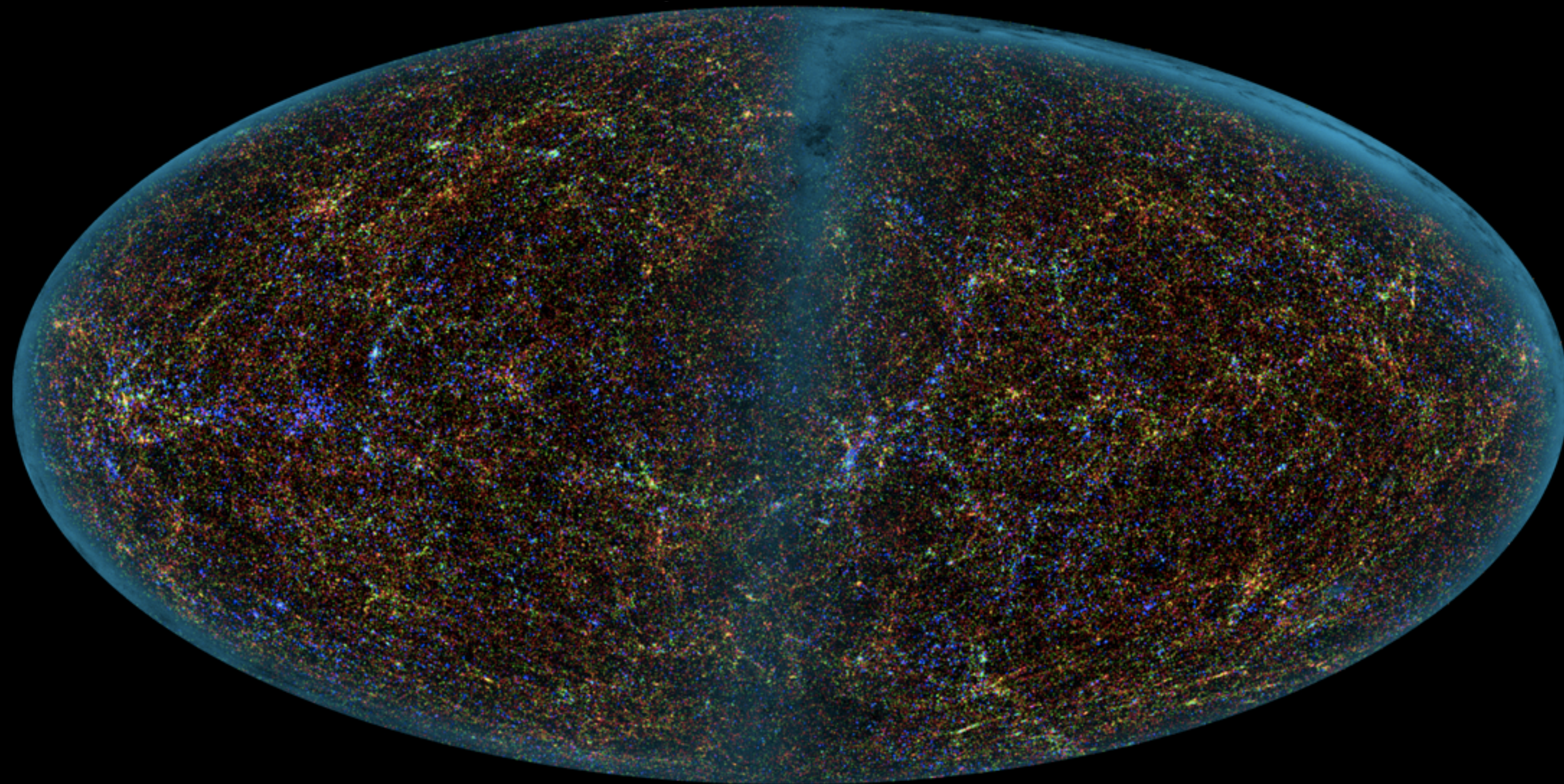


Cosmology

and Large Scale Structure



Today
Model Universes
Friedmann Eqn

First problem set due

Notional distances:

Hubble's "constant" is the current expansion rate $H = \frac{\dot{a}}{a} = \frac{1}{a} \frac{da}{dt}$

Proper distance: real separation in km or Mpc or furlongs.

The proper distance between galaxies increases as the universe expands.

$$D_{\text{proper}} = a(t) d_{\text{comoving}}$$

Comoving distance: separation on a grid of coordinates that expands along with the universe.

The comoving separation remains fixed as the universe expands.

Luminosity distance: equivalent to inverse-square distance in a Euclidean geometry.

