### **Cosmology and Large Scale Structure**



**15 October 2024** http://astroweb.case.edu/ssm/ASTR328/

#### **Today** Age constraints Early Universe



- Oldest stars
	- Globular clusters
- White dwarfs
	- cooling curves & luminosity function
- Radioactive chronometers
	- Thorium/Europium ratio
- Interstellar dust grains
	- Oxygen isotope ratios

Expansion time scale set by the Hubble time  $H_0^{-1}$ . What about the ages of observed objects? The contents of the universe should not be older than the universe itself!

Sets lower limit on the age of the universe that is [just barely] consistent with the Hubble time  $H_0^{-1}$ .

$$
t_{GC} = 13.32 \pm 0.1 \text{ (stat)} \pm 0.5 \text{ (sys) Gyr}
$$
  

$$
t_{WD}^{max} = 12.5_{-3.5}^{+1.4} \text{ Gyr}
$$
  

$$
t_{*,\text{Th}} = 12.8 \pm 3 \text{ Gyr}
$$
  

$$
t_{MW} = 13.7 \pm 1.3 \text{ Gyr}
$$

Independent age estimates are pretty consistent

- White dwarfs
	- cooling curves & luminosity function

White dwarfs follow a simple radiative cooling curve — these cores of former stars fade as they cool in a predictable way. One therefore expects a truncation of the white dwarf luminosity function at the lowest luminosity to which they have had time to cool. That, plus the lifetime of the star that gave them birth, gives an age limit.



$$
L_{WD} = 4\pi R_{WD}^2 \sigma_{SB} T^4 = \frac{dE}{dt} = \frac{d}{dt} \left(\frac{3}{2} N k_B T\right) \rightarrow \frac{dT}{dt} \approx \frac{\sigma_{SB} T^4 m_p}{2\pi R_{WD}^2 k_B M_{WD}}
$$



Gaia G absolute magnitude



- 
- White dwarfs<br>• cooling curves & luminosity function

**Figure 6.** The white dwarf luminosity function from the deep proper motion survey (points with error bars, Munn et al. 2017) using a disk scale height range of 200–900 pc. The top panel shows the model fits assuming a population of 100% thin disk stars, whereas the bottom panel shows the fits using a composite population where the ratio of thick disk to thin disk white dwarfs is 35%. Dashed and dotted lines show the contribution from the thin disk and thick disk white dwarfs, respectively.



# Number density Number density

#### *t WD*  $\frac{max}{WD} \approx 10 \text{ Gyr}$

Can distinguish the age of the disk …

#### For white dwarfs in the disk of the Milky Way

- White dwarfs
	- cooling curves & luminosity function



… from the age of the stellar halo …

$$
t_{WD}^{max} = 12.5_{-3.5}^{+1.4} \text{ Gyr}
$$

For halo white dwarfs

**Figure 9.** Munn et al. (2017) luminosity function for the km s−1 halo white dwarf sample. Solid, dashed, and dotted lines show model luminosity functions for 12.5, 13.9, and 15.0 Gyr old halo samples, respectively. This luminosity function implies a halo age of  $12.5^{+1.4}_{-3.5}$  Gyr.

r-process elements created in SN explosions. Some are radioactive and decay with a half-life that is well-known from the laboratory. 232Thorium has a half-life of 14.05 Gyr, so makes a good cosmochronometer. 238Uranium has a half-life of 4.5 Gyr so is also useful, but is even harder to measure in stars. These are typically referenced to stable elements like Europium.



Fig. 7-30 Neutron-capture paths for the  $s$  process and the  $r$  process. The  $s$  process follows a path in the  $NZ$  plane along the line of beta stability. The neutron-rich progenitors to the stable r-process nuclei, which are here shown as small circles, are formed in a band in the neutron-rich area of the NZ plane, such as the shaded area shown here. This r-process path was calculated for the case  $T_9 = 1.0$  and  $\log n_n = 24$ . After the synthesizing event the nuclei in this band beta-decay to the stable r-process nuclei. The abundance peaks at  $A = 80$ , 130, and 195 are attributed to abundance peaks in the neutron-rich progenitors having  $N = 50$ , 82, and 126. Neutron capture flows upward from the lower left-hand corner along the shaded band until neutron-induced fission occurs near  $A = 270$ . [P. A. Seeger, W. A. Fowler, and D. D. Clayton, Astrophys. J. Suppl., 11:121 (1965). By permission of The University of Chicago Press. Copyright 1965 by The University of Chicago.]

- Radioactive chronometers
	- Uranium/Europium ratio
	- Thorium/Europium ratio

SN explosions provide neutron-rich environments that generate neutron-rich nuclei at high atomic number through the "r-process." These isotopes are highly unstable and quickly decay to the nuclear valley of stability.



#### • Radioactive chronometers

• Uranium/Europium ratio

r-process elements created in SN explosions. Some are radioactive and decay with a half-life that is well-known from the laboratory. 232Thorium has a half-life of 14.05 Gyr, so makes a good cosmochronometer. 238Uranium has a half-life of 4.5 Gyr so is also useful, but is even harder to measure in stars. These are typically referenced to stable elements like Europium.



