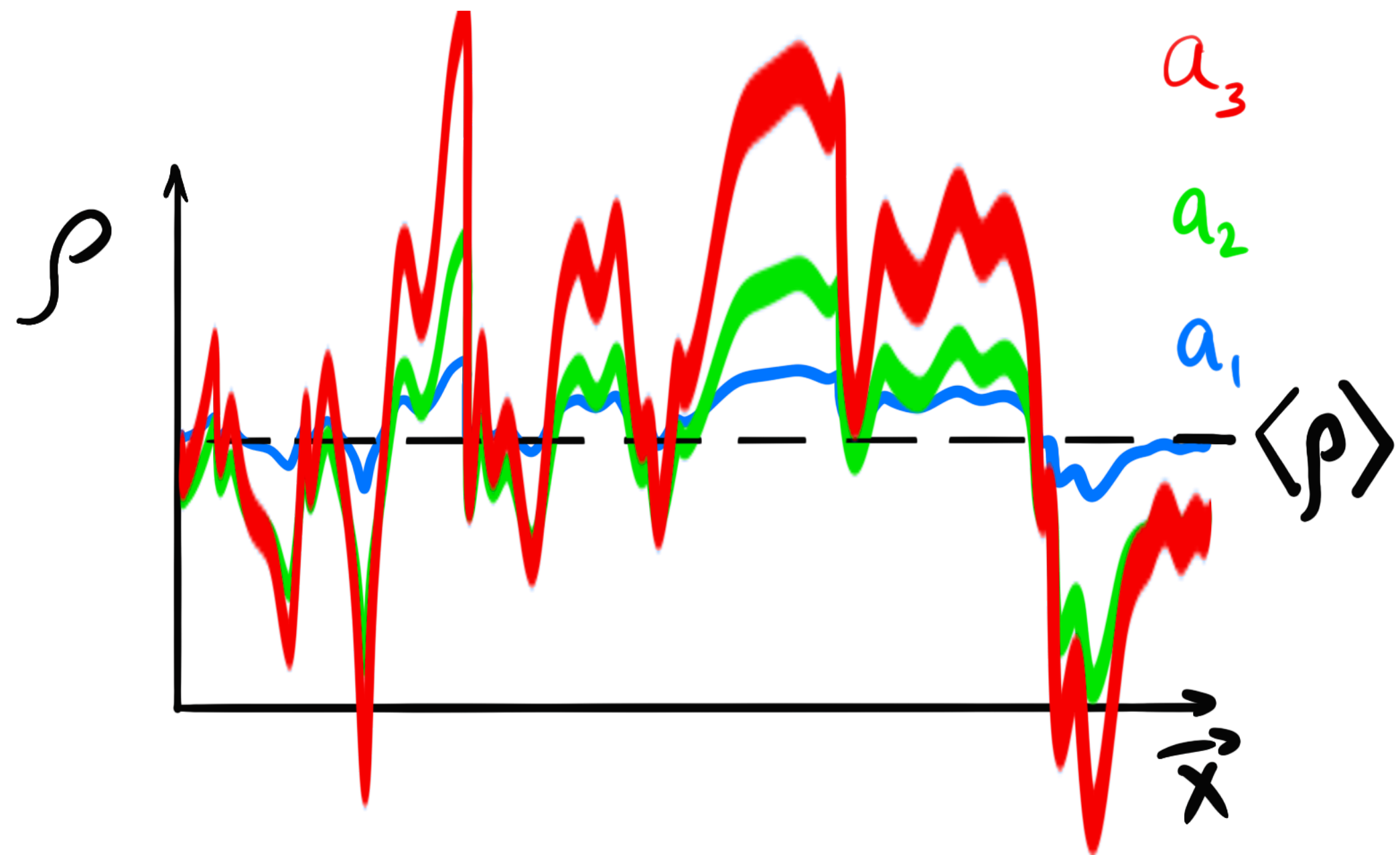
Density perturbations $\delta = \frac{\rho - \langle \rho \rangle}{\langle \rho \rangle}$

grow as $\delta(t) \sim a(t)$.

In the early universe, $\langle \rho \rangle = \rho_{\text{crit}}$.



over-density δ
grows with time



You can't get here from there

The factor of 100 offset in density and temperature fluctuations is a prime motivation for non-baryonic **cold dark matter** — a substance for which perturbations δ can grow sufficiently large while not leaving an imprint of corresponding magnitude on the CMB.

