MOND: an Empiricist's Perspective

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A Field of Two Attitudes

\bullet ΛCDM is obviously correct

- CMB, LSS, etc. "cosmology solved!" (Turner 1998)
- considering alternatives is a waste of time

• Λ CDM is obviously wrong

- cusp/core, missing satellites, no direct detection of CDM, DE just another fudge
- much more reasonable to consider alternatives if ΛCDM is falsified
 - ... if that is even possible



Heck, I reckon you wouldn't even be human beings if you didn't have some pretty strong personal feelings about cosmology.

WIMP detection experiments



- WIMPs have been excluded at > 95% c.l. Repeatedly.
- Usual fix has been to make more massive the particle through which WIMPs interact, thus lowering their interaction cross-section.
 - Use to be the usual weak force carriers.
 - Now down to WIMPs that exchange Higgs particles.

One can imagine DM particles that don't interact with anything but gravity. Might be true, but isn't falsifiable. Tuning the interaction cross-section ever-further down is *the express elevator to hell*.



3 Laws of Galactic Rotation

- Rotation curves tend towards asymptotic flatness
- Baryonic mass scales as the fourth power of rotation velocity (Baryonic Tully-Fisher)
- Gravitational force correlates with baryonic surface density

Just the facts, mam. Just the facts.





2. Baryonic Tully-Fisher relation: $M_b = 47 V^4$



The data specify a particular acceleration scale: $a = \frac{V_f^4}{GM_b}$



histogram: data line: distribution expected from observational uncertainties.

The data are consistent with zero intrinsic scatter.





UGC 128

The BTFR is just the zeroth moment, as it were total baryonic mass vs. characteristic circular velocity. There is more information in the distribution of mass.



Same (M,V) but very different size and surface density which is strange, since $V^2 = \frac{GM}{R}$



For disk galaxies

Newton says:
$$\frac{V^2}{R} = \frac{GM}{R^2} = G(\Sigma_b + \Sigma_{DM})$$

surface density $\Sigma = \frac{M}{R^2}$

So we infer that $\Sigma_{DM} \gg \Sigma_b$

(i.e., that all disks are dark matter dominated) in order to explain the lack of TF residuals with luminous size R_p

There is still more information in the mass distribution...

Just looking at the peak radius





central baryonic surface density

But wait - before we decided $\Sigma_b \ll \Sigma_{DM}$

Now it has to matter. Is this a contradiction?

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Lack of TF residuals says baryon distribution does not matter.

Correlation of dynamical force with observed surface density says the baryon distribution does matter.

$$\frac{V^2}{R} = G\left(\Sigma_b + \Sigma_{DM}\right)$$





with the acceleration (gravitational force per unit mass)

Renzo's Rule: (2004 IAU; 1995 private communication) "When you see a feature in the light, you see a corresponding feature in the rotation curve."





Renzo's Rule:

"When you see a feature in the light, you see a corresponding feature in the rotation curve."

The distribution of mass is coupled to the distribution of light.

Quantify by defining the Mass Discrepancy:

$$\mathcal{D} = \frac{V^2}{V_b^2} = \frac{V^2}{\Upsilon_\star v_\star^2 + V_g^2}$$

The Mass Discrepancy correlates with acceleration and baryonic surface density



McGaugh (2004)

Can see the effect directly in the data with no assumption on M*/L







Can fit a fcn to the data $D(\Sigma/\Sigma_{\dagger})$ or $D(a/a_0)$



and use it to map between V_b and V_{tot} .





3 Laws of Galactic Rotation

- I. Rotation curves tend towards asymptotic flatness
- Baryonic mass scales as the fourth power of rotation velocity (Baryonic Tully-Fisher)
- Gravitational force correlates with baryonic surface density

No theory so far - just data.

Can always be interpreted in terms of dark matter (with sufficient fine-tuning).

Might stem more naturally from a universal force law.

MOND

 $a \gg a_0 \qquad a \to g_N \qquad a_0 \approx 10^{-10} \text{ m s}^{-2} \sim cH_0 \sim c\Lambda^{1/2}$ $a \ll a_0 \qquad a \to \sqrt{g_N a_o} \qquad \longrightarrow a_0 GM = V_f^4$ Milgrom (1983)

modified inertia (F=ma) OR modified gravity

Bekenstein & Milgrom (1984): $\nabla \cdot \left[\mu(|\nabla \phi|/a_0) \nabla \phi \right] = 4\pi G \rho$

$$\mu \to 1 \qquad a \gg a_0$$
$$\mu \left(\frac{a}{a_0}\right) = \frac{g_N}{a} \qquad \mu \to \frac{a}{a_0} \qquad a \ll a_0$$

Relativistic extensions: TeVeS, bimetric gravity, ???



Always emotions MOND is.

