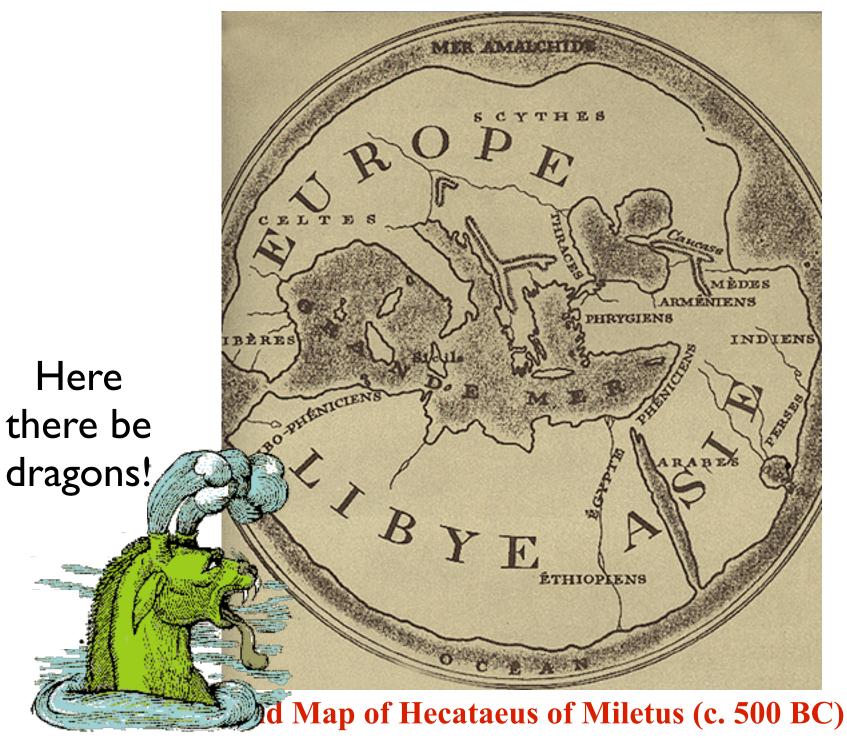
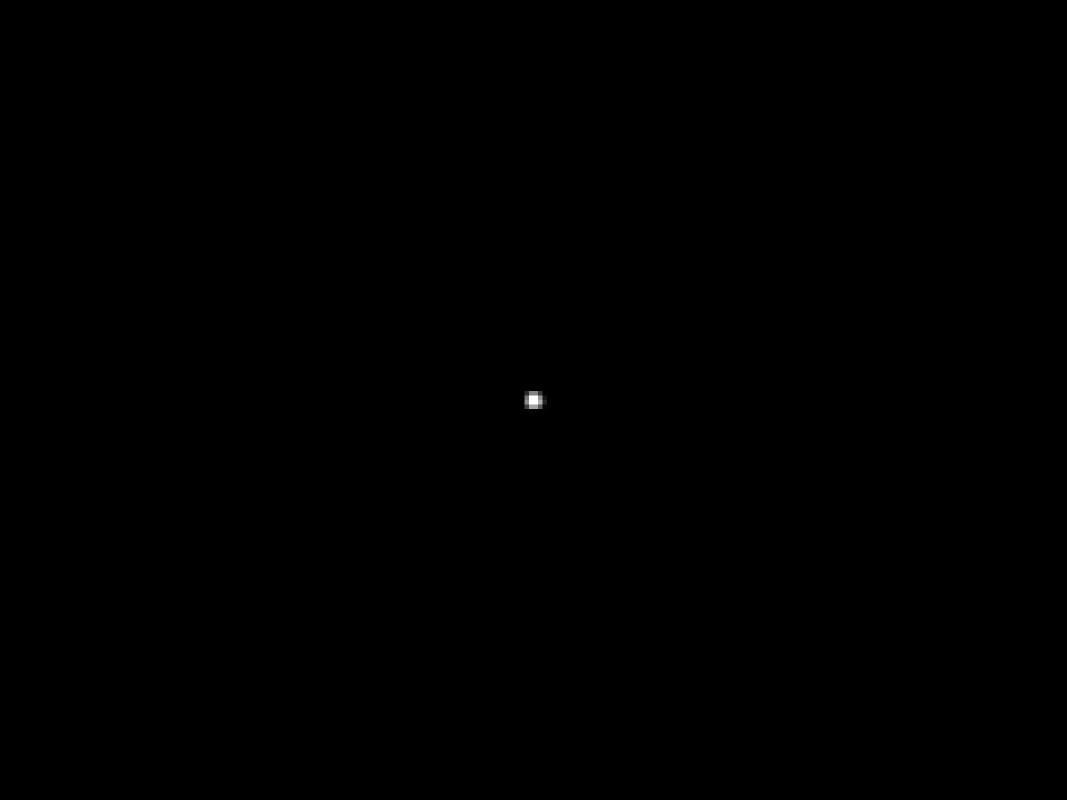
Mass discrepancies in the Universe: missing mass or new physics?

Stacy McGaugh University of Maryland

Ancient Cosmology: A Flat Earth







The modern idea for the origin of this expanding universe is called the



Three Pillars of the

BIG BANG

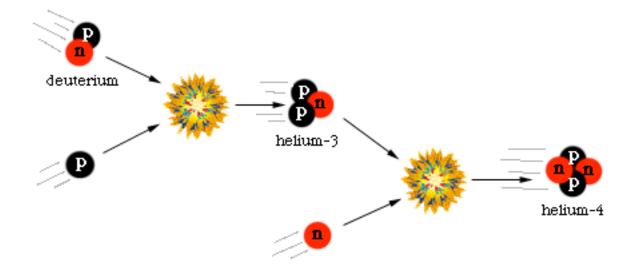
- Expanding Universe
- Primordial Nucleosynthesis
 Origin of the Light Elements (H, He, Li)
- Relic Radiation Field
 The Cosmic Microwave Background

Primordial Nucleosynthesis:



Gamow

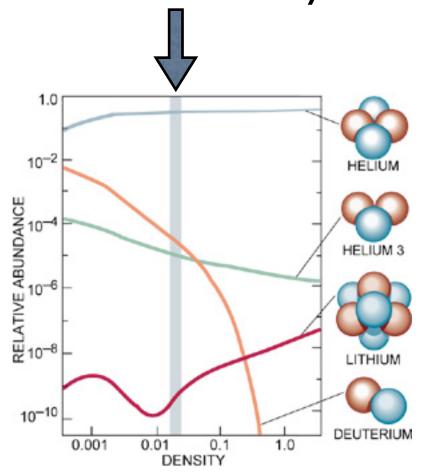
When the universe is just a few minutes old, the Temperature and Density are just right for it to be one Big Nuclear Furnace:



The light elements
Hydrogen, Helium, and Lithium
are made at this time.

The amount of each element to emerge from the big bang depends on the amount of raw material (protons + neutrons = "baryons")

Measured abundances point to a modest amount of "normal" baryonic matter



 $\Omega_b \approx 0.04$

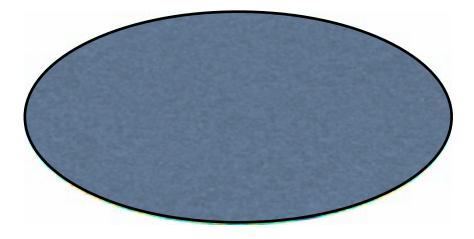
Relic Radiation Field:

The residual heat of the Big Bang should leave an echo - a relic glow of the cosmic fireball.

This was discovered in 1963; now called the Cosmic Microwave Background

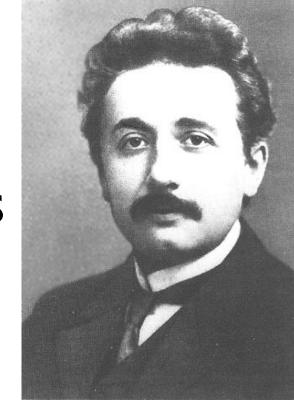


Wilson Penzias
Nobel Prize



Baby picture of the planterse it in 1980,000 years of above absolute zero.

- The Expanding Universe
- Primordial Nucleosynthesis
- a Relic Radiation Field



all fit beautifully with the idea of a Hot Big Bang which makes sense in terms of Einstein's theory of General Relativity. Just need to know two numbers...

$$H_0 \& \Omega_m$$

and we know everything.