Radiation and baryon plasma tightly coupled at recombination, so a fluctuation in density is reflected by one in temperature: $\frac{\delta \rho}{\rho} \propto \frac{\Delta T}{T}$.

Large Scale Structure

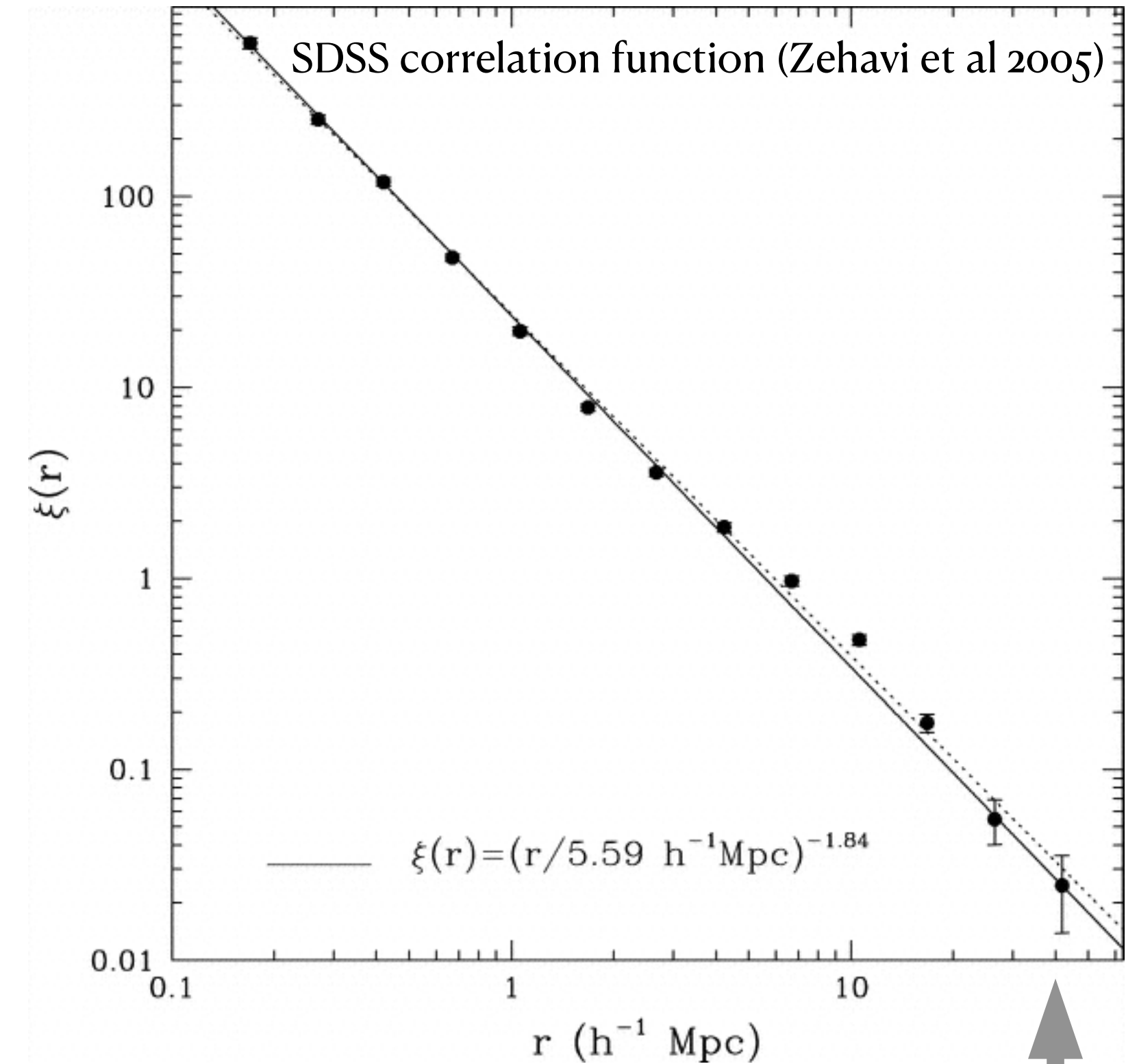
Quantified with the **correlation function** $\xi(r)$ which is the Fourier transform of the **power spectrum** $P(k)$.

The correlation function is the excess probability of finding a galaxy near another galaxy over that in a random distribution.

