

Example feedback scheme (Dutton 2009 MNRAS, 396, 121)

Supernova drives outflows, which are assumed to move at the local escape velocity to minimize mass removal. ($< v_{esc}$ doesn't escape
 $> v_{esc}$ moves less mass for same energy)

Energy driven wind model:

$$\Delta M_{\text{eject}}(R) = \frac{2 E_{\text{EFB}} \eta_{\text{SN}} E_{\text{SN}} \Delta M_{\text{SN}}(R)}{v_{\text{esc}}^2(R)}$$

↑
←

ejected mass from radius R mass of stars formed at radius R

$$E_{\text{SN}} \approx 10^{51} \text{ erg} = 5 \times 10^7 \text{ km}^2 \text{ s}^{-2} M_{\odot}$$

$$\eta_{\text{SN}} = 8.3 \times 10^{-3} \text{ \# SN per } M_{\odot} \text{ of stars formed (this \# for a Chabrier IMF)}$$

E_{EFB} = fraction of kinetic energy injected into wind

usually a large # in simulations (0.25 - 1)

usually a small # observed (0.02 - 0.1)

Momentum driven wind model

$$\Delta M_{\text{eject}}(R) = \frac{E_{\text{MFB}} P_{\text{SN}} \eta_{\text{SN}} \Delta M_{\text{SN}}(R)}{v_{\text{esc}}(R)}$$

$P_{\text{SN}} = 3 \times 10^4 M_{\odot} \text{ km s}^{-1}$ is momentum produced by one SN

E_{MFB} is again the coupling efficiency to the ISM

This formulation maximizes the impact of SN.

Luminosity function & Missing Satellite Problem

Schechter Fun:

$$\Phi(M) = \Phi_* e^{-\left(\frac{M}{M_*}\right)} \left(\frac{M}{M_*}\right)^\alpha \quad \# \text{ of galaxies per Mpc}^3 \text{ per mass}$$

3 parameters: Φ_* characteristic density (Mpc^{-3})
 M_* characteristic mass / luminosity
 α faint end slope

traditionally used to fit the galaxy luminosity fun
 captures gross feature; can over-sm some details
 may vary with environment (cluster LF steeper than field LF)

Bill et al (2003)

find

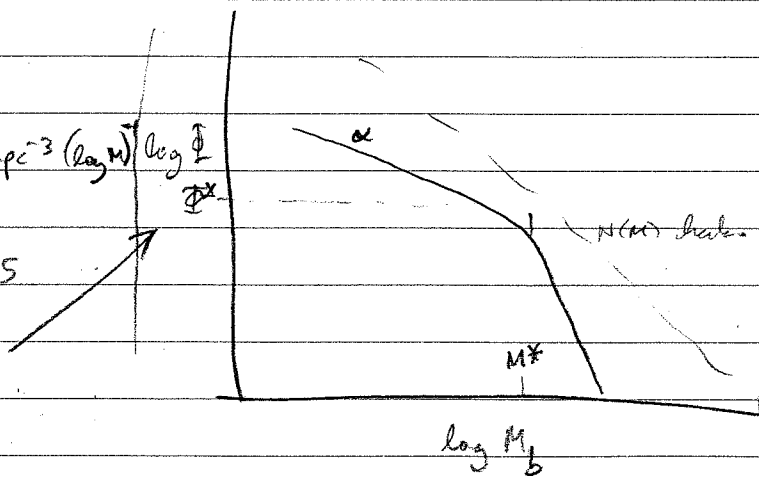
$$\Phi_* = 4 \times 10^{-3} \text{ Mpc}^{-3} (\log M) \log \Phi$$

$$M_* = 10^{11} M_\odot$$

$$\alpha = -1.21 \pm 0.05$$

Sometimes $\frac{dN}{dM}$

sometimes $\frac{d \log N}{d \log M}$



Waller, Ostriker & Bode (2005)

give $M_* = 1.1 \times 10^{15} M_\odot$

$$\alpha = -1.9$$

for CDM halo mass fun

Integrated luminosity density

$$j = \int_0^\infty L \Phi(L) dL$$

$$j = L_* \Phi_* \Gamma(\alpha + 2)$$

Abundance matching: map $\overset{\text{expected}}{\Phi(M_{\text{halo}})}$ to $\overset{\text{observed}}{\Phi(M_*)}$

The expected halo abundance is usually determined by matching the clustering of DM halos in simulations to that observed for large galaxy survey like SDSS.

Leads to $M_* - M_{\text{halo}}$ relations like that of Moster et al. (2010):

$$\frac{m}{M} = 2 \left(\frac{m}{M} \right)_0 \left[\left(\frac{M}{M_1} \right)^{-\beta} + \left(\frac{M}{M_1} \right)^{\gamma} \right]^{-1}$$

where m is stellar mass and M is halo mass.

$\left(\frac{m}{M} \right)_0$, M_1 , β , & γ are fit parameters of a broken power law

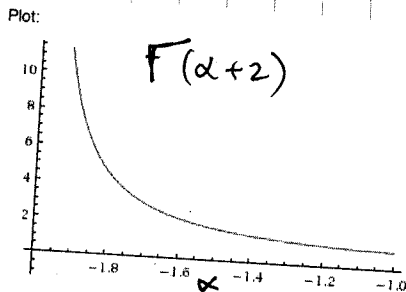
Moster et al (2010)

$$\log M_1 = 11.884$$

$$\left(\frac{m}{M} \right)_0 = 0.0282$$

$$\beta = 1.057$$

$$\gamma = 0.556$$



Note $\Gamma(\alpha+2) = 1+2$ for $\alpha = -1 \rightarrow -1.5$
 blows up as $\alpha \rightarrow -2$
 (> 10 for $\alpha < -1.9$;
 goes to ∞ at $\alpha \rightarrow -2$)

$$j = L^* \Phi^* \Gamma(\alpha+2)$$