But there is more:

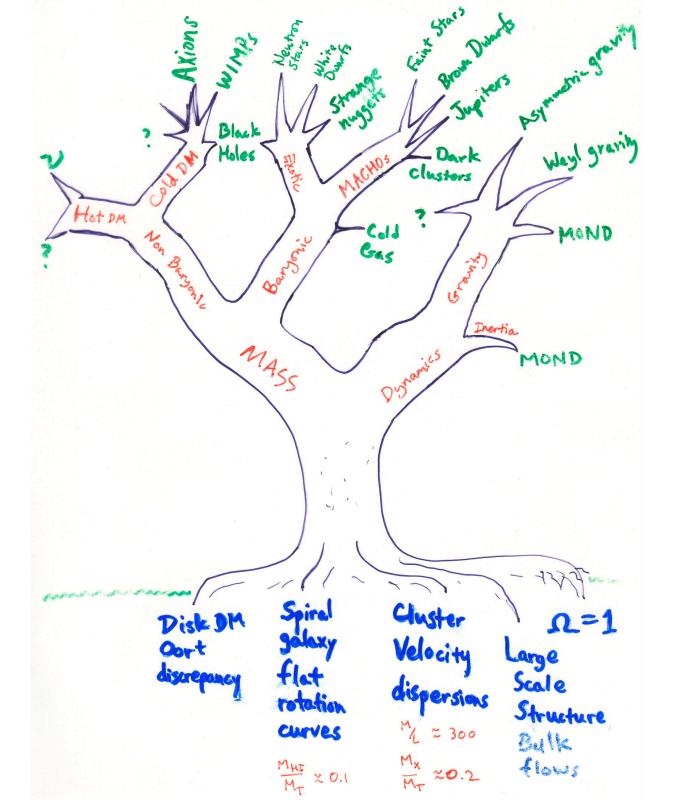
- Dark Matter
- Dark Energy

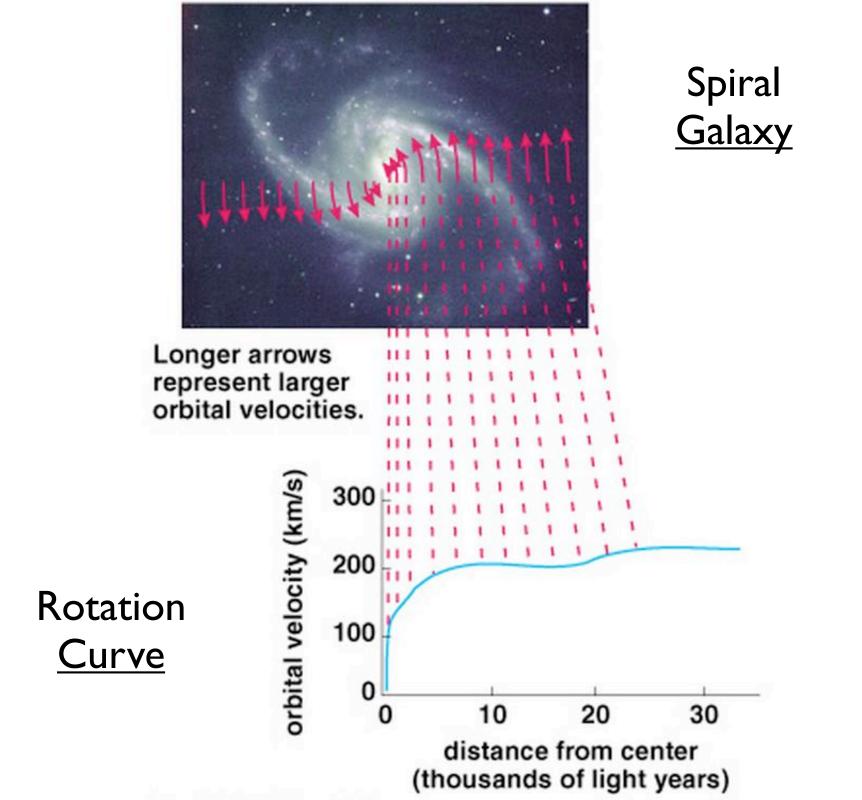
How many tooth fairies should a theory be allowed?

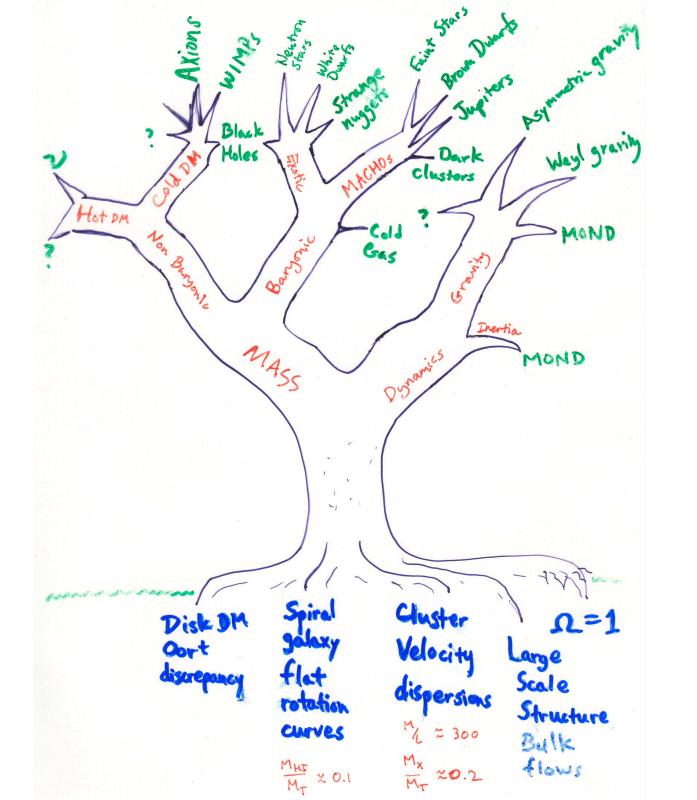


so maybe we need a few more numbers:

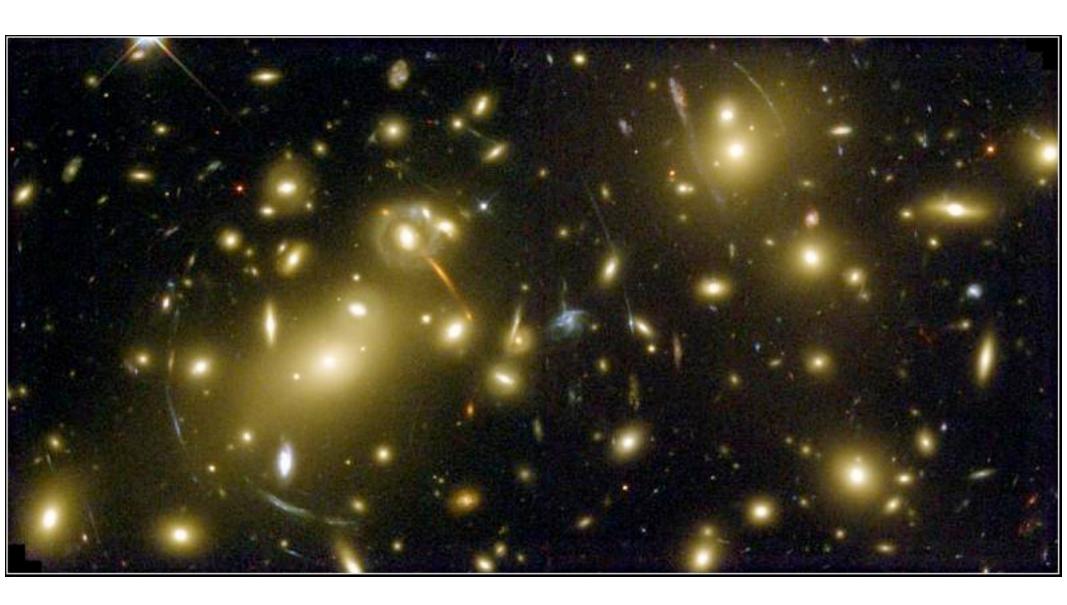
$$\Omega_b \Omega_\Lambda \sigma_8 n$$

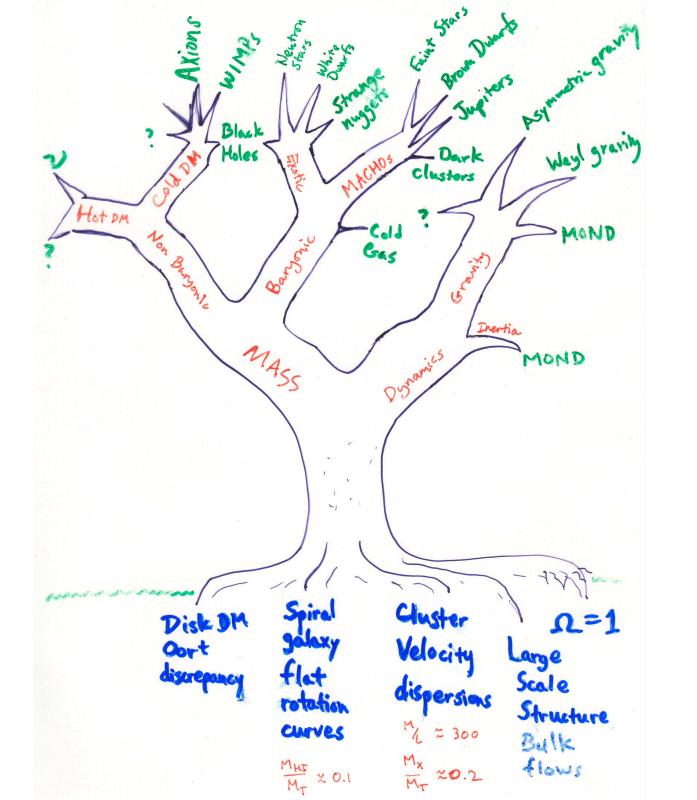




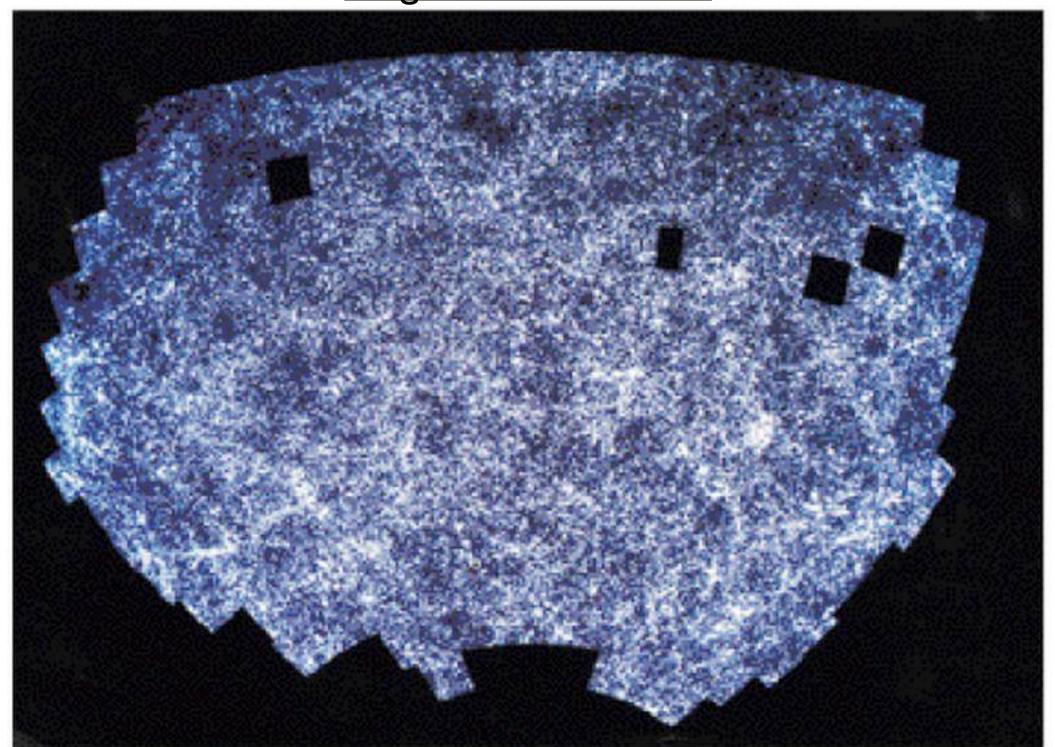


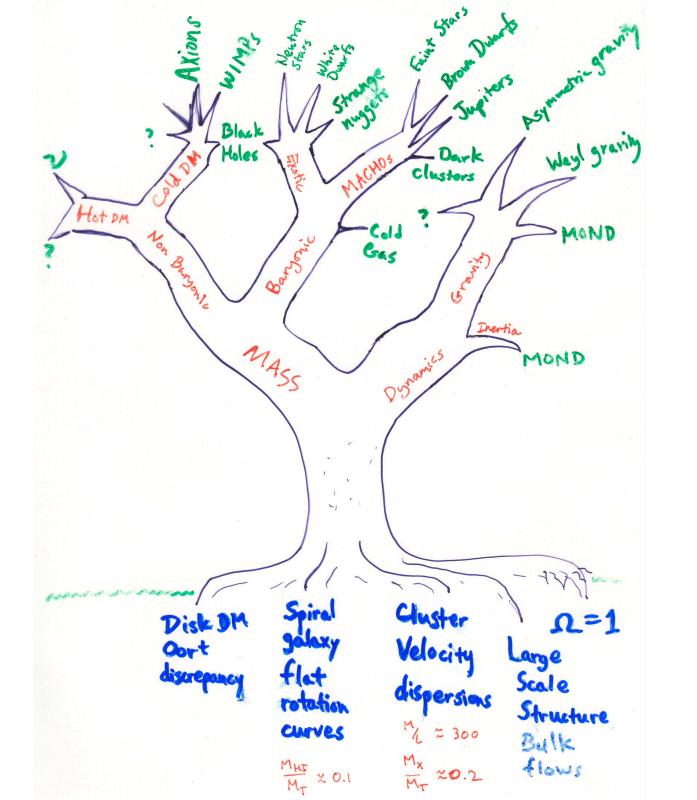
Galaxy Cluster

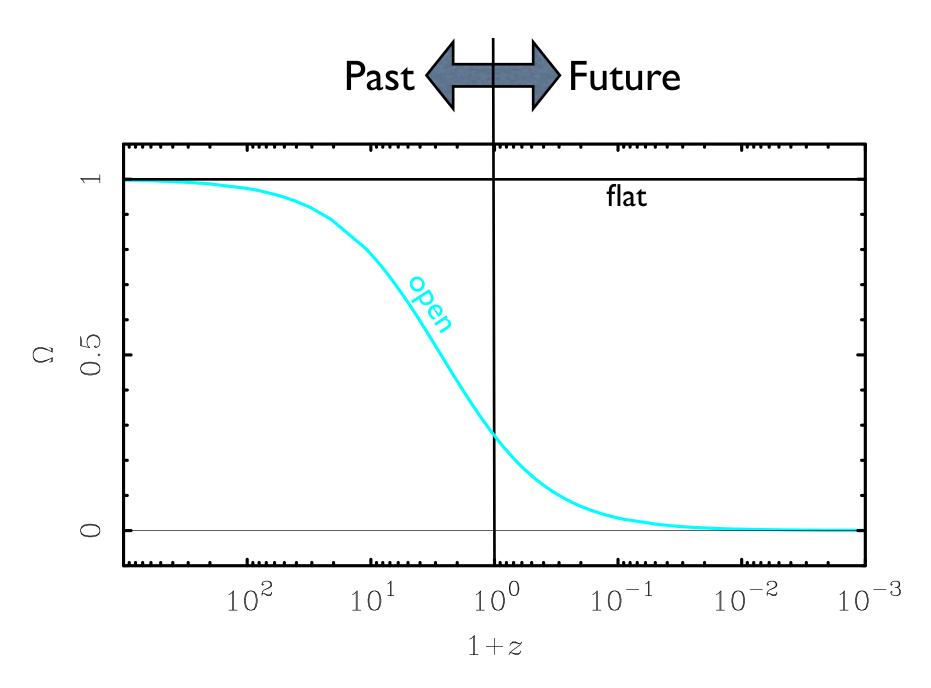




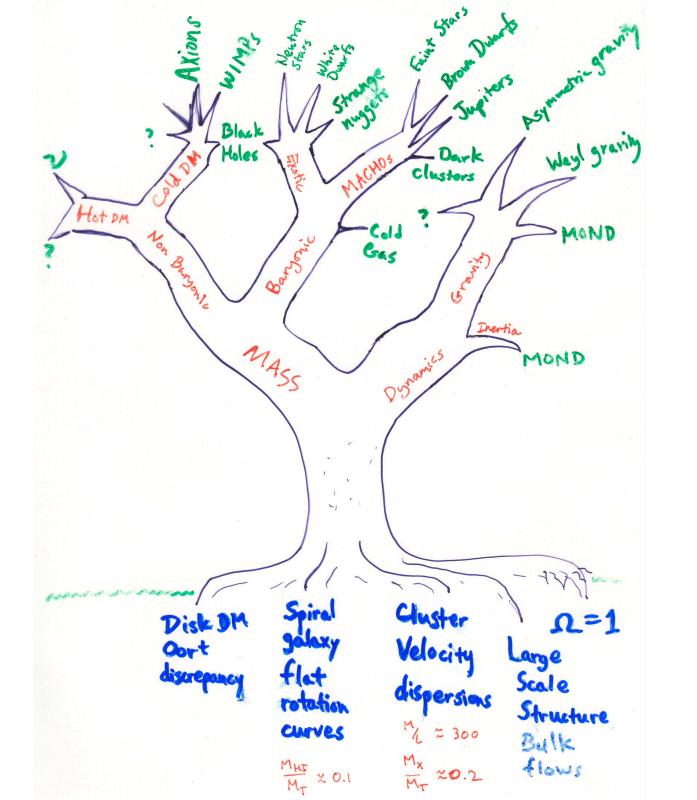
Large Scale Structure







Coincidence Problem solved iff $\,\Omega_m=1\,$



What is the Dark Matter?

Baryonic Dark Matter

Norman things:

very i int stars, brown dwarfs other hard-to-see objects (planets, gas)

Hot Dark Matter

neutrinos - got mass, but not enough

Cold Dark Matter

Some new fundamental particle doesn't interact with light, so quite invisible.