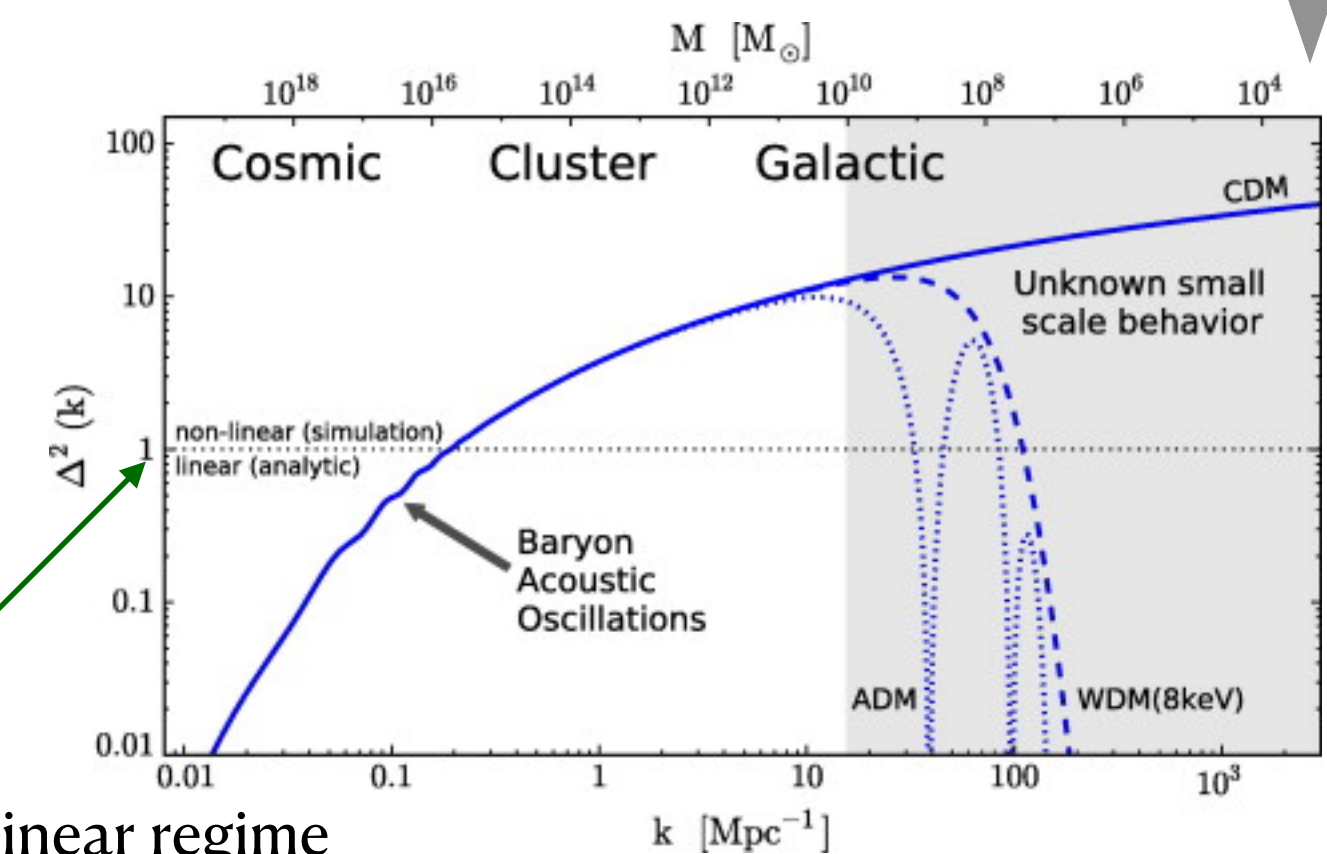
$$\frac{dN}{N} = [1 + \xi(r)]dV \quad \xi(r) = \frac{V}{(2\pi)^3} \int P(k) e^{-\vec{k} \cdot \vec{r}} d^3k$$

$$P(k) \propto |\delta(k)|^2 \propto k^n \quad \xi(r) \propto r^{-(n+3)}$$

Harrison-Zeldovich spectrum has $n = 1$, which is a Gaussian random field. Inflation predicts $n \approx 1$, but different flavors of Inflationary theory predict slightly different values depending on the shape of the Inflationary potential (the Inflaton). Planck measures $n = 0.965 \pm 0.004$