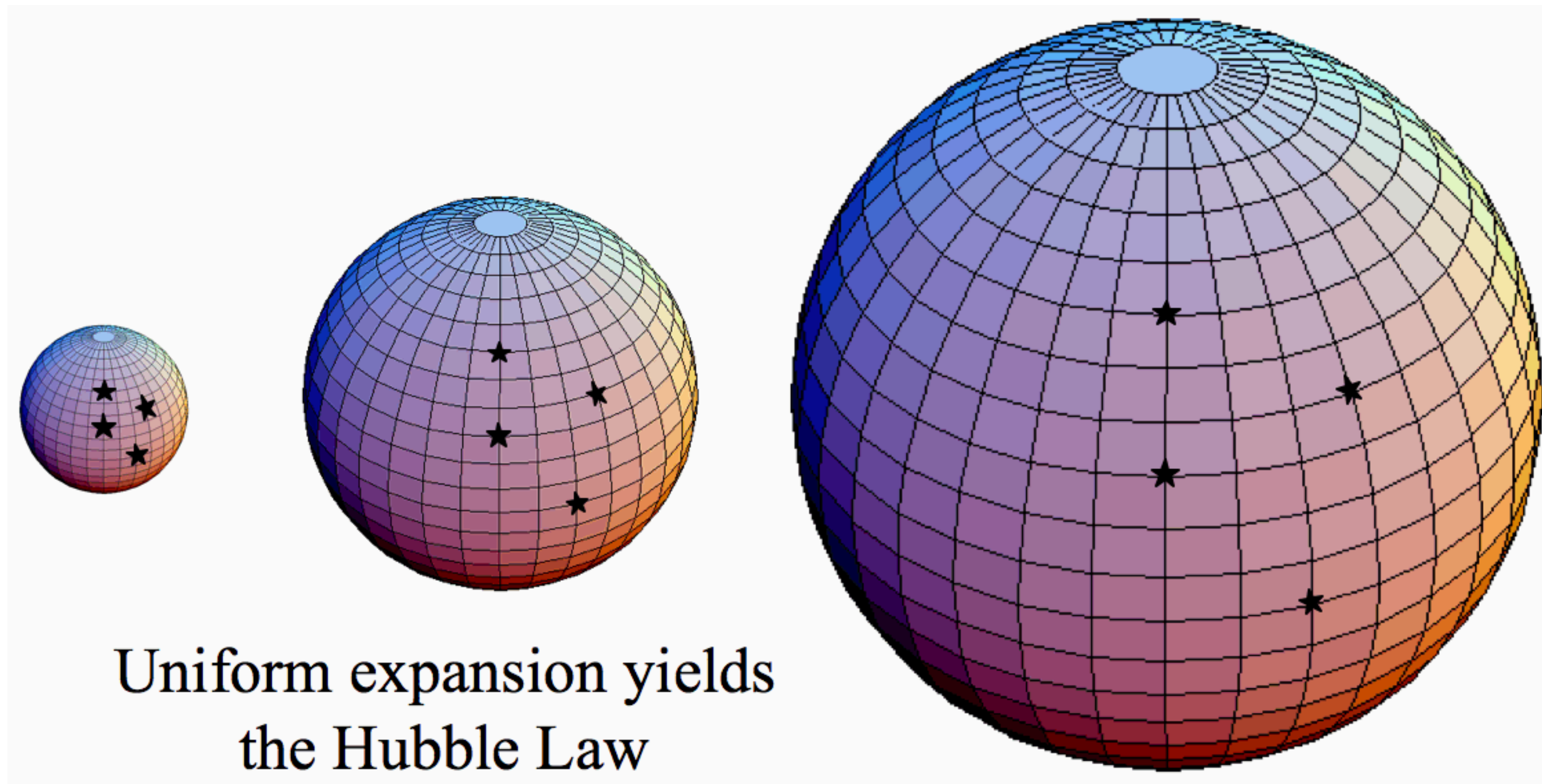
measured flux $f = \frac{L}{4\pi d_L^2}$

Angular-diameter distance: equivalent angular scale to that in a Euclidean geometry.

$$d_L = (1 + z) D_{\text{proper}}$$

angular size $\theta = \frac{\ell}{d_A}$

$$d_A = \frac{D_{\text{proper}}}{(1 + z)}$$



To get the proper distance to a galaxy we observe, we need to integrate over the expansion since the time of photon emission:

$$c dt = a(t) dr$$

$$c \int_{t_e}^{t_0} \frac{dt}{a(t)} = \int_0^r dr = r$$

where we make use of the fact that for photons, $ds = 0$

we know the expansion factor from the redshift

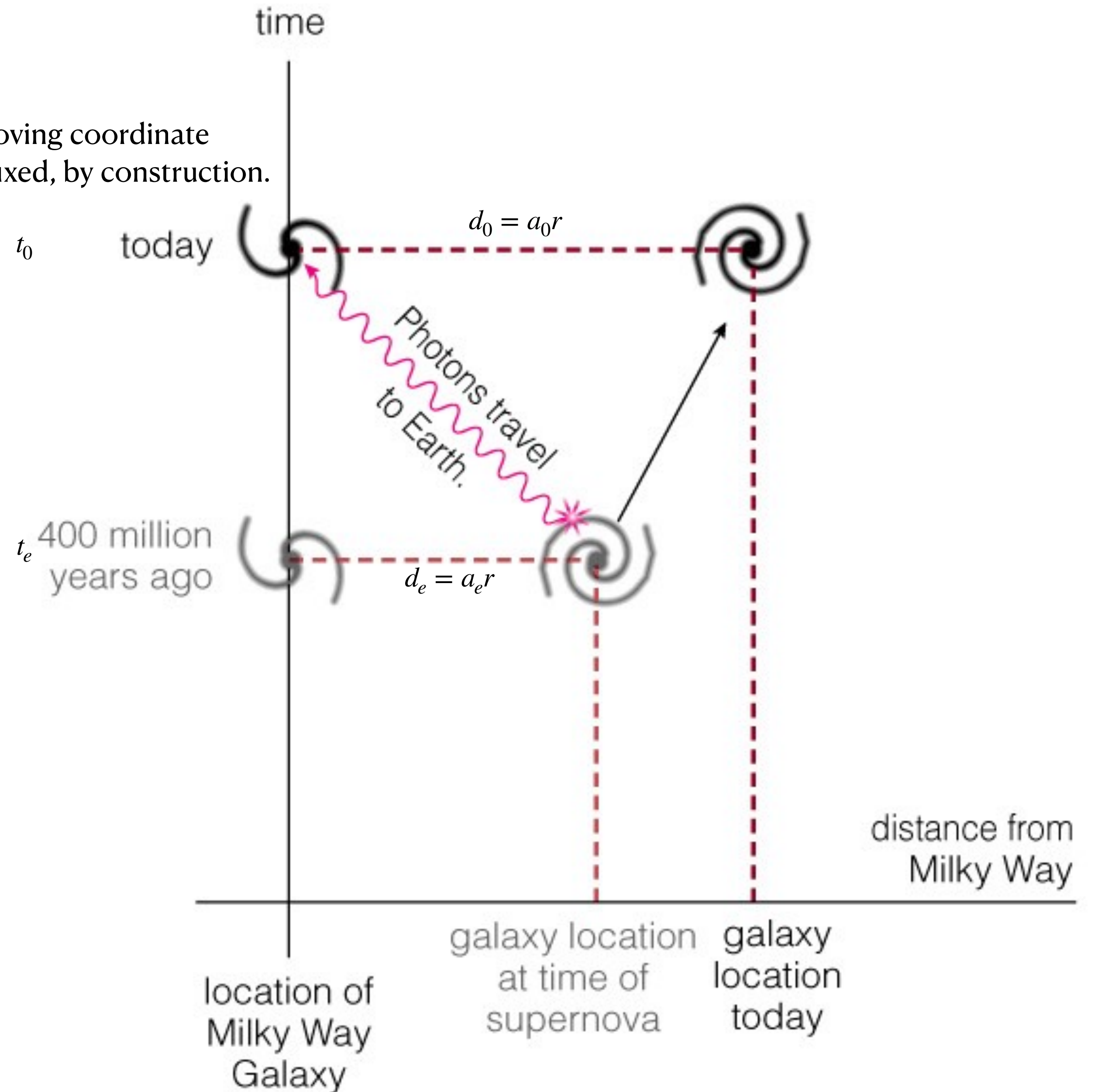
$$\frac{a(t_0)}{a(t_e)} = \frac{1}{1+z}$$