Two big motivations:

- I) total mass outweighs normal mass from BBN
- 2) needed to grow cosmic structure

(I)

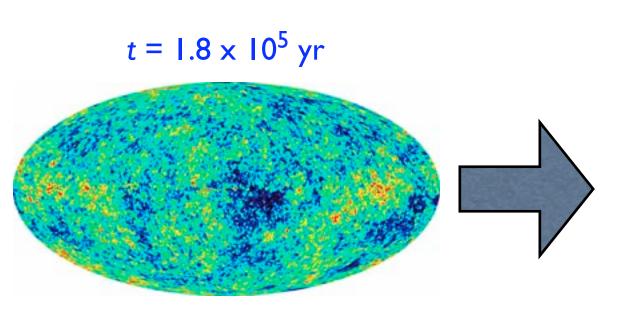
Normal baryonic mass $\Omega_b \approx 0.04$ from Primordial Nucleosynthesis

Total mass density $\Omega_m \approx 0.25$ from gravity

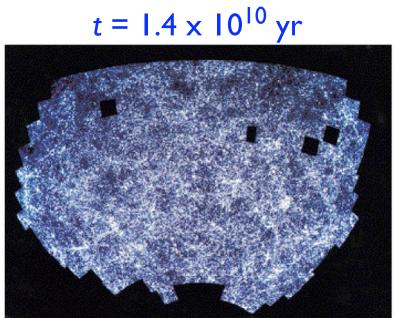
gravitating mass > normal mass

Most of the mass needs to be in some brand new form!

(2) There isn't enough time to form the observed cosmic structures from the smooth initial conditions unless there is a component of mass independent of photons.



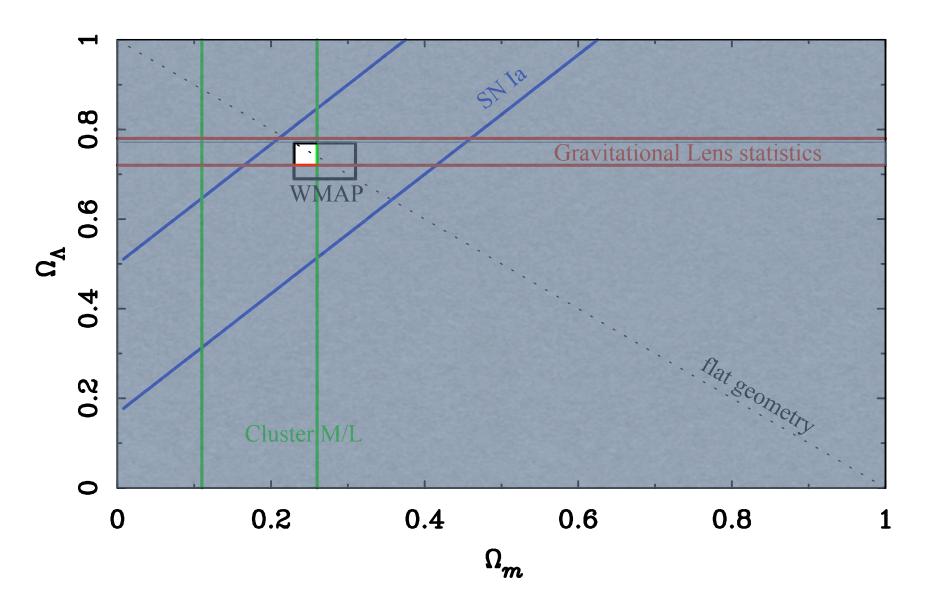
very smooth: $\delta \rho / \rho \sim 10^{-5}$