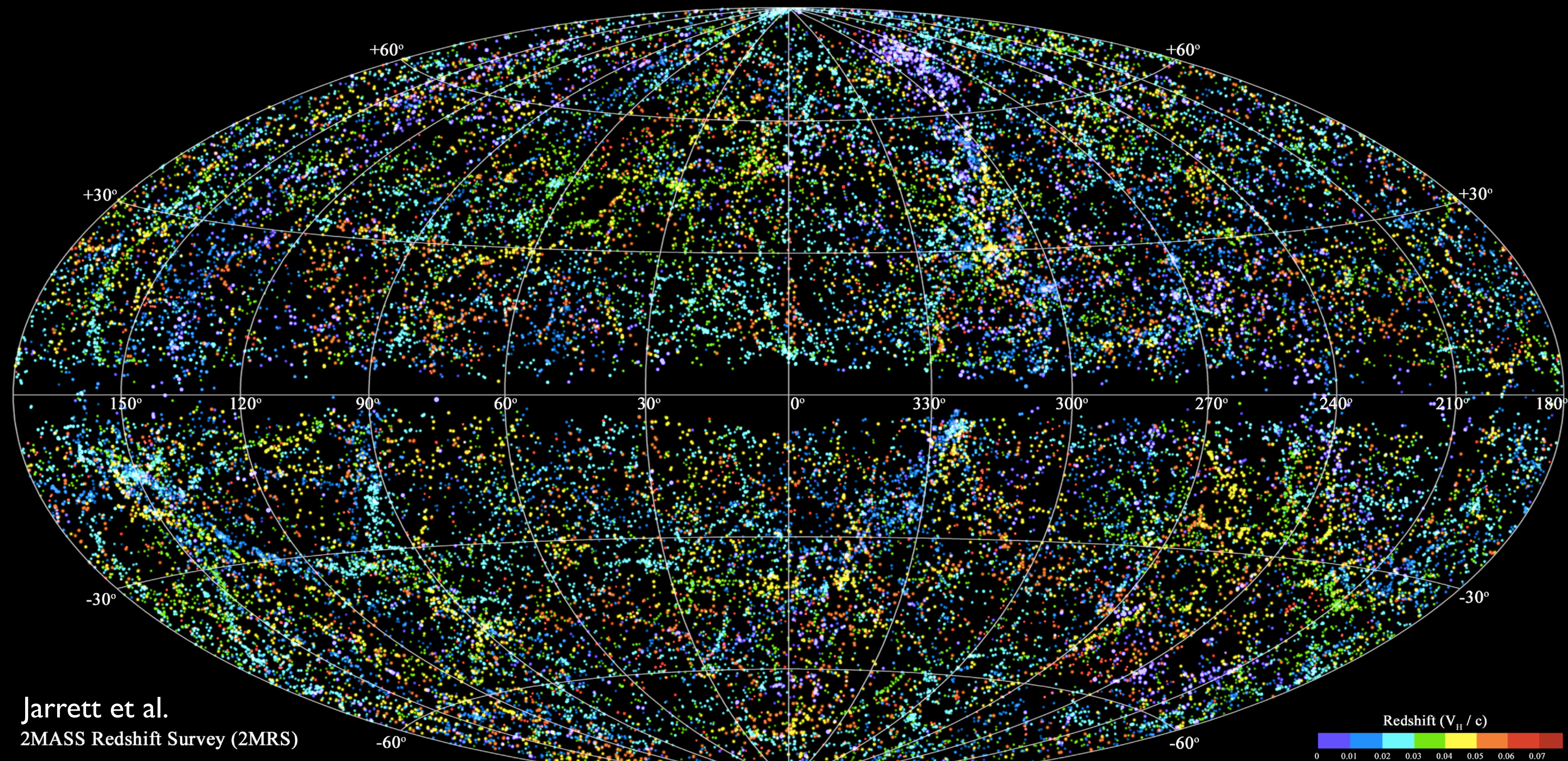
Power Spectrum



$\delta > 1$ marks the transition to the non-linear regime where perturbation theory no longer applies.

Measurements of the gravitating mass density

- Power spectrum of galaxies



- Power spectrum of galaxies

$$\delta \equiv \frac{\delta\rho}{\rho}$$

The power spectrum is commonly used to quantify large scale structure. It is related to the 2 point correlation function via Fourier transform.