A. Frebel & K.-L. Kratz

**Figure 2.** Spectral region around the U II line in HE  $1523-0901$  (filled dots) and CS 31082-001 (crosses; right panel only). Overplotted are synthetic spectra with different U abundances. The dotted line in the left panel corresponds to a scaled solar r-process U abundance present in the star if no U were decayed. Figure taken from (Frebel *et al.* 2007a).





- Radioactive chronometers
	- Thorium/Europium ratio

 $N_{\text{Th}} = N_{\text{Th}}(t_0) e^{-t/\tau_{\text{Th}}}$ 

Theoretically expect initial  $Th/Eu = 0.48$ . In CS 22892-052, Th/Eu =  $0.24$ . In the solar system,  $Th/Eu = 0.46$ .



$$
t_{*,\text{Th}} = 12.8 \pm 3 \text{ Gyr}
$$

Old stars are everywhere, not just in globular clusters. However, they are hard to identify in the field.

- Interstellar dust grains
	- Oxygen isotope ratios

Some dust grains have been captured from hyperbolic orbits: they have an origin external to the solar system. Such grains have isotopic ratios of oxygen indicating dredge-up during the red giant phase prior to the birth of the sun. Some grains with similar isotopic ratios were captured into meteorites at the dawn of the solar system. The required chain of events implies a minimum age:

$$
t_{MW} = t_{\odot} + t_{mix} + t_* + t_Z
$$

 $t_{\odot} = 4.55 \ \mathrm{Gyr} \quad$  age of solar system

 $t_{mix} \approx 0.1 \; \rm Gyr \quad$  mixing time in ISM

 $t_* = f(M_*, Z) \approx 5 \pm 1 \text{ Gyr}$  evolution of star that made the grain

 $t_Z \approx 4 \; \rm Gyr \; \;$  evolution of the stars that made the metals in the star than made the grain

$$
t_{MW} = 13.7 \pm 1.3
$$
 Gyr





FIG. 2. Comparison of oxide grain data for groups 1 and 3 (see Fig. 1 for symbol definitions) with predictions of  $O$ isotopic ratios following first dredge-up in red giant stars of initial mass  $(0.85-3)M_{\odot}$  and metallicity  $Z = 0.012-0.02$  [22]. For the sake of clarity, error bars on grain measurements are not shown. Each open circle corresponds to predictions for a distinct star. The dotted lines indicate interpolated values for masses and metallicities intermediate to those calculated. Oxide grains belonging to groups 1 and 3 have isotopic compositions consistent with these predictions, provided that they come from several different stars with distinct masses and initial compositions.



### **Cosmic Chronometers**

In principle, we can age date the stars in galaxies observed over a range of redshifts to empirically reconstruct the age-redshift relation.

LCDM age-redshift relation ( $H_0 = 70$ ,  $\Omega_m = 0.3$ ;  $t_0 = 13.5$  Gyr)





#### ivent Planck scale (*speculative*) GUT scale (*speculative*) Inflation (*speculative*) tandard Model forces emerge WIMPs decouple (*speculative*) quarks condense into baryons (*baryogenesis*) proton-antiproton annihilation ends neutrinos decouple electron-positron annihilation ends Big Bang Nucleosynthesis Matter-radiation equality Atoms form, CMB emerges Gas temperature decouples from radiation Dark Ages Cosmic dawn (first stars) Galaxies form Peak star formation Sun forms Multicellular life on earth You are now



### **Cosmic Timeline**

![](_page_10_Picture_0.jpeg)

Figure 17.4 This timeline summarizes conditions and transitions that marked the early eras of the universe.

Early U radiation dominated

*a* ∼  $t^{1/2}$  $T \sim a^{-1}$  $Tt^2 \sim \text{constant}$  $t \lesssim 10^5$  yr

Decoupling means to fall out of thermal equilibrium i.e., when it becomes impossible for the radiation field to spontaneously create particle-antiparticle pairs.

### **Early Universe**

![](_page_11_Figure_0.jpeg)

Time Event Planck scale (*speculative*)  $t \sim 10^{-43}$  s

 GUT scale (*speculative*)  $t \sim 10^{-38}$  s

Known physics breaks down at the Planck scale

 Inflation (*speculative*)  $t \sim 10^{-35}$  s

Period of exponential growth:  $a \sim e^{Ht}$ Must revert to radiation  $a \sim t^{1/2}$  after  $t \sim 10^{-24}$  s.

 Standard Model forces emerge  $t \sim 10^{-12}$  s

The four forces become distinct; one can begin to recognize "ordinary" particles that one might find in high energy particle accelerators:  $T \sim 10^{15} \text{ K} \sim 150 \text{ GeV}.$ Nb.: The LHC probes ~ 7 GeV, roughly equivalent to 10<sup>-16</sup> s

![](_page_11_Picture_10.jpeg)

GUT stands for Grand Unified Theory; this is the hypothetical scale at which the strong nuclear force becomes indistinguishable from the electroweak force.

$$
m_P = \sqrt{\frac{\hbar c}{G}} = 1.22 \times 10^{19} \text{ GeV c}^{-2} \approx 2 \times 10^{-5} \text{ g}
$$
  

$$
\ell_P = \sqrt{\frac{\hbar G}{c^3}} = 1.6 \times 10^{-35} \text{ m}
$$
  

$$
t_P = \frac{\ell_P}{c}
$$
; need a theory of quantum gravity.