ApJ, 270, 381

Milgrom 1983

No. 2, 1983

MODIFICATION OF NEWTONIAN DYNAMICS

A major step in understanding ellipticals can be made if we can identify them, at least approximately, with idealized structures such as the FRCL spheres discussed above. I have also studied isotropic and nonisotropic isothermal spheres, in the modified dynamics, as such possible structures. I found that they have properties which feed near resemble those of empticale and galactic ourge. I i so it there in Mit rot. (181)

VIII. PREDICTIONS

The main predictions conce lows.

1. Velocity curves calculated with the modified dynamics on the basis of the observed mass in galaxies should agree with the observed curves. Elliptical and S0 galaxies may be the best for this purpose since (a)practically no uncertainty due to obscuration is involved and (b) there is not much uncertainty due to the possible presence of molecular hydrogen.

2. The relation between the asymptotic velocity (V_{∞}) and the mass of the galaxy (M) $(V_{\infty}^4 = MGa_0)$ is an absolute one.

3. Analysis of the z-dynamics in disk galaxies using the modified dynamics should yield surface densities which agree with the observed ones. Accordingly, the same analysis using the conventional dynamics should yield a discrepancy which increases with radius in a predictable manner.

4. Effects of the modified dynamics are predicted to be particularly strong in dwarf elliptical galaxies (for review of properties see. e.g., Hodge 1971 and Zinn 1980). For example, those dwarfs believed to be bound to our Galaxy would have internal accelerations typically of order $a_{in} \sim a_0/30$. Their (modified) acceleration, g, in the field of the Galaxy is larger than the internal ones but still much smaller than $a_0, g \approx (8)$ kpc/d) a_0 , based on a value of $V_{\infty} = 220 \text{ km s}^{-1}$ for the Galaxy, and where d is the distance from the dwarf galaxy to the center of the Milky Way (d - 70-220)kpc). Whichever way the external acceleration turns out to affect the internal dynamics (see the discussion at the end of § II, the section on small groups in Paper III, and Paper I), we predict that when velocity dispersion data is available for the dwarfs, a large mass discrepancy will result when the conventional dynamics is used to determine the masses. The dynamically determined mass is predicted to be larger by a factor of order 10 or more than that which can be accounted for by stars. In case the internal dynamics is determined by the external acceleration, we predict this factor to increase with d

and be of order (d/8 kpc) (as long as $a_{in} \ll g$, $h_{50} = 1$). Prediction 1 is a very general one. It is worthwhile listing some of its consequences as separate predictions. numbered 5-7 below (note that, in fact, even prediction 2 is already contained in prediction 1). 5. Measuring local M/L values in disk galaxies (assuming conventional dynamics) should give the following results: In regions of the galaxy where $V^2/r \gg a_{\eta}$ the local M/L values should show no indication of hidden mass. At a certain transition radius, local M/L should start to increase rapidly. The transition radius

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the all operations $P(A \neq a)$ is solve as the following $P(A \neq a)$ is solve as the following $P(A \neq a)$ we have not require an absolute calibration of M/L as we are concerned only with variations of this quantity; (b) Effects of the modified dynamics manifest themselves more clearly in local cases

ior in the lisk only while the spheroid can be neglected. This makes the determination of mass from velocity more certain.

6. Disk galaxies with low surface brightness provide particularly strong tests (a study of a sample of such galaxies is described by Strom 1982 and by Romanishin et al. 1982). As low surface brightness means small accelerations, the effects of the modification should be more noticeable in such galaxies. We predict, for example, that the proportionality factor in the $M \propto V_{\perp}^4$ relation for these galaxies is the same as for the high surface density galaxies. In contrast, if one wants to obtain a correlation $M \propto V_{\infty}^4$ in the conventional dynamics (with additional assumptions), one is led to the relation $M \propto$ $\Sigma^{-1}V_{\infty}^{4}$ (see, for example, Aaronson, Huchra, and Mould 1979), where Σ is the average surface brightness. This implies that low surface density galaxies, of a given velocity, have a mass higher than predicted by the M-Vrelation derived for normal surface density galaxies.

We also predict that the lower the average surface density of a galaxy is, the smaller is the transition radius, defined in prediction 5, in units of the galaxy's scale length. In fact, if the average surface density is very small we may have a galaxy in which $V^2/r < a_0$ everywhere, and analysis with conventional dynamics should yield local M/L values starting to increase from very small radii.

7. As the study of model rotation curves shows, we predict a correlation between the value of the average surface density (or brightness) of a galaxy and the steepness with which the rotational velocity rises to its asymptotic value (as measured, for example, by the radius at which $V = V_{\infty}/2$ in units of the scale length of the disk). Small surface densities imply slow rise of V.