2 point correlation function: $\xi(r) = \langle \delta(\vec{x}) \cdot \delta(\vec{x} + \vec{r}) \rangle$

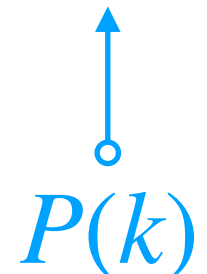
The 2 point correlation function is the probability of finding one galaxy near another in excess over a random distribution.

Power spectrum: $P(k) = \langle |\delta_k|^2 \rangle$ where $k = \frac{2\pi}{\lambda}$

where k is the wavenumber corresponding to the scale λ

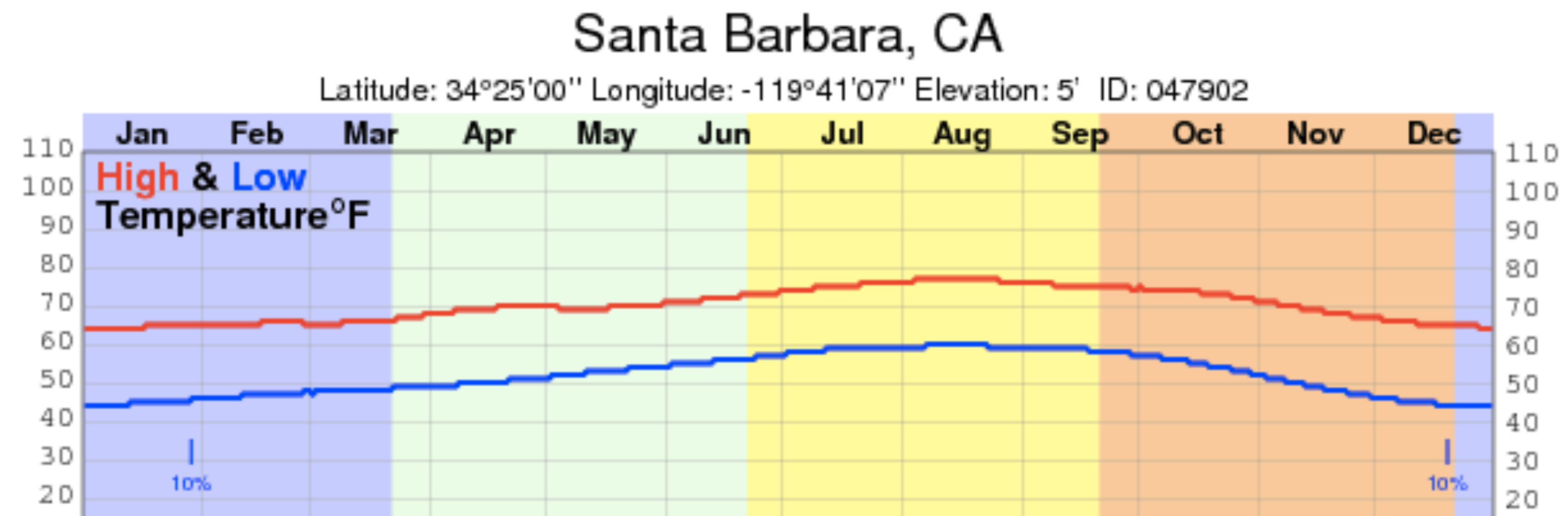
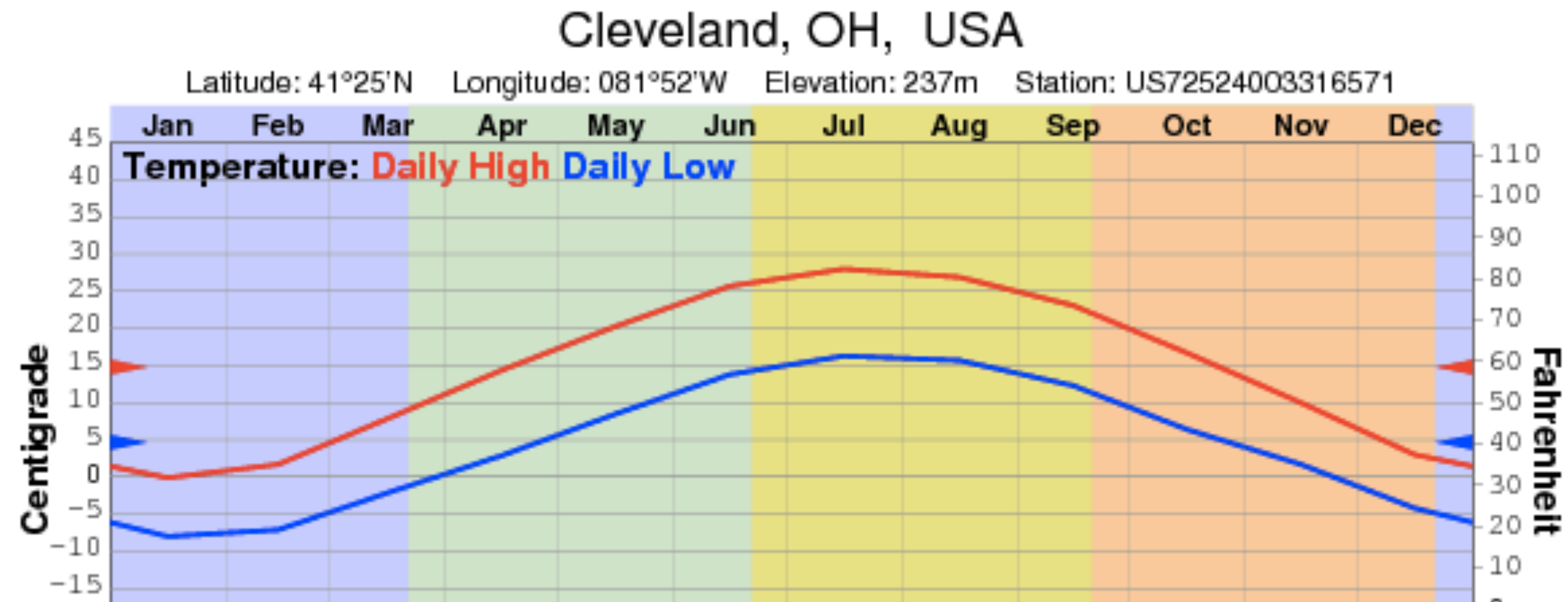
Fourier transform:

$$\xi(\vec{r}) = \frac{V}{(2\pi)^3} \int |\delta_k|^2 e^{-i\vec{k}\cdot\vec{r}} d^3k \quad \text{averaged over volume } V$$


 $P(k)$

Power Spectrum

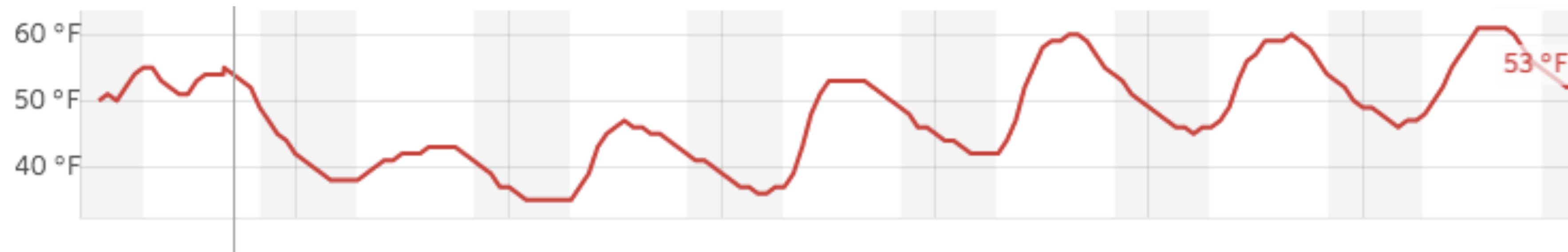
Example: weather in Cleveland and Santa Barbara
More power on long time scales in Cleveland (seasonal variation)