t_0 is now, so by construction $a(t_0) = 1$.

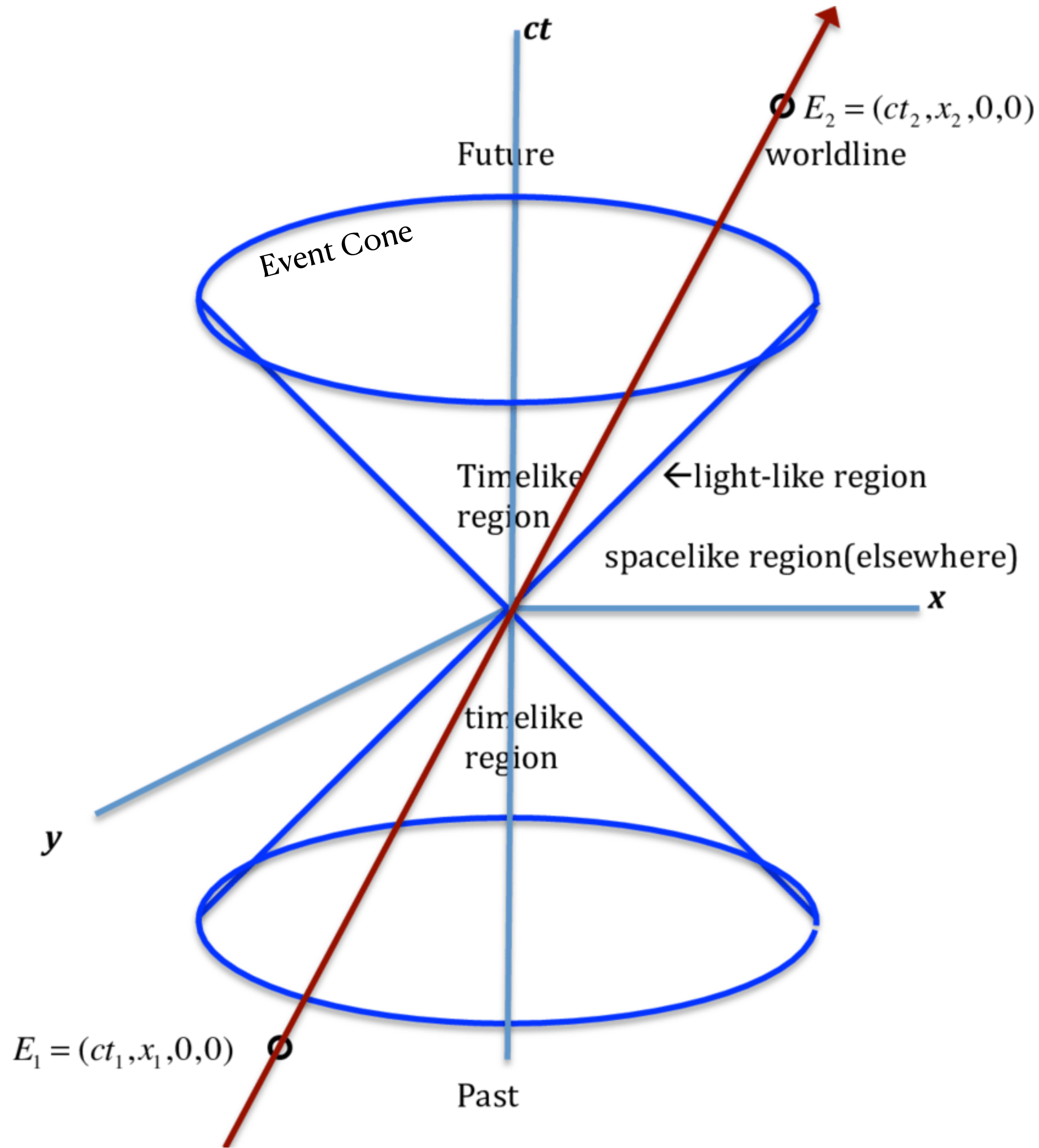
Cosmological parameters specify $a(t)$ through solution of the Friedmann equation

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{kc^2}{(aR_0)^2} + \frac{c^2}{3}\Lambda$$

