

Mass components

Stars in galaxies

- Spiral & Irregular galaxies - Exponential disks

$$\Sigma(R) = \Sigma_0 e^{-R/R_0}$$

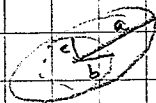
is a decent first approximation to the azimuthally averaged face-on surface brightness.

Can measure light profile, assume a mass-to-light ratio, and solve the Poisson equation to obtain the gravitational potential of the observed stars.

- Elliptical galaxies

spherical

$$a=b=c$$



3D ellipsoids seen in projection

can be oblate, prolate, or triaxial

$$a=b > c, \quad a > b = c, \quad a > b > c$$

(all can be a few of radius; position angle can twist)

Classical "r^{1/4}" or de Vaucouleurs profile

$$\Sigma(r) = \Sigma_e e^{-7.67 \left[\left(\frac{r}{R_e} \right)^{1/4} - 1 \right]}$$

R_e = effective radius (contains $\frac{1}{2}$ light)

$$\Sigma_e = \Sigma(R_e)$$

Sersic profile: generalized hybrid

$$\Sigma(r) = \Sigma_e 10^{-b_n \left[\left(\frac{r}{R_e} \right)^n - 1 \right]}$$

adds shape parameter n

$n=1 \rightarrow$ exponential disk

$n=4 \rightarrow$ de Vaucouleurs profile

Galaxies with both bulge & disk can be fit as separate components (e.g. exponential disk + "classical" r^{1/4} bulge) or as a single Sersic profile (with n managing the transition from bulge to disk).

The ISM

The Interstellar Medium contains gas in all phases - molecular, atomic, and ionized (plasma) as well as dust.

In the Milky Way, all of these are largely confined to the plane of the disk, though there probably exists a tenuous quasi-spherical corona of hot, ionized gas at radii far beyond the edge of the stellar disk, ~~extending~~ extending to ~ 300 kpc (very roughly).

This corona contains relatively little mass, though it might integrate up to $\sim 10^{10} M_{\odot}$ all the way out.

IN the disk of the Milky Way, in order of mass:

Atomic gas	HI	$\sim 60-70\%$	of ISM mass
Molecular gas	H ₂	$\sim 30\%$	" " "
Ionized gas			
Ionized gas	HII	$< 1\%$	" " "
Dust		$\ll 1\%$	

Total gas mass in MW $\sim 10^{10} M_{\odot}$
dust mass maybe $\sim 10^6 M_{\odot}$

HI extends to ~ 20 kpc

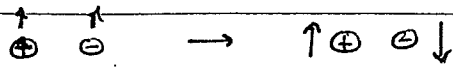
(The stellar $R_d \approx 2.2$ kpc, so almost all stellar mass is interior to 10 kpc)

$M_{\text{gas}} \approx 5 \text{ or } 6 \times 10^{10} M_{\odot}$

so the Milky Way has a gas fraction $f_g \approx 20\%$ which is typical for bright spirals.

HI Atomic gas

detected by spin-flip transition of H



transition frequency very low energy

$$\nu_{10} = \frac{8}{3} g_I \left(\frac{m_e}{m_p} \right) \alpha^2 (cR_H) = 1420.4 \text{ MHz}$$

nuclear g-factor
 $g_I = 5.6$ for proton

fine structure constant
 $\alpha = \frac{1}{137}$

the 21cm line

Rydberg frequency

Emission coefficient

$$A_{ul} = \frac{64 \pi^4 \nu_{ul}^3}{3 h c^3} |\mu_{10}^*|^2$$

Bohr magneton
(dipole moment of e^-)

numerically

$$A_{10} = 2.85 \times 10^{-15} \text{ s}^{-1}$$

$$|\mu_{10}^*| = \frac{\hbar}{2} \frac{e}{m_e c}$$

radiative half-life $\tau_{1/2} = \frac{1}{A_{10}} \approx 11 \times 10^6 \text{ yr}$
 $(3.5 \times 10^{14} \text{ s} = 11 \text{ Myr}) = \frac{1}{\tau_{1/2}}$