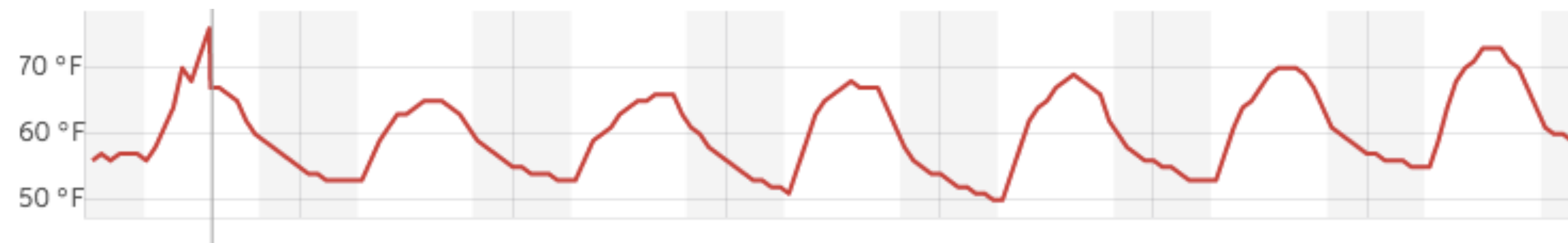
Power Spectrum

Example: weather in Cleveland and Santa Barbara
Similar power on short time scales in Santa Barbara (diurnal variation)