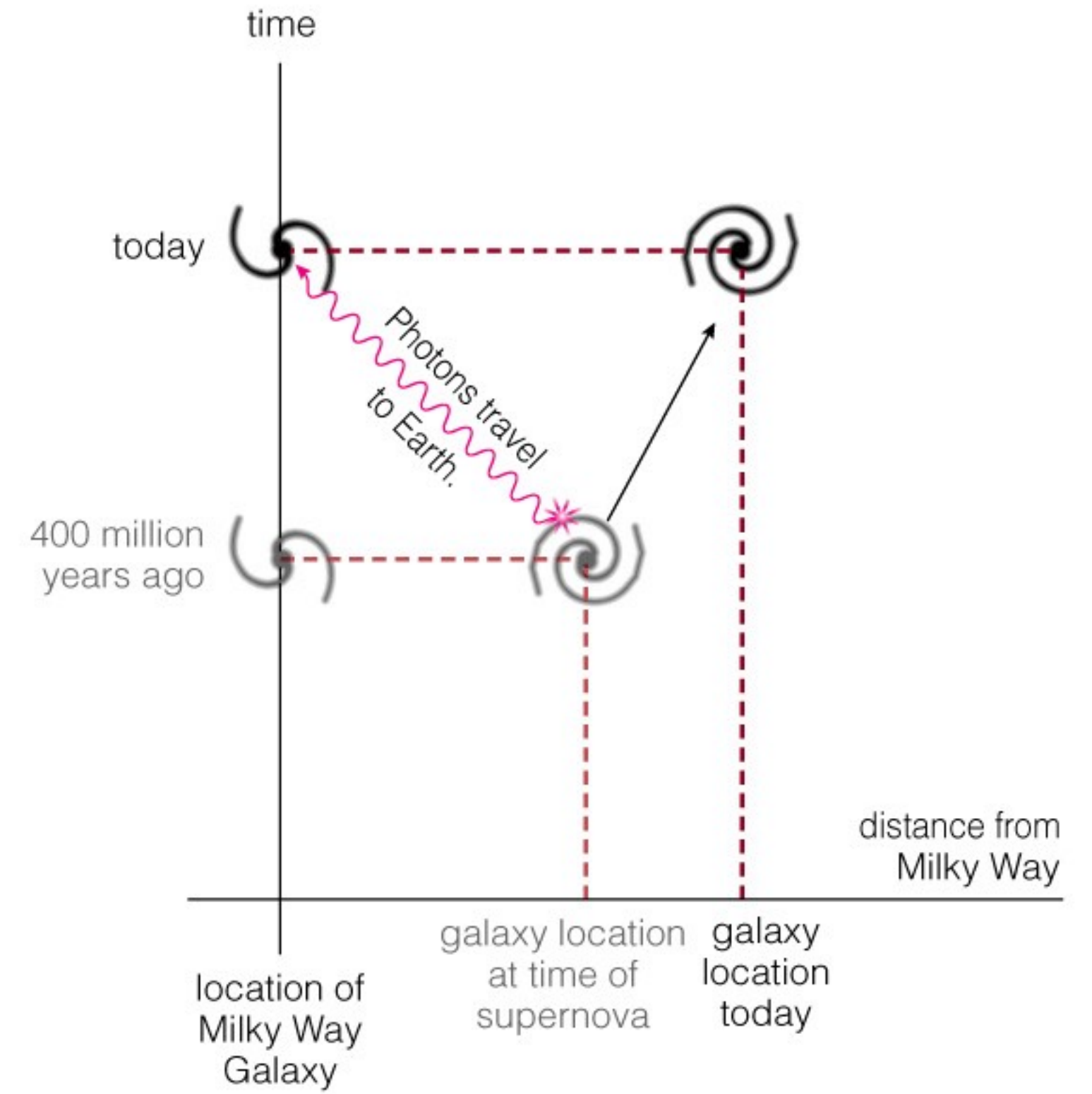
The comoving coordinate separation is fixed, by construction.



Spacetime diagram



Spacetime diagram for two galaxies



$$c^2 dt^2 > d\ell^2$$

Time like region - cone of causal connectivity

$$c^2 dt^2 = d\ell^2$$

Light like cone - traversed by photons in vacuum

$$c^2 dt^2 < d\ell^2$$

Space like region - out of causal contact

Model Universes

governed by

Einstein field equation

$$\mathcal{R}_{\mu\nu} - \frac{1}{2}\mathcal{R}g_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu} + \Lambda g_{\mu\nu}$$

