

## The Oort Discrepancy

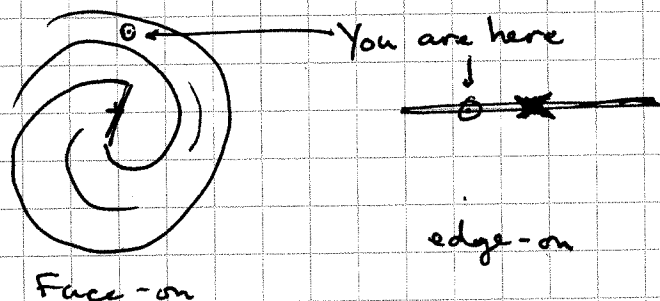
The Milky Way, like most spiral galaxies, is a thin, dynamically cold, rotating disk of stars & gas.

Thin:  $R_d: z_0 = 8:1$  is typical, with considerable scatter

Dynamically Cold:  $v/\sigma \gg 1$

For the Milky Way,  $v \approx 200 \text{ km s}^{-1}$   
 $\sigma \approx 20 \text{ km s}^{-1}$

This is typical. Understanding the stability & persistence of thin spiral disks is a challenge, but the universe is littered with them.



To a decent first approximation, the orbit of a star in the Milky Way is approximately circular, ~~with~~ ~~in~~ in the plane of the disk, while executing harmonic oscillation in the  $z$ -direction.

Oort (1932) considered the  $z$ -motions of stars in the solar neighborhood, extending and applying earlier work by Kapteyn and Jeans.

Poisson equation

$$\nabla^2 \Phi = 4\pi G \rho$$

$\Phi$  = gravitational potential

$\rho$  = 3D density ( $M_{\odot} \text{pc}^{-3}$ )

In cylindrical coordinates  $R, \theta, z$

$$\nabla^2 \Phi = \frac{1}{R} \frac{\partial}{\partial R} \left( R \frac{\partial \Phi}{\partial R} \right) + \frac{1}{R^2} \frac{\partial^2 \Phi}{\partial \theta^2} + \frac{\partial^2 \Phi}{\partial z^2}$$

Assume axis-symmetry, so  $\frac{\partial \Phi}{\partial \theta} = 0$ . Not really valid if spiral arms are strong (or bars)

integrate once over  $z$

$$\Sigma = \int_{-|z|}^{|z|} \rho dz$$

for distributions symmetric in  $z$  [ $\rho(z) = \rho(-z)$ ]

this is equivalent to  $\int_0^{|z|}$  times 2,

$$\text{so } 4\pi G \rho \rightarrow 2\pi G \Sigma |z|$$

assume separable:

radial force  $K_R = -\frac{\partial \Phi}{\partial R} = -\frac{v^2}{R}$

vertical force  $K_z = \frac{\partial \Phi}{\partial z} = \int \frac{\partial^2 \Phi}{\partial z^2} dz$

so now we've eliminated one derivative to have

$$2\pi G \Sigma = |K_z| - |z| \frac{1}{R} \frac{\partial}{\partial R} (R K_R)$$

$$R K_R = R \cdot \frac{v^2}{R} = v^2$$

so

$$K_z = 2\pi G \Sigma + \frac{z}{R} \frac{\partial v^2}{\partial R}$$

restoring force to plane

Surface density of plane w/  $z < |z|$   
depends on disk surface density

radial force term amounts to square of rotation velocity also depends on DM halo

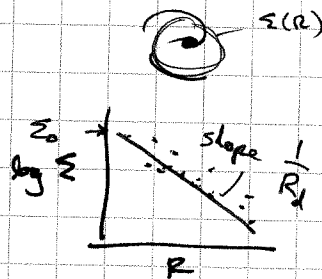
By convention, vertical force has been measured at  $|z| = 1.1 \text{ kpc}$ , which contains most of the disk mass, but still has stars left to measure

## Double exponential model

The azimuthally averaged light profile of spiral galaxies can to a level be approximated by an "exponential disk"

$$\Sigma(R) = \Sigma_0 e^{-R/R_d}$$

$\Sigma_0$  central surface brightness  
 $R_d$  disk scale length



This describes the 2D face-on distribution of star light. The  $z$ -distribution can also be approximated as exponential, so in 3D

$$\rho(R, z) = \rho_0 e^{-R/R_d} e^{-z/z_0}$$

We can use this to integrate the Poisson equation and replace  $K_z$  with the vertical velocity dispersion  $\sigma_z^2$

The result is

$$\sigma_z^2 = K \pi G z_0 \Sigma_0$$

Very local approximation  
Get cross terms if  $\delta R$  large

where the constant  $K$  depends on the mass distribution.

Usually it is taken to be  $K=2$  (~~isothermal slab~~) isothermal

or  $K=1.5$  (pure exponential vertical  $dz$ )

In general, this applies to all components - every population of stars (A V, K III, etc) as well as gas, and dark matter.

The constant  $K$  can be determined numerically for arbitrary  $\Sigma(z)$ , but for realistic cases  $1.5 \leq K \leq 7$

Will return to this with the Jeans equations and the concept of the phase space. These give a different approach to individual tracer populations, as young stars are thinner & dynamically colder than older stellar populations

## MEASUREMENTS

There have been many attempts to measure the local surface density of the disk since Oort:  
e.g.,

Kuijken & Gilmore 1991

Holmberg & Flynn 2004

Bienaymé et al 2014

Bovy & Rix 2013

to  $z \approx 2$  kpc!

$K_{11}$  over large range in radius,  
not just locally!

numerically, the accounting goes (at the solar circle,  $R=R_0$ )

$$\begin{array}{l} \text{at } R_0 \\ \Sigma_{\star} = 38 M_{\odot} \text{pc}^{-2} \\ \Sigma_{\text{gas}} = 14 M_{\odot} \text{pc}^{-2} \\ \Sigma_{\text{dyn}} = 74 M_{\odot} \text{pc}^{-2} \end{array} \left. \vphantom{\begin{array}{l} \Sigma_{\star} \\ \Sigma_{\text{gas}} \\ \Sigma_{\text{dyn}} \end{array}} \right\} \Sigma_{\text{b}} = 52 M_{\odot} \text{pc}^{-2}$$
$$\Sigma_{\text{dyn}} = 74 M_{\odot} \text{pc}^{-2} > \Sigma_{\text{b}}$$

That  $\Sigma_{\text{dyn}} > \Sigma_{\text{b}} = \Sigma_{\star} + \Sigma_{\text{g}}$  is a mass discrepancy.

The dynamics wants more mass than is directly seen.

The difference  $- 22 M_{\odot} \text{pc}^{-2} -$  is presumably the integrated portion of a quasi-spherical dark matter halo between  $-1.1 \leq z \leq 1.1$  kpc.

For comparison, the local V-band surface brightness is  $\sim 27 L_{\odot} \text{pc}^{-2}$ . So the mass-to-light ratio of the stars of the Milky Way is a reasonable

$$\nu_{\star} = 1.4 M_{\odot}/L_{\odot}$$