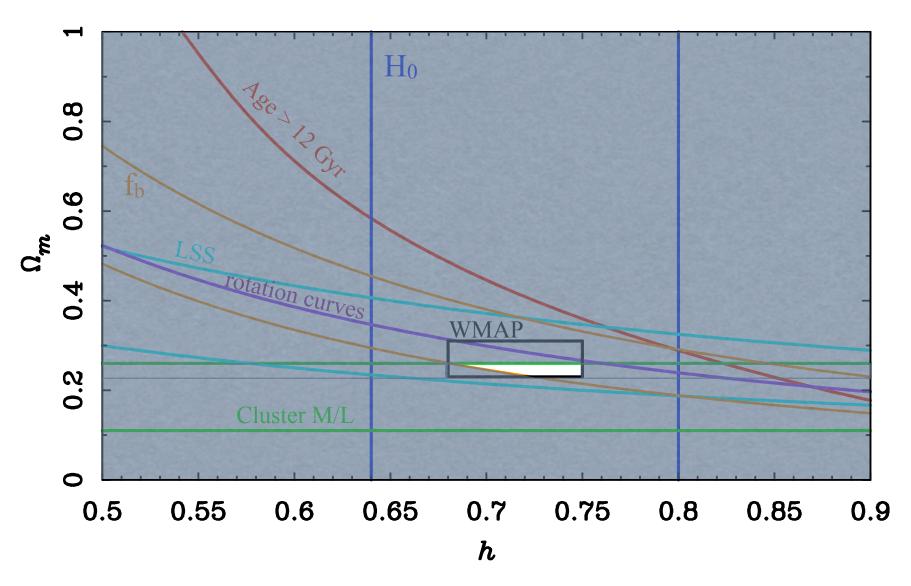
very lumpy: $\delta \rho / \rho \sim 1$

 $\delta \rho / \rho \propto t^{2/3}$

Cosmic parameters now constrained by multiple independent data sets

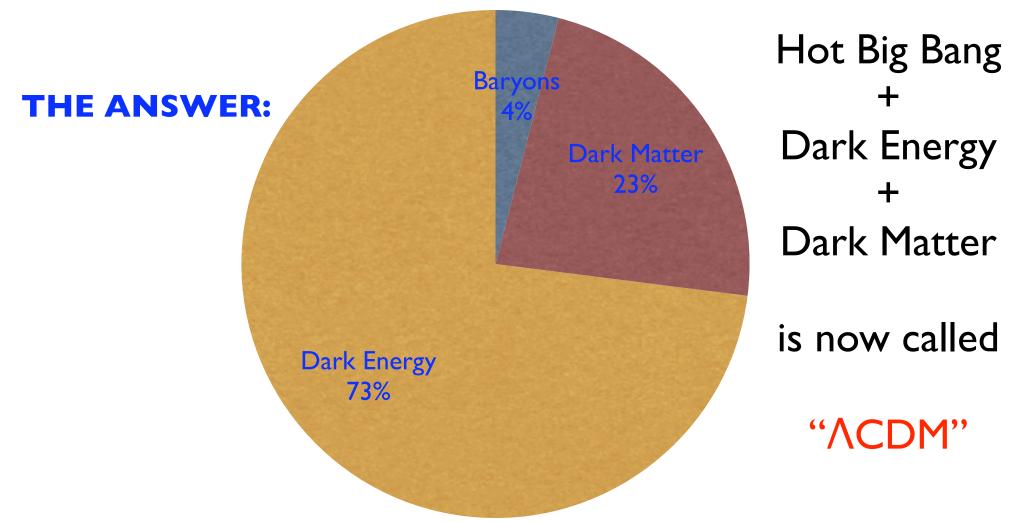


Cosmic parameters now constrained by multiple independent data sets



Is open window the right answer? Or merely the least improbable?

Cosmological Mass-Energy Budget



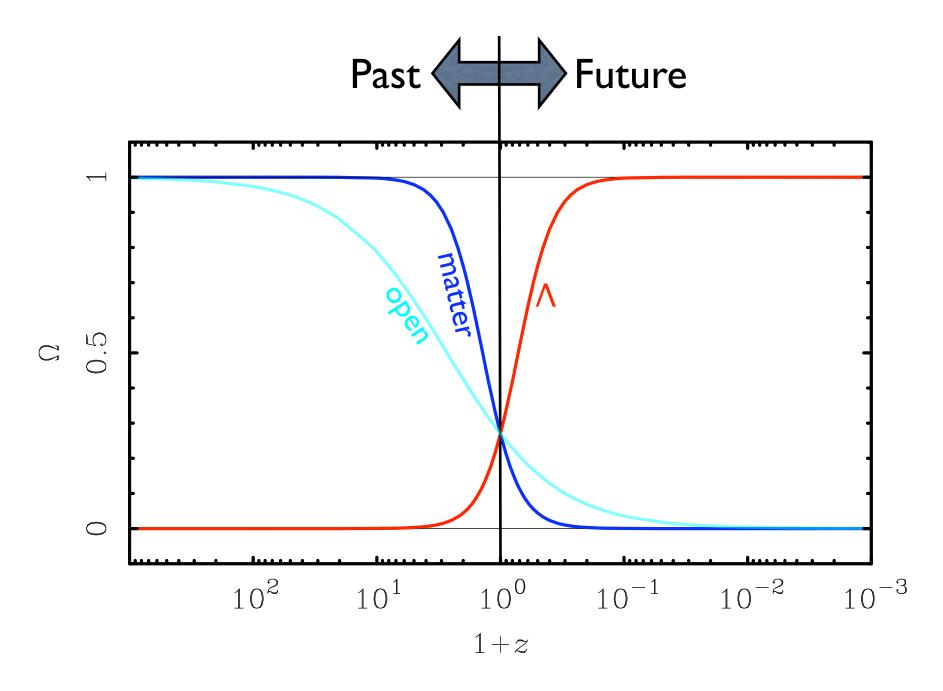
"The molye imagnest interpredale and the possibility of their ever discoveries is exceedingly. Kennote. (2004) Our future discoveries must be looked for in the sixth place of decimals."

- A. Michelson (1894)



You got a problem with Λ CDM?



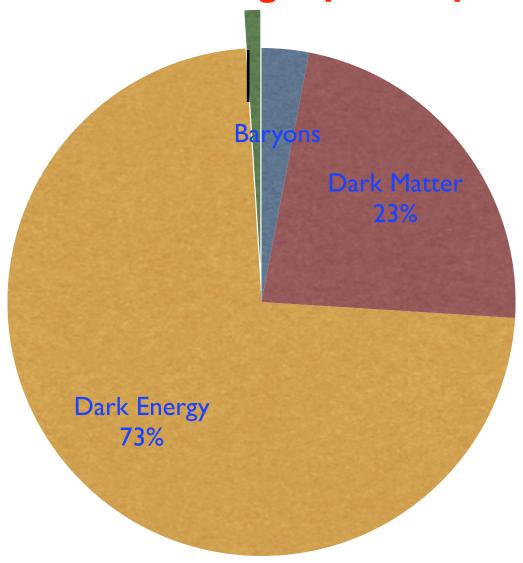


Coincidence Problem has gotten worse

Coincidence Problem Anthropomorphic Arguments



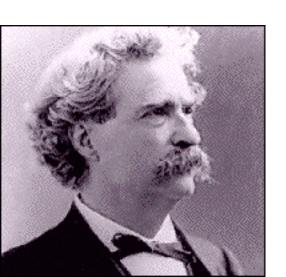
Known Baryons
We have direct knowledge of < 1% of this stuff.



Do these the Dark components really exist?

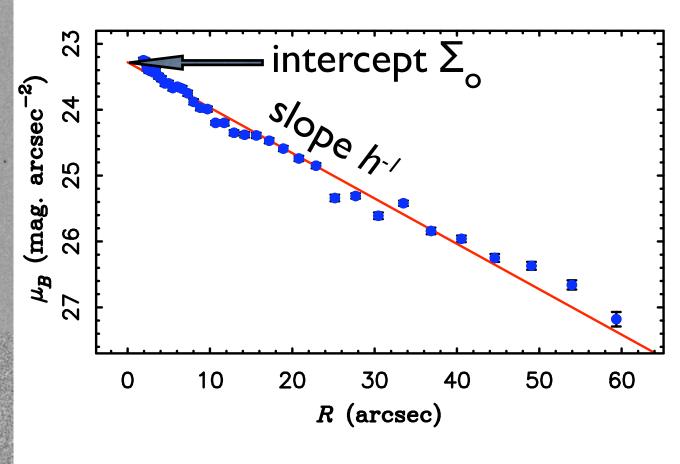
What gets us into trouble is not what we don't know.

It's what we know for sure that just aint so.



- Mark Twain

High Surface Brightness (HSB)