IX. DISCUSSION

The main results of this paper can be summarized by the statement that the modified dynamics eliminates the need to assume hidden mass in galaxies. The effects in galaxies which I have considered, and which are commonly attributed to such hidden mass, are readily explained by the modification. More specifically:

MOND predictions

• The Tully-Fisher Relation

surfigee brightness

- Normalization = 1/(a,G)
 Fundamentally a relation between Disk
 - Fundamentally a relation between Disk Mass and V_{flat}
 - No Dependence on Surface Brightness
- Dependence of conventional M/L on radius and surface brightness
- Rotation Curve Shapes
- Surface Density ~ Surface Brightness
- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios



• The Tully-Fisher Relation



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- $M_* > M_g (MOND fits)$ McGaugh (2005)
- $\begin{tabular}{ll} $$ $M_{*} > M_{g}$ (H-band popsynth)$ \\ $$ Sakai$ (2000); Gurovich et al. (2010)$ \\ \end{tabular}$

star dominated

gas dominated

- $M_* < M_g \sin(i_{opt}) < 1.12 \sin(i_{HI})$ Begum et al. (2008)
- $M_* < M_g$ Stark et al. (2009)
- $M_* < M_g$ Trachternach et al. (2008)

Position on BTFR independent of stellar M*/L for M* < Mg $\,$



• The Tully-Fisher Relation



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• The Tully-Fisher Relation



Dependence of conventional M/L on radius and surface brightness



- Surface Density ~ Surface Brightness
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Dependence of conventional M/L on radius and surface brightness

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Sanders & McGaugh 2002, ARA&A, 40, 263



Sanders & McGaugh 2002, *ARA&A*, **40**, 263



Sanders & McGaugh 2002, *ARA&A*, **40**, 263



Residuals for 74 galaxies

This procedure is generally successful (including the bumps and wiggles) given M*/L as a fit parameter

• The Tully-Fisher Relation

Dependence of conventional M/L on radius and surface brightness

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Line: stellar population model (mean expectation)

• The Tully-Fisher Relation

Dependence of conventional M/L on radius and surface brightness

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Observational Test	Successful	Promising	Unclear	Problematic
Rotating Systems				
solar system			Х	
galaxy rotation curve shapes	х			
surface brightness $\propto \Sigma \propto a^2$	X			
galaxy rotation curve fits	Х			
fitted M _* /L	Х			
Tully-Fisher Relation				
baryon based	X			
slope	X			
normalization	X			
no size nor Σ dependence	X			
no intrinsic scatter	х			
Galaxy Disk Stability				
maximum surface density	Х			
spiral structure in LSBGs	X			
thin & bulgeless disks		х		
Interacting Galaxies				
tidal tail morphology		х		
dynamical friction			Х	
tidal dwarfs	Х			
Spheroidal Systems				
star clusters			Х	
ultrafaint dwarfs			Х	
dwarf Spheroidals	X			
ellipticals	X			
Faber-Jackson relation	X			
Clusters of Galaxies				
dynamical mass				X
velocity (bulk & collisional)		х		
Gravitational Lensing				
strong lensing	X			
weak lensing				Х
Cosmology				
expansion history			х	
geometry			х	
big bang nucleosynthesis	X			
Structure Formation				
galaxy power spectrum			Х	
empty voids		х		
early structure		Х		
Background Radiation				
first:second acoustic peak	X			
second:third acoustic peak				X
detailed fit				X
early re-ionization	X			

Table 1: Observational tests of MOND.

1E 0657-56 - "bullet" cluster (Clowe et al. 2006)

residual mass discrepancy in clusters is real... the bullet cluster is a special case of a more general problem. Data for groups & cluster offset from MOND prediction, but slope pretty good over many decades in baryonic mass.

$$a_0 = 1.2 \times 10^{-10} \text{ m s}^{-2} \approx \frac{cH_0}{2\pi} \approx c\Lambda^{1/2}$$

 $\Sigma_{\dagger} = 860 \ \mathrm{M}_{\odot} \ \mathrm{pc}^{-2}$

The missing baryon problem

MOND MOND z = 105 Ż

ΛCDM

5

 \boldsymbol{z}

3

 \boldsymbol{z}

ACDM

Ż

Expect big clusters at high redshift

Other MOND tests

- Disk Stability
 Freeman limit in surface brightness distribution thin disks
 velocity dispersions
 LSB disks not over-stabilized
- Dwarf Spheroidals ?
- Giant Ellipticals
- X• Clusters of Galaxies
- **?** Structure Formation
 - Microwave background
 1st:2nd peak amplitude; BBN
 early reionization
 enhanced ISW effect
 3rd peak

X No Metric [±]/₅ [∞]
X Don't know expansion history[∞]

Logical possibilities

- ACDM is fine; puzzling observations will be explained by complicated feedback processes.
- MOND gets predictions right because there is something to it --- dark matter doesn't exist.
- We have no clue what is going on.