Cleveland forecast



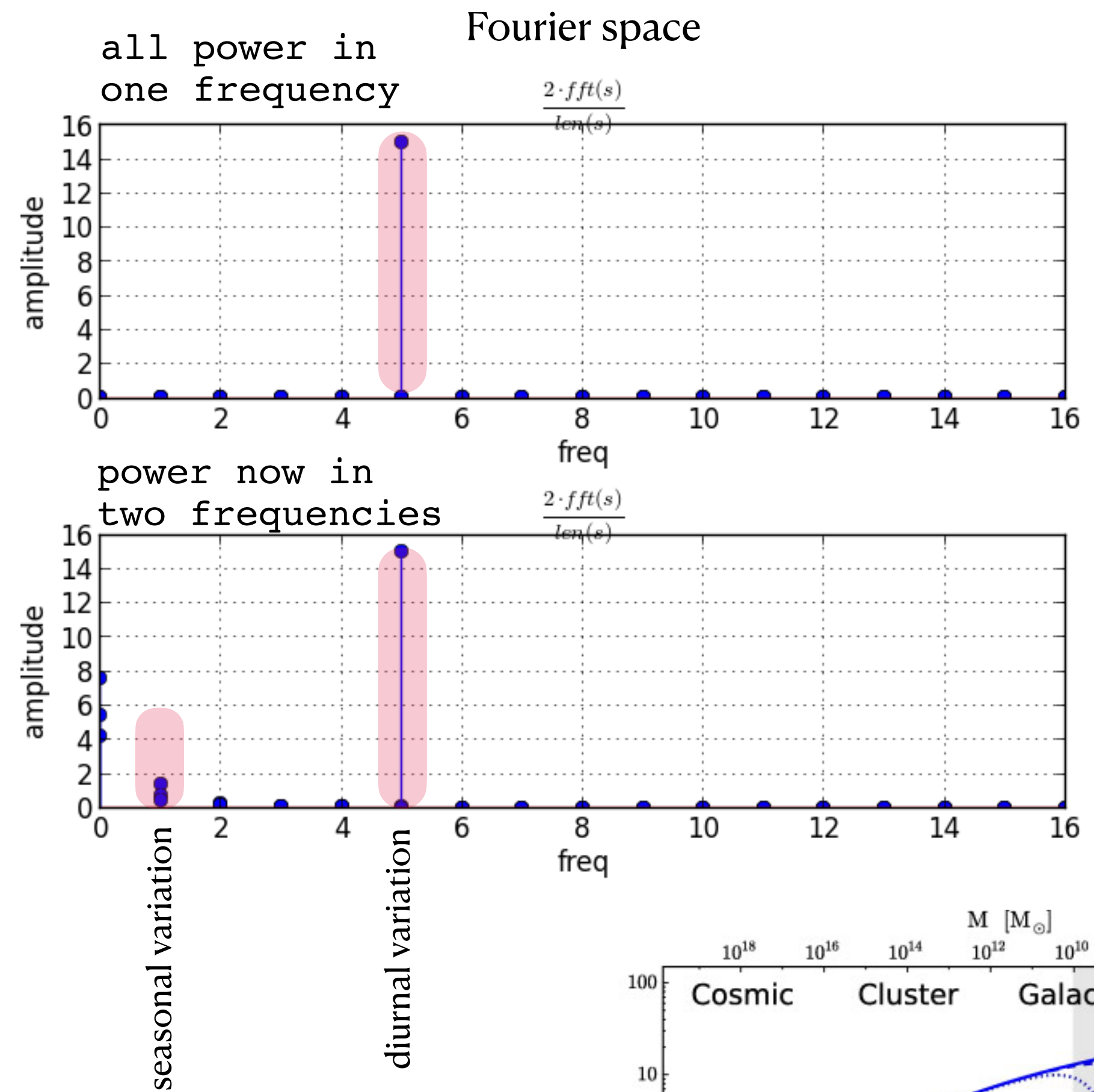
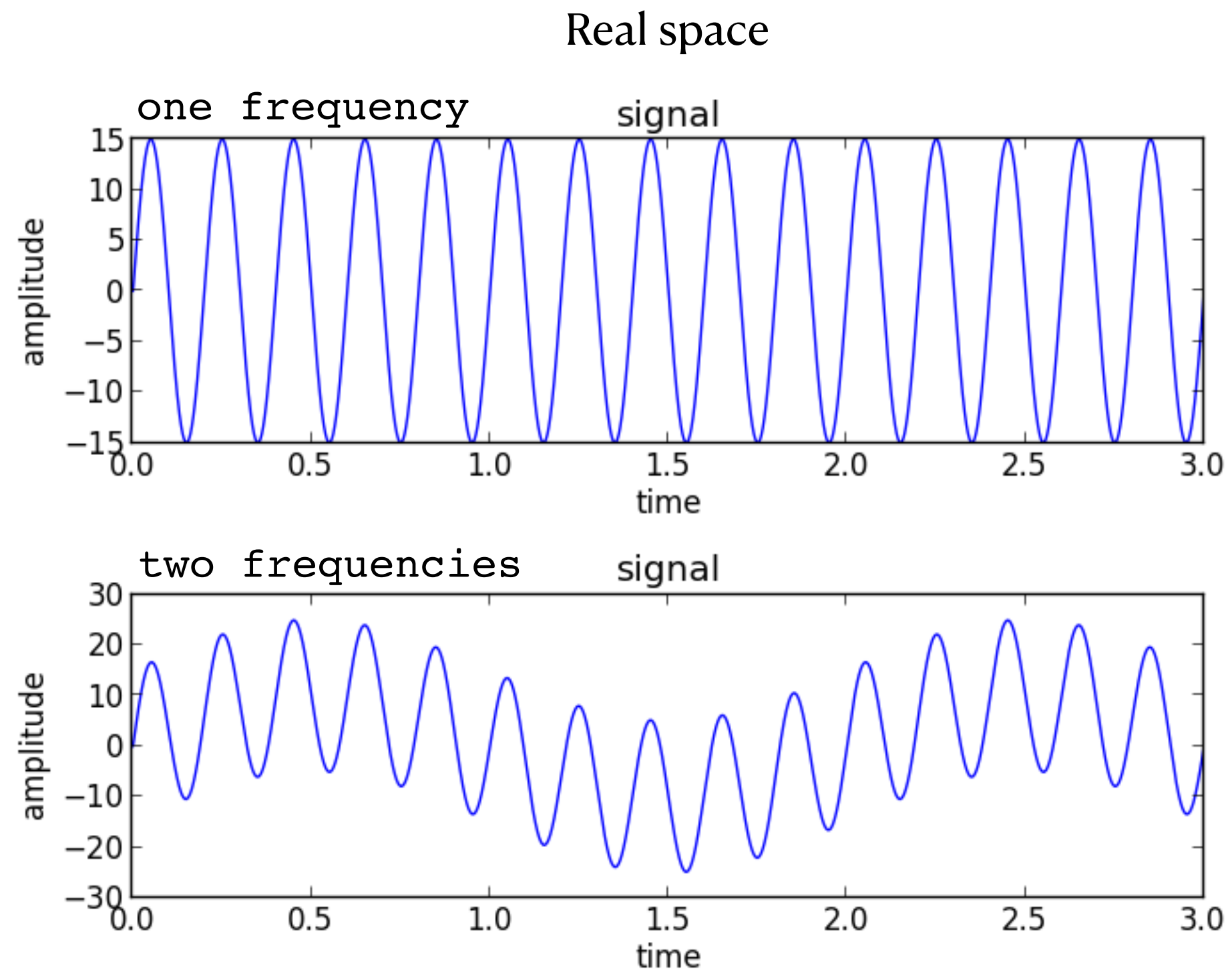
Santa Barbara forecast



A power spectrum is a Fourier transform that quantifies the relative variability on different scales

Superposition of two sinusoids

(e.g., diurnal and annual temperature variation)



So a smooth power spectrum has contributions from all frequencies, but also picks out which are more common.