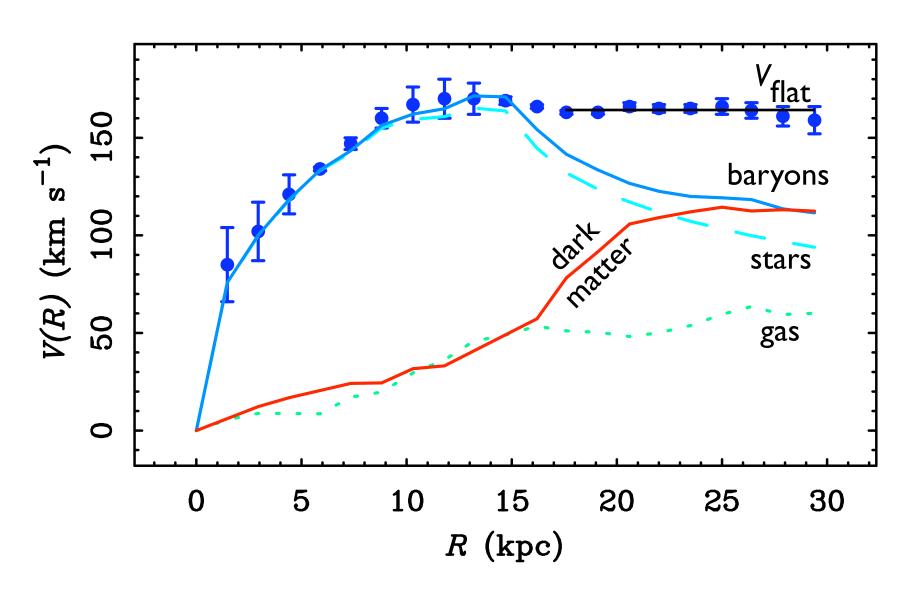


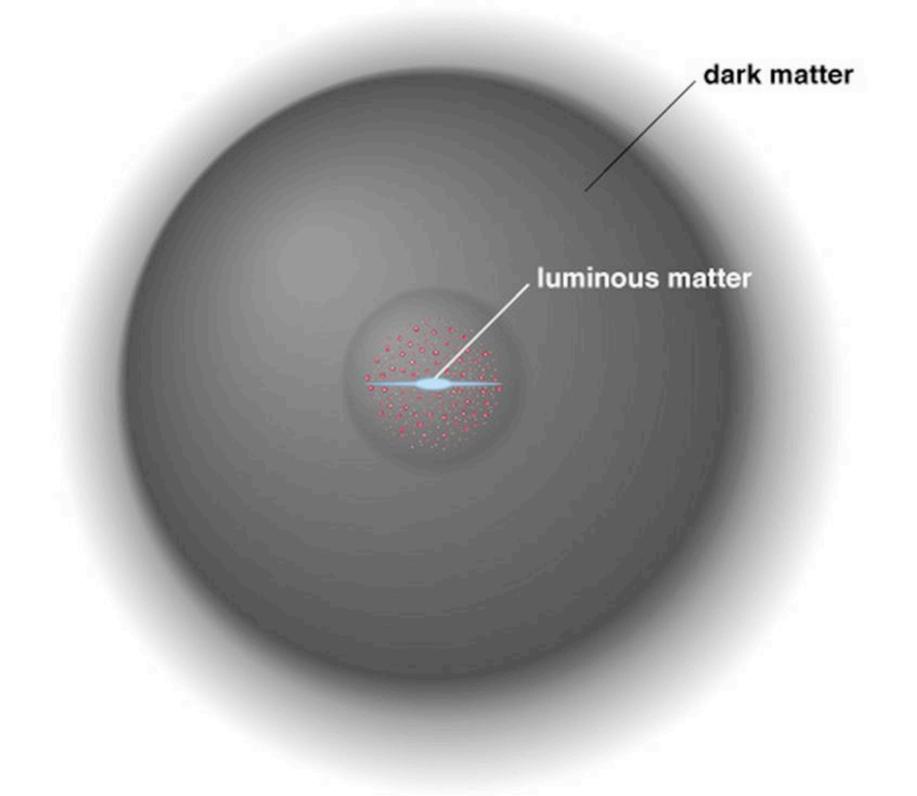
$$\Sigma(R) = \Sigma_o e^{-R/h}$$

Azimuthally averaged light distribution typically exponential for spiral disks.

Low Surface Brightness (LSB)

NGC 6946: $\mathcal{M}_*/L_B = 1.1 \ \mathcal{M}_{\odot}/L_{\odot}$





Dark Matter Halos in Cosmology

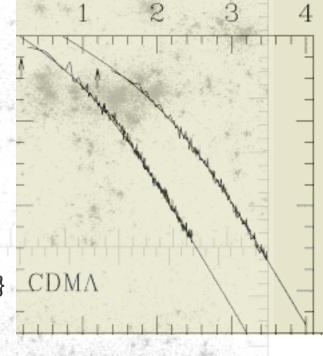
Simulations of structure formation find a characteristic halo form (the "NFW halo"):

$$\rho(r) = \delta_c \rho_c / [(r/r_s)(1+r/r_s)^2]$$

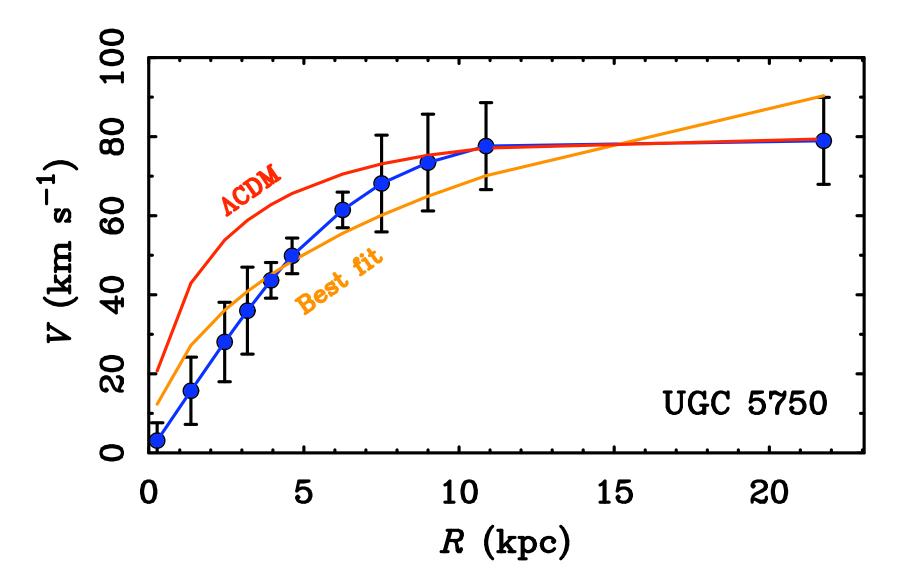
note: $\rho(r) \sim r^{-1}$ at small radii

$$V^{2}(r) = V^{2}_{200}[ln(1+cx)-cx/(1+cx)]/\{x[ln(1+c)-c/(1+c)]\} CDM\Lambda$$

with
$$x = r/r_{200}$$
 and $c = r_{200}/r_{s}$



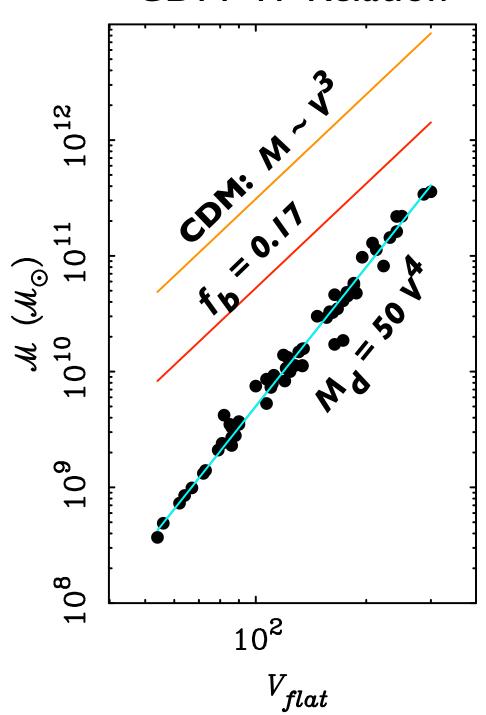
Predicts mass-velocity relation and rotation curve shape.



 Λ CDM predicts: c = 10; $V_{200} = 67$ km/s

Best NFW fit: c = 2.6; $V_{200} = 123 \text{ km/s}$

CDM TF Relation



The **Frenk Principle**: "If the Cold Dark Matter Model does not agree with observations, there must be physical processes, no matter how bizarre or unlikely, that can explain the discrepancy."

The **Strong Frenk Principle**: (2 versions)

Copyright @ Addison Wesler

1: "The physical processes must be the most bizarre and unlikely..."

2: "If we are incapable of finding any physical processes to explain the discrepancy between CDM models and observations, then observations are wrong."

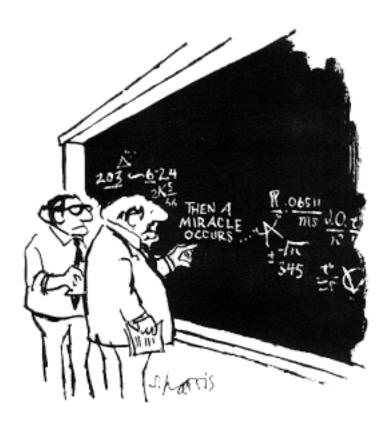
- George Efstathiou

step 1: dark matter halo

step 2: feedback? gastrophysics?