which bequeath us the

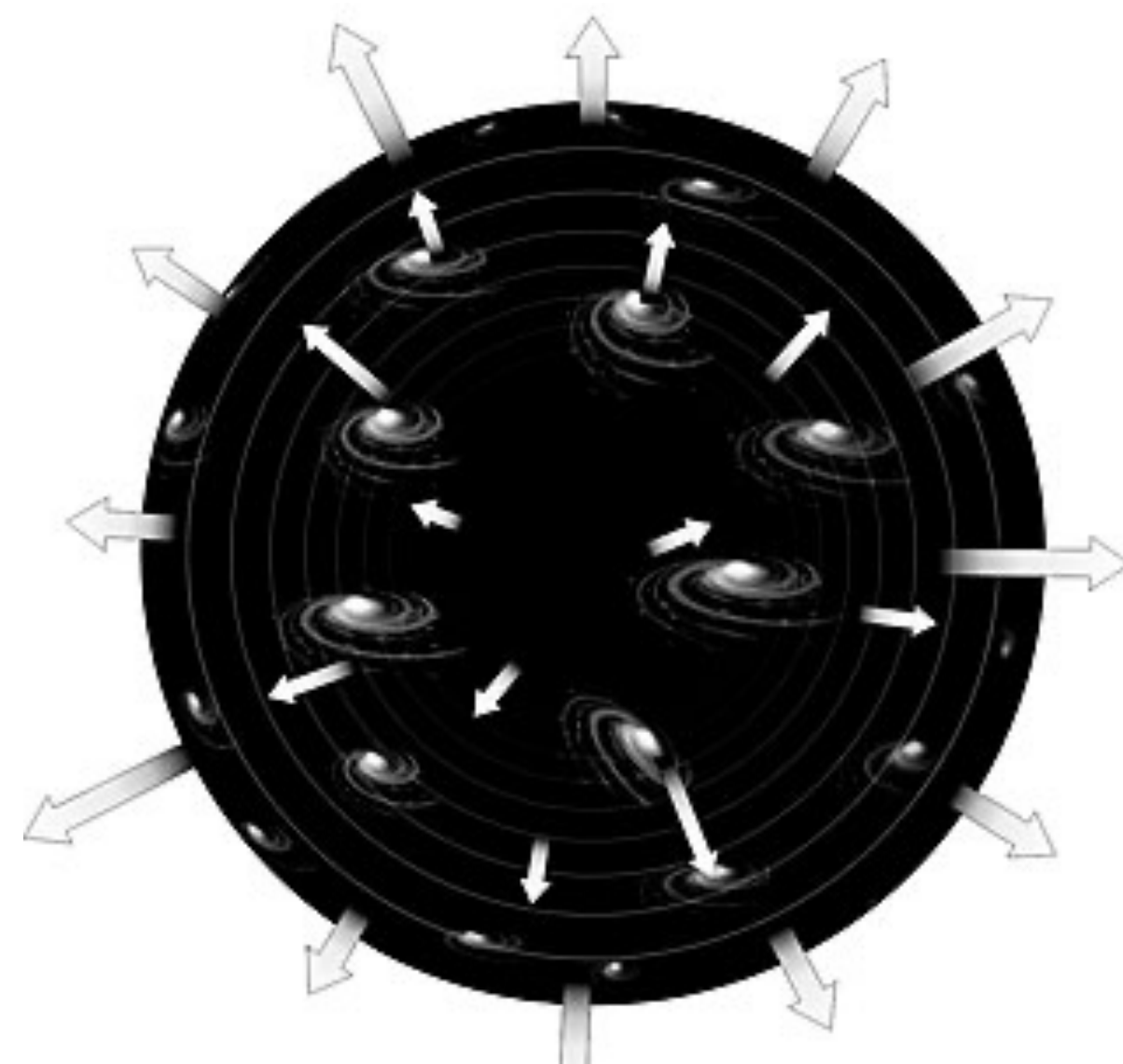
Roberston-Walker metric

$$ds^2 = -c^2 dt^2 + a^2(t)[dr^2 + S_k^2(r)d\Omega^2]$$

mostly just care about
 $c dt = a(t) dr$
 for events tied to the comoving
 coordinate system

and the

Friedmann equation



$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{kc^2}{(aR_0)^2} + \frac{c^2}{3}\Lambda$$

expansion rate

gravitating
mass-energy

geometry

R_0 is the radius of curvature

cosmological constant

Friedmann equation from GR

Ignoring the cosmological constant for the moment,

$$\mathcal{R}_{\mu\nu} - \frac{1}{2}\mathcal{R}g_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu}$$

$$T_{\mu\nu} = \begin{bmatrix} \rho & & 0 \\ & P & \\ 0 & & P \end{bmatrix}$$

Ricci Tensor

Stress-Energy Tensor

no rotationally invariant 4-vectors

$$\mathcal{R}_{ti} = 0$$

$$T_{ti} = 0$$

curvature of space related to the gravitational potential

$$\mathcal{R}_{tt} = U(t)$$

$$T_{tt} = \rho(t)$$

mass density Poisson equation in Newtonian gravity

$$\nabla^2\Phi = 4\pi G\rho$$

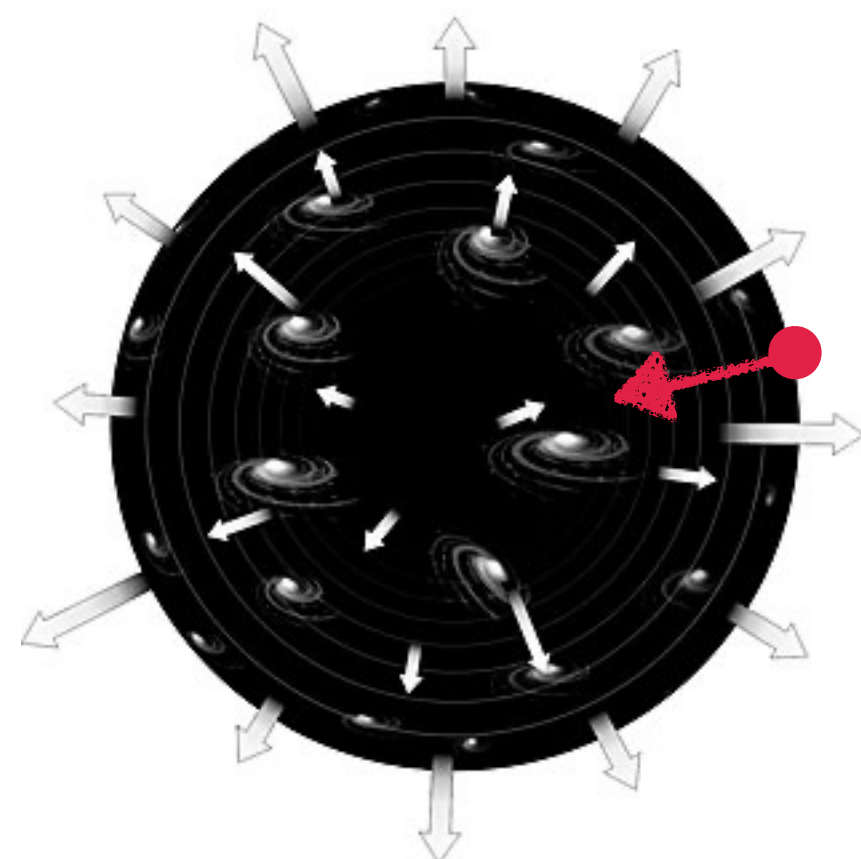
energy gravitates too

$$\mathcal{R}_{ij} = g_{ij}V(t)$$

$$T_{ij} = g_{ij}P(t)$$

pressure stemming from the energy density in relativistic components

Friedmann equation from GR



$$U(t) = -3\frac{\ddot{a}}{a} = 8\pi G\left(\frac{1}{2}\rho + \frac{3}{2}P\right)$$

The units of density and pressure are taken to be the same here; they are related by c^2 in conventional units — pressure is an energy density.