Very low critical density
($\ll 1 \text{ cm}^{-3}$)

so almost always in LTE through
collisional excitation

Can also define "spin temperature"

$$\frac{N_1}{N_0} = \frac{g_1}{g_0} e^{-\frac{h\nu_{10}}{kT}}$$

$g_1 = 3$ upper states

$g_0 = 1$ ground state

Note $\frac{h\nu}{k} \ll 1 \text{ K}$ very small, so $e^{-\frac{h\nu}{kT}} = 1$ if $T_s \approx T$

so $\frac{N_1}{N_0} = \frac{3}{1}$ & $N_H = N_0 + N_1 = 4N_0$ good for counting atoms

anywhere near equilibrium

counting γ rays photons \rightarrow counting HI atoms

$$M_{\text{HI}} = \frac{16\pi M_{\text{H}} D^2}{3A_{\text{ul}} hc} \int F_{\nu} dv$$

flux integral in $\text{Jy} \cdot \text{km/s}$

Numerically

$$1 \text{ Jansky} = 10^{-26} \text{ W m}^{-2} \text{ Hz}^{-1}$$

$$M_{\text{HI}} = (2.34 \times 10^5 M_{\odot}) D^2 \int F_{\nu} dv$$

$$M_{\text{atomic gas}} = \frac{1}{X} M_{\text{HI}}$$

Can measure HI masses

Hydrogen mass fraction

accurately to a few % IF care is taken.

D in Mpc (luminosity distance)

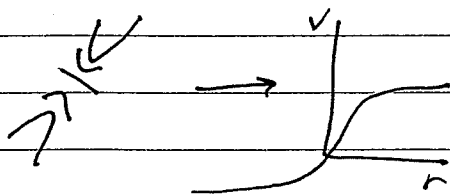
in cosmology

because it's a line

HI emission ALSO gives velocity field

from Doppler effect - important tracer of the gravitational potential

Fit velocity field with tilted ring model to obtain rotation curve



Molecular ISM

cold $\sim 30\text{K}$

"dense" $\sim 100\text{ cm}^{-3}$

very clumpy; need to be to self-shield against interstellar radiation field (with UV can photodissociate many molecules)

most mass in GMCs: Giant Molecular Clouds $\sim 10^6 M_{\odot}$

... this is where stars are born

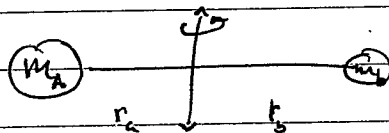
Diatomic molecules boring - or at least hard to excite

$\hookrightarrow \text{H}_2, \text{O}_2, \text{N}_2$ because symmetry prevent them from having a dipole moment.

Hence no radiation if they rotate (vibrational states much higher energy)

Polar molecules (like CO) have permanent dipole moment thanks to asymmetry

leading to rich rotational spectrum (often at mm wavelengths & cm)



Angular momentum

$$L = \left(\frac{m_a m_b}{m_a + m_b} \right) (r_a + r_b) \omega \quad \text{is quantized}$$

rotational energy

$$E = \frac{J(J+1) \hbar^2}{2I}$$

moment of inertia

$$I = m_a r_a^2 + m_b r_b^2$$

need I to be large to generate detectable emission in cold molecular gas

quantum mechanics requires $\Delta J = \pm 1$ so

$$v = \frac{hJ}{4\pi m_e r_e^2}$$

$$m_e = \frac{m_a m_b}{m_a + m_b}$$

$$r_e = r_a + r_b$$

Rotational transitions lead to ladder spectrum

gets more complicated for multi-atom molecules

Can play same tricks as with HI:

$$M_{H_2} = 1.1 \times 10^4 D^2 F_{co}$$

$$X_{co} = 2.8 \times 10^{20} \text{ cm}^{-2} (\text{K km/s})$$