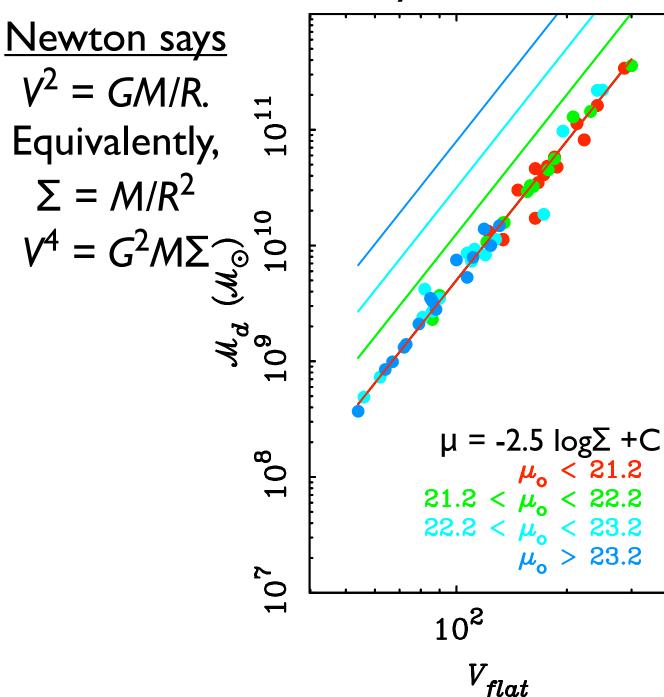
step 3: a real galaxy



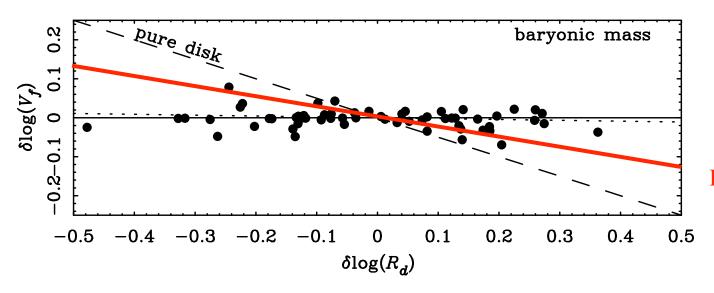


"I think you should be more explicit here in step two."

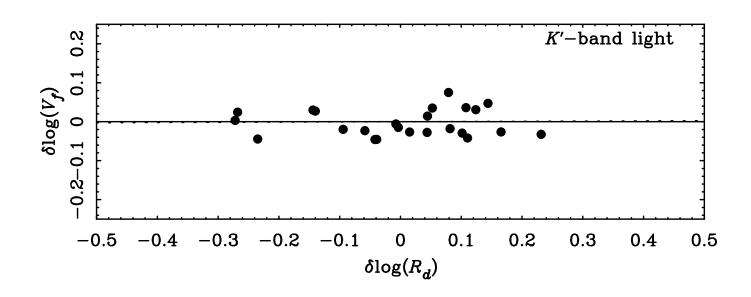
Tully-Fisher Relation

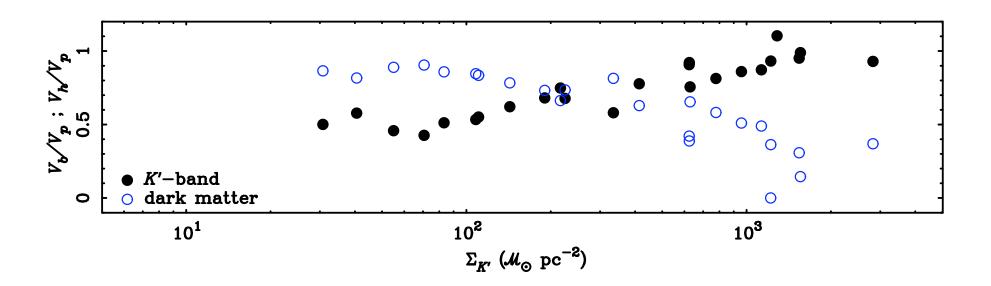


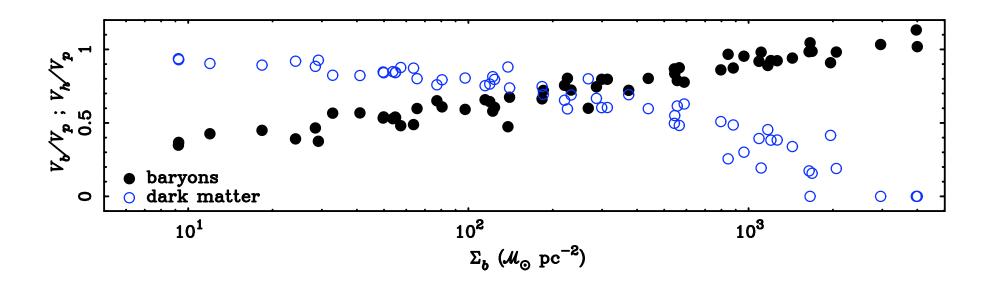
Therefore Galaxies of different surface brightnesses should form distinct Tully-Fisher sequences.



Dutton et al. (2005)

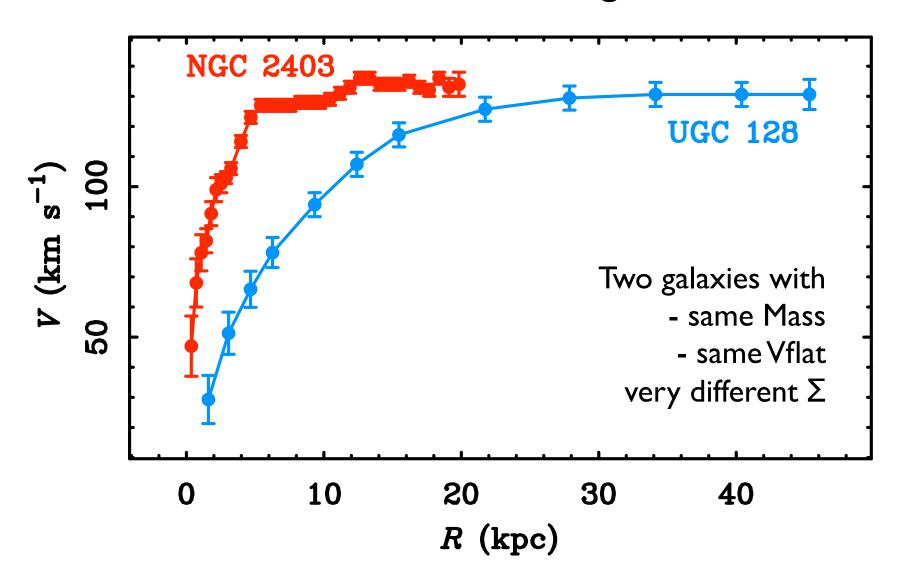




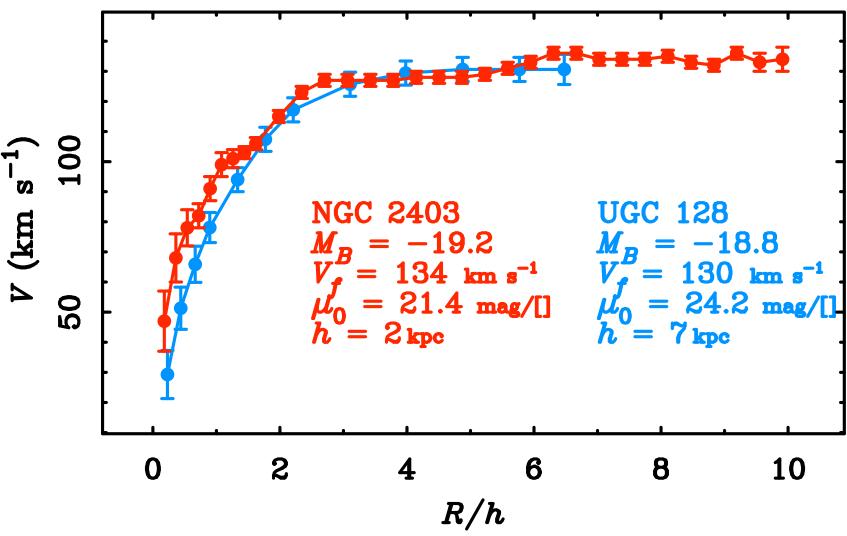


fine-tuning unavoidable

Galaxies which are indistinguishable in TF:



Radius measured by size of visible disk



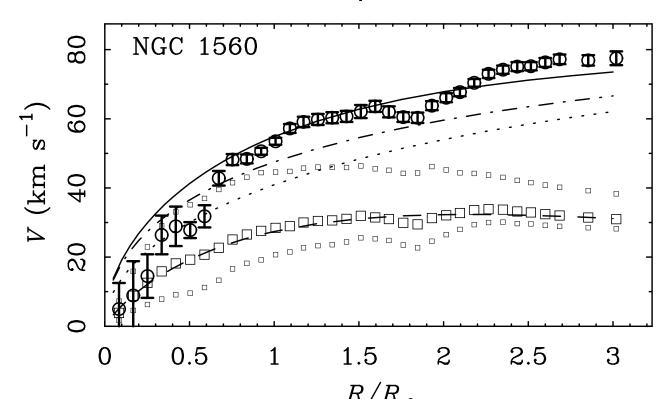
Dynamics knows about the distribution of light as well as the total mass.

