Today:

- MOND
- Review

# Next time: Exam





#### MOND

Modified Newtonian Dynamics (Milgrom 1983)

Instead of invoking dark matter, modify gravity (or inertia). Milgrom suggested a modification at a particular acceleration scale  $a_0$ 

Newtonian regimeMOND regime $a = g_N$  for  $a \gg a_0$  $a = \sqrt{g_N a_0}$  for  $a \ll a_0$ 



Regimes smoothly joined by

$$\mu\left(\frac{a}{a_0}\right)a = g_N$$

 $\mu(x) \to 1 \text{ for } x \gg 1$   $\mu(x) \to x \text{ for } x \ll 1$   $x = \frac{a}{a_0}$ 

Modified Poisson equation

$$\nabla \left[ \mu \left( \frac{\nabla \Phi}{a_0} \right) \nabla \Phi \right] = 4\pi G \rho$$

Derived from aquadratic Lagrangian of Bekenstein & Milgrom (1984) to satisfy energy conservation.





1983 Milgrom

MODIFICATION OF NEWTONIAN DYNAMICS

No. 2, 1983

A major step in understanding ellipticals can be made if we can identify them, at least approximately, with idealized structures such as the FRCL spheres discussed above. I have also studied isotropic and nonisotropic isothermal spheres, in the modified dynamics, as such possible structures. I found that they have properties which the properties to complete those of empticals and

#### VIII. PREDICTIONS

The main predictions conce lows.

 Velocity curves calculated with the modified dynamics on the basis of the observed mass in galaxies should agree with the observed curves. Elliptical and S0 galaxies may be the best for this purpose since (a) practically no uncertainty due to obscuration is involved and (b) there is not much uncertainty due to the possible presence of molecular hydrogen.

2. The relation between the asymptotic velocity  $(V_m)$ and the mass of the galaxy (M)  $(V_m^4 = MGu_0)$  is an absolute one.

3. Analysis of the z-dynamics in disk galaxies using the modified dynamics should yield surface densitie which error with the loss of one councilings of the same as with the galaxies of one councilings of the same as with the same as a statistic of the same as a same as a statistic of the same as a statistic of th

4. Effects of the be particularly stro A review of property & a 1980). For example, those dwarfs believed to be bound to our Galaxy would have internal accelerations typically of order ain - an/30. Their (modified) acceleration, g, in the field of the Galaxy is larger than the internal ones but still much smaller than  $a_0, g = (8)$ kpc/d ) $a_0$ , based on a value of  $V_{w} = 220$  km s<sup>-1</sup> for the Galaxy, and where d is the distance from the dwarf Jalaxy to the center of the Milky Way (d = 70-220kpc). Whichever way the external acceleration turns out to affect the internal dynamics (see the discussion at the end of § II, the section on small groups in Paper III, and Paper I), we predict that when velocity dispersion data is available for the dwarfs, a large mass discrepancy will result when the conventional dynamics is used to determine the masses. The dynamically determined mass is predicted to be larger by a factor of order 10 or more than that which can be accounted for by stars. In case the internal dynamics is determined by the external acceleration, we predict this factor to increase with d

and be of order (d/8 kpc) (as long as a<sub>in</sub> ≪ g, h<sub>30</sub> = 1). Prediction 1 is a very general one. It is worthwhile listing some of its consequences as separate predictions, numbered 5-7 below (note that, in fact, even prediction - is already contained in prediction 1). 5. Measuring local M/L values in disk galaxies (assuming conventional dynamics) should give the following results: In regions of the galaxy where  $V^2/r \gg a_0$ the local M/L values should show no indication of hidden mass. At a certain transition radius, local M/L should start to increase rapidly. The transition radius

Canonation of M/L as we are concerned only with variation of M/L as we are concerned only with variations of this quantity; (b) Effects of the modified dynamics manifest themselve more clearly in ford d a miny m is a set of the matrix and d is

ior in the 1sk only while the spheroid can be neglected. This makes the determination of mass from velocity more certain.

6. Disk galaxies with low surface brightness provide particularly strong tests (a study of a sample of such galaxies is described by Strom 1982 and by Romanishin et al. 1982). As low surface brightness means small accelerations, the effects of the modification should be more noticeable in such galaxies. We predict. for example, that the proportionality factor in the M  $\propto V^4$  relation for these galaxies is the same as for the high surface tensity grides. In the rate is the high surface of the strong strong test is the same as for the high surface tensity grides. In the rate of the high surface of the strong strong test is the average surface brightness. This is been as unpletened. As a way of the relation of the strong strong test is the average surface brightness. This is plies that low surface tensity galaxies. We also predict that the lower the average surface density of a galaxy is, the smaller is the transition rate. As predict that the lower the average surface density of a galaxy is, the smaller is the transition rate. As predict that the lower the average surface density of a galaxy is, the smaller is the transition rate. As predict that the lower the average surface density of a galaxy is, the smaller is the transition rate. As predict that the lower the average surface density of a galaxy is, the smaller is the transition rate. As predict that the lower the average surface density of a galaxy is, the smaller is the transition rate. As predict that the lower the average surface density of a galaxy is, the smaller is the transition rate. As predict that the lower the average surface density of a galaxy is average surface tensity of scale long. As some the starting to increase from very small we may have the starting to increase from very small radii.

7. As the study of model rotation curves shows, we predict a correlation between the value of the average surface density (or brightness) of a galaxy and the efferness with which the rotational velocity rises to its asymptotic value (as measured, for example, by the radius at which  $V = V_{\rm sc}/2$  in units of the scale length of the disk). Small surface densities imply slow rise of V.