![](_page_12_Figure_0.jpeg)

![](_page_13_Figure_1.jpeg)

Fig. 6.1 The thermal history of the standard Hot Big Bang. The radiation temperature decreases as  $T_r \propto R^{-1}$  except for abrupt jumps as different particle-antiparticle pairs annihilate at  $kT \approx mc^2$ . Various important epochs in the standard model are indicated. An approximate time scale is indicated along the top of the diagram. The neutrino and photon barriers are indicated. In the standard model, the Universe is optically thick to neutrinos and photons prior to these epochs.

Neutrinos drop out of equilibrium. They lose energy with expansion the same as the radiation field,  $T_\nu \thicksim a^{-1}$ , with an initial energy density  $\varepsilon_{\nu} \sim T_{\nu}^{4}$  fixed at this point.

This energy is deposited in the radiation field (which becomes the CMB). Only about one proton is left over for every billion proton-antiproton pairs (this is the matterantimatter asymmetry).

There are  $10^9 + 1$  protons for every  $10^9$  antiprotons

![](_page_13_Picture_210.jpeg)

 neutrinos decouple *t* ∼ 1 s

Being less massive than protons, electrons freeze out from positrons at this later time. The excess energy feeds the radiation background but not the neutrino background.

 electron-positron annihilation ends *t* ∼ 4 s

Note the brief pause in the decline of the temperature of the radiation field as first protonantiproton annihilation, then later electron-positron annihilation dump energy into radiation.

![](_page_14_Figure_0.jpeg)

M.S. Longair: The Physics of Background Radiation

Fig. 6.1 The thermal history of the standard Hot Big Bang. The radiation temperature decreases as  $T_r \propto R^{-1}$  except for abrupt jumps as different particle-antiparticle pairs annihilate at  $kT \approx mc^2$ . Various important epochs in the standard model are indicated. An approximate time scale is indicated along the top of the diagram. The neutrino and photon barriers are indicated. In the standard model, the Universe is optically thick to neutrinos and photons prior to these epochs.

432

 $\sigma_{\rm{eff}}$ 

![](_page_14_Figure_6.jpeg)

![](_page_14_Figure_7.jpeg)

![](_page_14_Picture_185.jpeg)

![](_page_14_Figure_8.jpeg)

Time Event Gas temperature decouples from radiation  $t \sim 5 \times 10^6$  yr

After recombination, the kinetic temperature of matter departs from that of the radiation field. They start out identical, and it takes a while for the matter to relax to fall as  $T_{mat} \sim a^{-2} \sim (1+z)^2$  after  $z \approx 200$ .

![](_page_15_Figure_0.jpeg)

 Dark Ages (no stars)  $t \sim 10^7$  yr

From  $z \approx 1000$  to  $z \approx 20$ , the universe is composed of neutral, primordial gas. Sources of light have yet to form, so this period is known as the Dark Ages.

 $t \sim 5 \times 10^8$  yr Cosmic dawn / re-ionization

Formation of the first stars (and maybe quasars?) These flood the universe with UV radiation that re-ionizes the universe. The gas in the intergalactic medium remains hot and highly ionized to this day.

![](_page_15_Figure_7.jpeg)

![](_page_15_Figure_8.jpeg)

![](_page_15_Figure_9.jpeg)

Time Event  $t \sim 4 \times 10^5$  yr Atoms form, CMB emerges  $t \sim 5 \times 10^6$  yr Gas temperature decouples from radiation

 Dark Ages *t* ∼ 5 × 10<sup>8</sup> yr Cosmic dawn / reionization  $t \sim 10^7$  yr

 Galaxies form  $t \sim 10^9$  yr

Baryons can only begin to gather together and form structures after decoupling from the radiation

Free fall time for Milky Way mass  $\sim 1 \; \rm Gyr$ Probably a messy process involving the merger of smaller protogalactic fragments that can collapse more quickly.

 $t \sim 4 \times 10^9$  yr Peak star formation

The star formation rate of the universe peaked at  $z\approx2;$ declines precipitously after *z* < 1

 $t \sim 9 \times 10^9$  yr Sun forms

A relative late comer, merely 4.5 Gyr old.

 $t \sim 13 \times 10^9$  yr multicellular life on earth

Singe-celled life appeared fairly early, but the Cambrian explosion didn't occur until 0.6 Gyr ago

![](_page_16_Picture_16.jpeg)

![](_page_16_Figure_0.jpeg)

### Empirical Pillars of Hot Big Bang

- Hubble expansion
- Big Bang Nucleosynthesis
- Relic Radiation (CMB)