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P)$$

Looks like the Newtonian acceleration of a point on the surface of a sphere with a term added for pressure.

Aside: Newtonian Cosmology

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\rho$$

$$\ddot{a} = -\frac{GM}{a^2}$$

$$\ddot{a} = -\frac{4\pi G}{3}\rho a \quad \text{for} \quad M = \frac{4\pi G}{3}\rho$$

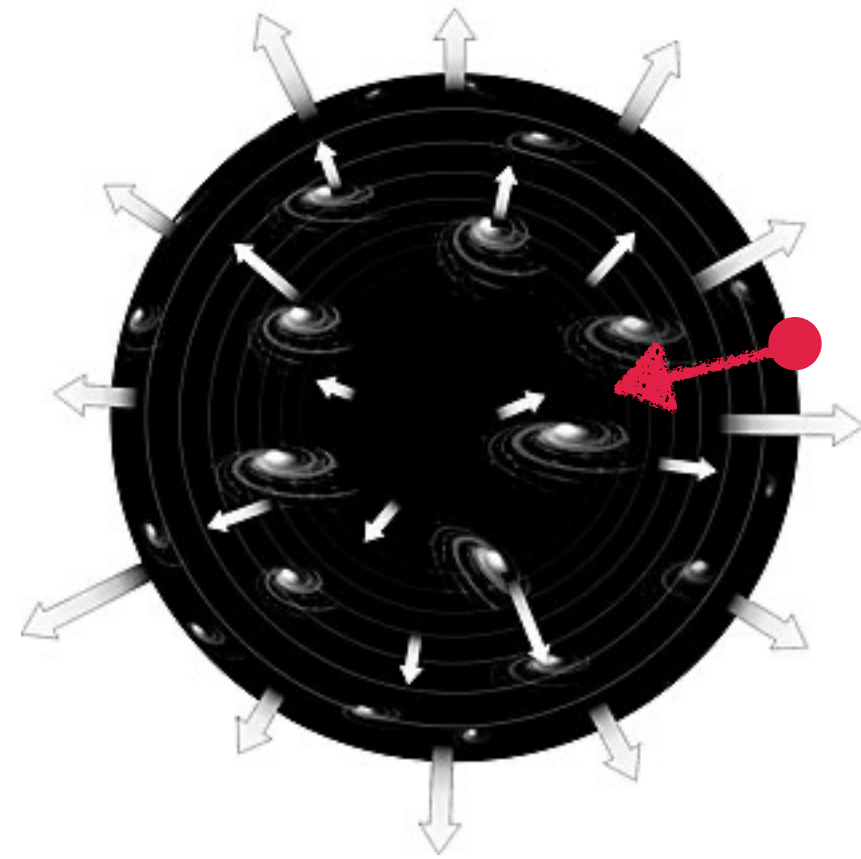
ρ uniform, by assumption (homogeneity)

Pick up integration constant with units of energy, U . This determines whether the expansion is “bound” or not - i.e., the sphere recollapses or expands forever.

$$\dot{a}^2 - \frac{8\pi G}{3}\rho a^2 = 2U$$

The behavior is the same for any uniform sphere, so its size ℓ is irrelevant

Friedmann equation from GR



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Looks like the Newtonian acceleration of a point on the surface of a sphere with a term added for pressure.

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\rho$$

$$\mathcal{R}_{ij} = g_{ij}V(t)$$

$$V(t) = a\ddot{a} + 2\dot{a} + 2k = 4\pi G(\rho - P)a^2$$

Combine equations for $U(t)$ and $V(t)$ to eliminate the second time derivative, leaving the Friedmann equation

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{kc^2}{(aR_0)^2}$$

The book replaces mass density with energy density (eq. 4.20)

$$\rho = \frac{\varepsilon}{c^2}$$

Friedmann equation

Including the cosmological constant,
the Friedmann equation becomes

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{kc^2}{(aR_0)^2} + \frac{c^2}{3}\Lambda$$