#### IX. DISCUSSION

The main results of this paper can be summarized by the statement that the modified dynamics eliminates the need to assume hidden mass in galaxies. The effects in galaxies which I have considered, and which are commonly attributed to such hidden mass, are readily explained by the modification. More specifically:

### MOND predictions

• The Tully-Fisher Relation

#### • Slope = 4 • Slo

ary strong langerally arelation between Diskelass and V<sub>flat</sub>

> • No Dependence on Surface Brightness

**OWING** Automotion Husters and Surface of conventional **O**/**Dop** radius and surface brightness **SB galaxies were widely** • Rotation Curve Shapes

not to exist.

- Surface Density ~ Surface Brightness
- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios



• The Tully-Fisher Relation



Brightness

- Dependence of conventional M/L on radius and surface brightness
- Rotation Curve Shapes
- Surface Density ~ Surface Brightness
- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios

#### In MOND limit of low acceleration

$$a = \sqrt{g_N a_0}$$



$$V^4 = a_0 GM$$

#### observed TF!

#### Why? Physics of the BTFR scaling relation

dark matter

halos:  $M_{tot} \propto V^3$ 

baryons:  $M_d \propto V^x$ 

 $x \geq 3$  depending on  $m_d(V)$ 

Should depend on disk scale length, unless all disks submaximal

Should work as long as object not tidally disrupted

 $\frac{\text{MOND}}{M_{tot}} = M_b = \frac{V^4}{a_0 G}$ an absolute consequence of the force law for a  $\langle a_0 \rangle$ :  $g_N = \mu\left(\frac{g}{a_0}\right)g$ Newtonian regime:  $\mu \rightarrow 1$  for  $g \gg a_0$  so  $g = g_N$ MOND regime:  $\mu \rightarrow g/a_0$  for  $g \ll a_0$  so  $g = \sqrt{g_N a_0}$ 

Should only work for objects in MOND regime



- The Tully-Fisher Relation
  - Slope = 4
    Normalization = 1/(a<sub>0</sub>G)
    Fundamentally a relation between Disk Mass and V<sub>flat</sub>
    - No Dependence on Surface Brightness
- Dependence of conventional M/L on radius and surface brightness
  - Rotation Curve Shapes
  - Surface Density ~ Surface Brightness
  - Detailed Rotation Curve Fits
  - Stellar Population Mass-to-Light Ratios



• The Tully-Fisher Relation

Slope = 4
Normalization = 1/(a<sub>0</sub>G)
Fundamentally a relation between Disk Mass and V<sub>flat</sub>

- No Dependence on Surface Brightness
- Dependence of conventional M/L on radius and surface brightness



- Surface Density ~ Surface Brightness
- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios



•	MOND predictions The Tully-Fisher Relation
	<ul> <li>Slope = 4</li> <li>Normalization = 1/(a<sub>0</sub>G)</li> <li>Fundamentally a relation between Disk Mass and V<sub>flat</sub></li> <li>No Dependence on Surface Brightness</li> </ul>
•	Dependence of conventional M/L on radius and surface brightness
/.	Rotation Curve Shapes
•	Surface Density ~ Surface Brightness
•	Detailed Rotation Curve Fits

• Stellar Population Mass-to-Light Ratios

















#### Residuals of MOND fits



Famaey, B., & McGaugh, S.S. 2012, Living Reviews in Relativity, 15, 10



• The Tully-Fisher Relation

Slope = 4
Normalization = 1/(a<sub>0</sub>G)
Fundamentally a relation between Disk Mass and V<sub>flat</sub>

 No Dependence on Surface Brightness

Dependence of conventional M/L on radius and surface brightness

Rotation Curve Shapes

• Surface Density ~ Surface Brightness

- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios



#### Line: stellar population model (mean expectation)







• The Tully-Fisher Relation

Slope = 4
Normalization = 1/(a<sub>0</sub>G)
Fundamentally a relation between Disk Mass and V<sub>flat</sub>

 No Dependence on Surface Brightness

Dependence of conventional M/L on radius and surface brightness

- Rotation Curve Shapes
  - Surface Density ~ Surface Brightness
  - Detailed Rotation Curve Fits
  - Stellar Population Mass-to-Light Ratios





Newtonian	regime	MOND regime	
$g_{in} > a_0$	$M = \frac{RV^2}{G}$	$g_{in} < a_0$	$M = \frac{V^4}{a_0 G}$
	e.g., surface of the Earth	e.g., remote dwarf Leo I	© Anglo-Australian Observatory
External Field Newtonian	d dominant regime	Exter quasi-	nal Field dominant -Newtonian regime
External Field Newtonian $g_{in} < a_0 < g_{ex}$	d dominant regime $M = \frac{RV^2}{G}$	Exter quasi- $g_{in} < g_{ex} <$	nal Field dominant -Newtonian regime $a_0 \qquad M = \frac{a_0}{g_{ex}} \frac{RV^2}{G}$

#### A new test: the dwarf satellites of Andromeda





Velocity dispersions of the dwarf satellites of Andromeda







Pairs of photometrically identical dwarfs should have different velocity dispersion depending on whether they are isolated are dominated by the external field effect.



There is no EFE in dark matter - this is a unique signature of MOND.



### Tidal Debris Dwarfs - should be devoid of Dark Matter



Bournaud et al. (2007) Science, 316, 1166





### I find your lack of faith disturbing.

- You don't know the Power of the Dark Side
- Can MOND explain large scale structure?
- Can it provide a satisfactory cosmology?
- Can it be reconciled with General Relativity?



# Questions

• Don't know answers