Energy conservation $D_\mu T_{\mu\nu} = 0$ in an expanding universe

$$\frac{d}{da}(a^3\rho) = -a^2P \quad \text{just } PdV \text{ work}$$

total mass-energy “work” done in expansion

Solutions depend on the equation of state

$$P = w\rho$$

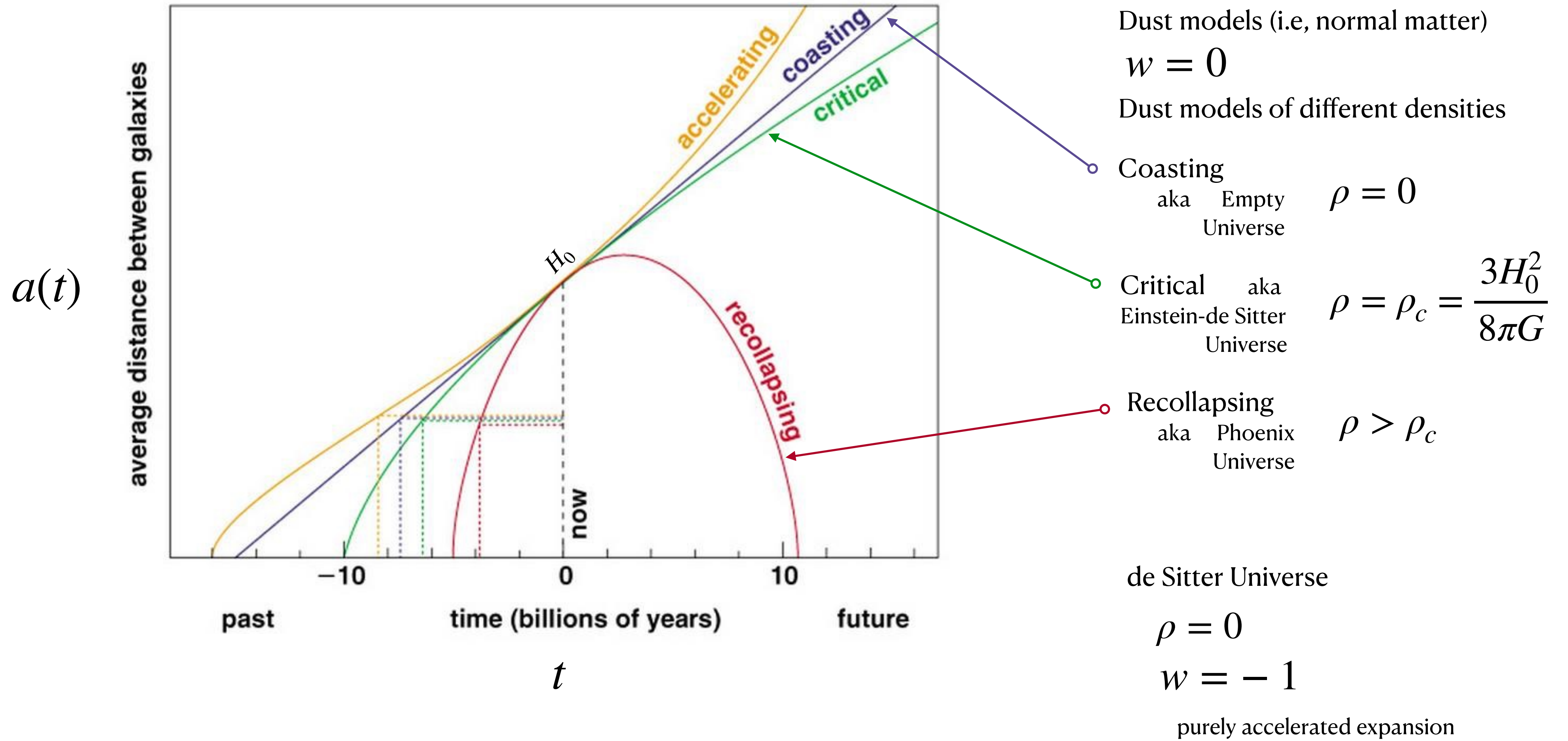
$w = 0$ non-relativistic mass (“dust”)

$w = \frac{1}{3}$ photons

$w = -1$ cosmological constant

The cosmological constant is repulsive because of the negative sign in its equation of state.

Possible expansion histories



The expansion started by the Big Bang is resisted by the attraction of gravity.

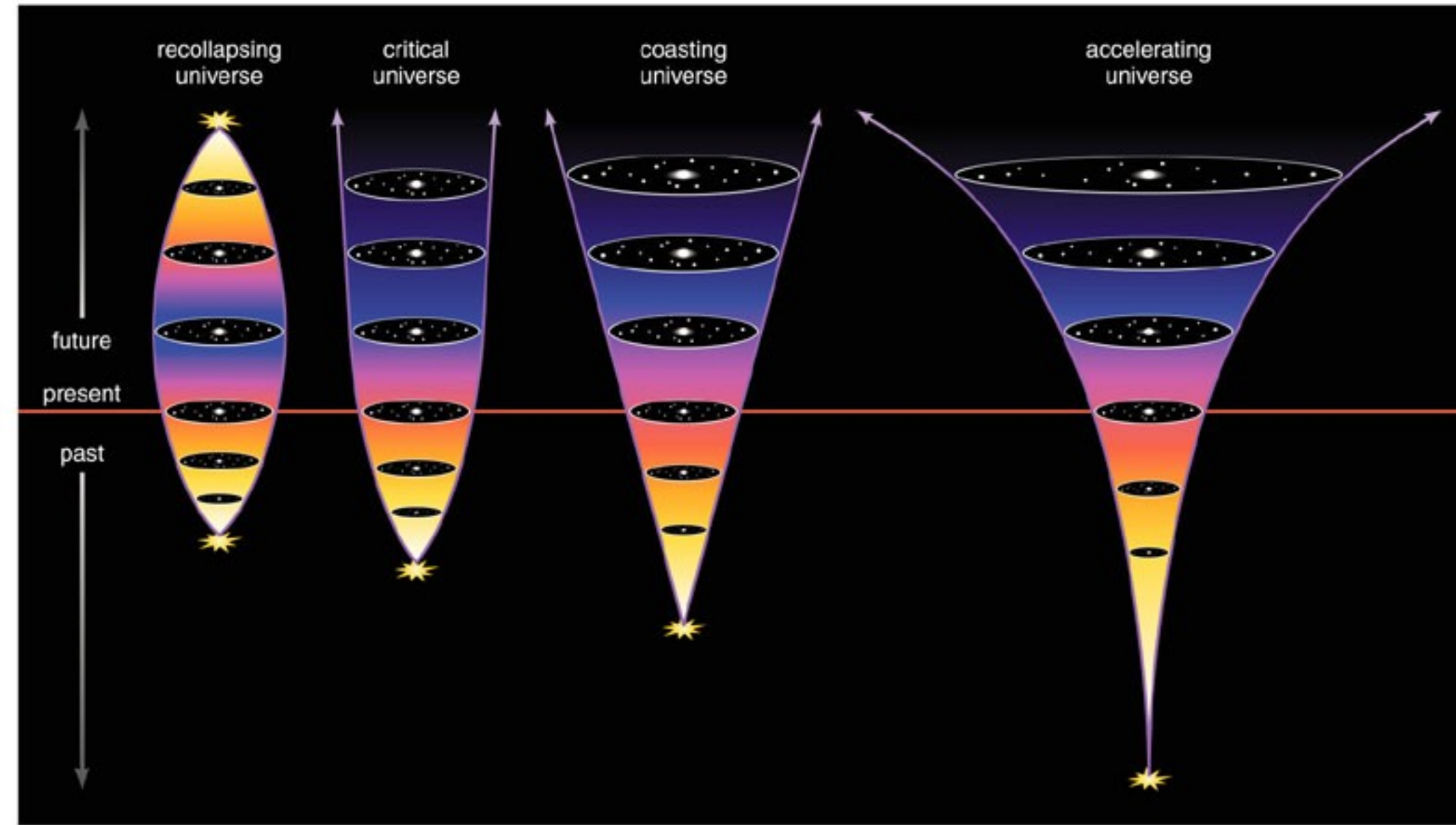
The more dense the universe, the more gravity... a balance is reached at a critical density:

	IF	the universe is
$k = -1$	$\rho < \rho_{crit}$	OPEN: expands forever
$k = 0$	$\rho = \rho_{crit}$	FLAT
$k = +1$	$\rho > \rho_{crit}$	CLOSED: eventually re-collapses

These are the traditional options in the absence of a repulsive force like a cosmological constant/dark energy

Density is destiny

It determines the age, geometry, and ultimate fate of the universe



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CLOSED

CRITICAL

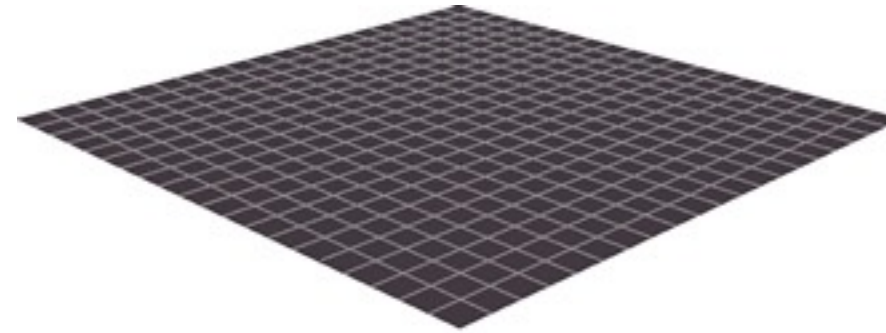
OPEN

ACCELERATING

$$k = 0$$

FLAT

Density = Critical

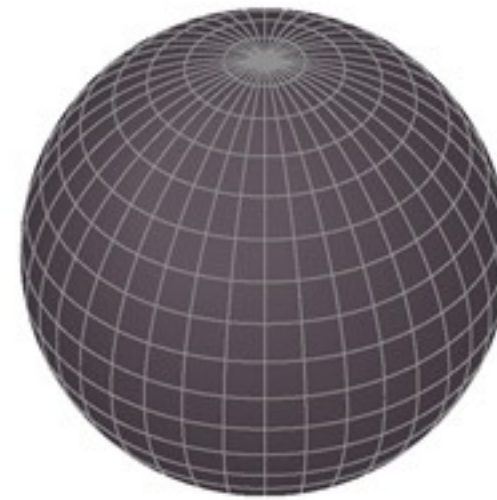


flat (critical) geometry

$$k = + 1$$

CLOSED

Density > Critical

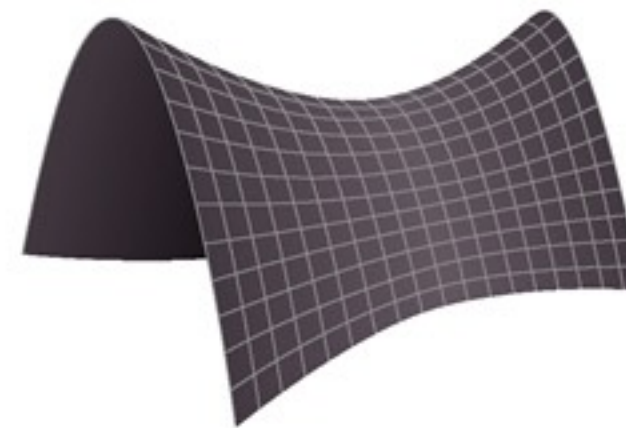


spherical (closed) geometry

$$k = - 1$$

OPEN

Density < Critical



saddle-shaped (open) geometry

Space can be curved.

The overall geometry of the universe is determined by the total density of matter and energy.

Friedmann equation

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{kc^2}{(aR_0)^2} + \frac{c^2}{3}\Lambda$$

where

$$\Omega_m = \frac{8\pi G}{3H^2}\rho$$

mass density

$$\Omega_r = \frac{\varepsilon c^{-2}}{\rho_c} \quad \text{radiation density}$$

$$\Omega_k = -\frac{kc^2}{(aR_0H)^2}$$

curvature

Flat cosmologies have $k = 0$ so $\Omega_k = 0$

$$\Omega_\Lambda = \frac{c^2\Lambda}{3H^2}$$

cosmological constant

Λ is constant but Ω_Λ evolves as H evolves

can be written

$$H^2 = H^2(\Omega_m + \Omega_r + \Omega_k + \Omega_\Lambda)$$

$H \equiv \frac{\dot{a}}{a}$ does not remain constant, so the Hubble “constant” is just the current value of the Hubble parameter $H(z)$.

the sum of density parameters so defined must be unity:

$$\Omega_m + \Omega_k + \Omega_\Lambda = 1$$

Expansion dynamics

The Acceleration equation with the cosmological constant:

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P) + \frac{1}{3}\Lambda$$

$$H = \frac{\dot{a}}{a}$$

$$a = (1+z)^{-1}$$

can usually be replaced with a single variable, as $P = w\rho$ for a single medium.

The Pressure P is zero when matter dominates.
It is simply related to the energy density when radiation dominates.

You can see why the cosmological constant leads to acceleration!

$$\ddot{a} \sim \Lambda$$

Solutions depend on the equation of state

$$P = w\rho$$

$w = 0$ non-relativistic mass ("dust")

$w = \frac{1}{3}$ photons

$w = -1$ cosmological constant