Renzo's Rule:

"When you see a feature in the light, you see a corresponding feature in the rotation curve."

(Sancisi 1995, private communication 2003, published in IAU proceedings)

The distribution of mass is coupled to the distribution of light.



<u>MOND</u>

MOdified Newtonian Dynamics

introduced by Moti Milgrom in 1983



instead of dark matter, suppose the force law changes such that

$$\mu(a/a_o) a = g_N$$
.

Above a critical acceleration a_0 everything is normal. Below that scale, gravity in effect becomes stronger.

$$a_0 = 1.2 \times 10^{-10} \text{ m s}^{-2}$$

Milgrom 1983

No. 2, 1983

MODIFICATION OF NEWTONIAN DYNAMICS

38

A major step in understanding ellipticals can be made if we can identify them, at least approximately, with idealized structures such as the FRCL spheres discussed above. I have also studied isotropic and nonisotropic isothermal spheres, in the modified dynamics, as such possible structures. I found that they have properties which the properties which the properties is the properties of compticals and galactic being I to set it the property for the properties.

VIII. PREDICTIONS

The main predictions concerns and or fi-

- 1. Velocity curves calculate with the modified dynamics on the basis of the observed mass in galaxies should agree with the observed curves. Elliptical and SO galaxies may be the best for this purpose since (a) practically no uncertainty due to obscuration is involved and (b) there is not much uncertainty due to the possible presence of molecular hydrogen.
- 2. The relation between the asymptotic velocity (V_{∞}) and the mass of the galaxy (M) $(V_{\infty}^4 = MGa_0)$ is an absolute one.
- 3. Analysis of the z-dynamics in disk galaxies using the modified dynamics should yield surface densities which yer with the base of one of dilingly the same in a situate of the carentiol trying mics standy yield a discrepancy which increases with radius in a predictable manner.
- 4. Effects of the n be particularly stro g review of propert's s. 1980). For example, those dwarfs believed to be bound to our Galaxy would have internal accelerations typically of order $a_{in} \sim a_0/30$. Their (modified) acceleration, g, in the field of the Galaxy is larger than the internal ones but still much smaller than a_0 , $g \approx (8)$ kpc/d) a_0 , based on a value of $V_{\infty} = 220 \text{ km s}^{-1}$ for the Galaxy, and where d is the distance from the dwarf galaxy to the center of the Milky Way ($d \sim 70-220$ kpc). Whichever way the external acceleration turns out to affect the internal dynamics (see the discussion at the end of § II, the section on small groups in Paper III, and Paper I), we predict that when velocity dispersion data is available for the dwarfs, a large mass discrepancy will result when the conventional dynamics is used to determine the masses. The dynamically determined mass is predicted to be larger by a factor of order 10 or more than that which can be accounted for by stars. In case the internal dynamics is determined by the external acceleration, we predict this factor to increase with dand be of order (d/8 kpc) (as long as $a_{in} \ll g$, $h_{50} = 1$).

Prediction 1 is a very general one. It is worthwhile listing some of its consequences as separate predictions, numbered 5-7 below (note that, in fact, even prediction 2 is already contained in prediction 1).

- 5. Measuring local M/L values in disk galaxies (assuming conventional dynamics) should give the following results: In regions of the galaxy where $V^2/r \gg a_0$ the local M/L values should show no indication of hidden mass. At a certain transition radius, local M/L should start to increase rapidly. The transition radius should start to increase rapidly. The transition radius should occur where $W/r \approx 10^{-4} L/L$ as the fraction of W/L as we are concerned only with variations of this quantity; (b) Effects of the modified dynamics manifest themselve times clearly in local W/L as minators in the first equation of those which is equipped to the modified dynamics manifest themselve times clearly in local W/L as the equipped W/L as the matter of the modified dynamics manifest themselve times clearly in local W/L as the equipped W/L and W/L as the equipped W/L and W/L as the equipped W/L as the matter of the modified dynamics manifest themselve times clearly in local W/L as the equipped W/L as the matter W/L and W/L as the equipped W/L and W/L as the equipped W/L as the matter W/L and W/L as the equipped W/L as the following W/L and W/L as the following W/L and W/L as the following W/L as the following W/L and W/L as the following W/L as the following W/L and W/L as the following W/L as the following W/L and W/L and W/L are the following W/L and W/L and W/L are the following W/L
- 6. Disk galaxies with low surface brightness provide particularly strong tests (a study of a sample of such galaxies is described by Strom 1982 and by Romanishin et al. 1982). As low surface brightness means small accelerations, the effects of the modification should be more noticeable in such galaxies. We predict, for example, that the proportionality factor in the $M \propto V_{\perp}^4$ relation for these galaxies is the same as for the high surface density g_{\perp}^2 xies. In to trast, if one wards to obtain a continuous surface lensity g_{\perp}^2 (i.e. for example, Aaronson, Huchra, and would 1979), where Σ is the average surface brightness. This implies that low surface lensy galaxies, of g_{\perp}^2 with the continuous surface lensy galaxies, of g_{\perp}^2 and g_{\perp}^2 in the average surface brightness. This implies that low surface lensy galaxies, of g_{\perp}^2 and g_{\perp}^2 is the average surface brightness. This implies that low surface lensy galaxies of g_{\perp}^2 is the average surface brightness of g_{\perp}^2 in the lower than the lower than

density of a galaxy is, the smaller in the transition rate $S_{\rm c}$ and $S_{\rm c}$ and $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate $S_{\rm c}$ and $S_{\rm c}$ are rate $S_{\rm c}$ are rate

7. As the study of model rotation curves shows, we predict a correlation between the value of the average surface density (or brightness) of a galaxy and the steepness with which the rotational velocity uses to its asymptotic value (as measured, for example, by the radius at which $V = V_{\infty}/2$ in units of the scale length of the disk). Small surface densities imply slow rise of V.

IX. DISCUSSION

The main results of this paper can be summarized by the statement that the modified dynamics eliminates the need to assume hidden mass in galaxies. The effects in galaxies which I have considered, and which are commonly attributed to such hidden mass, are readily explained by the modification. More specifically:

MOND predictions

• The Tully-Fisher Relation

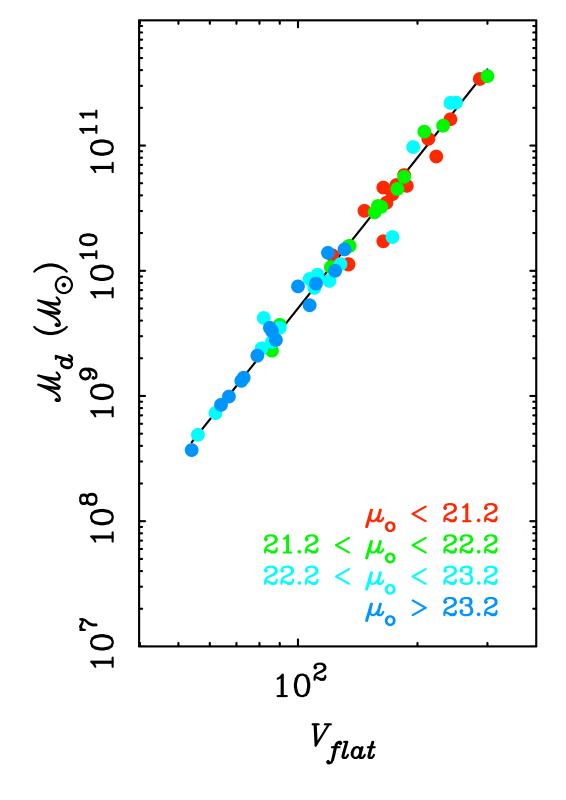
y strongmental stream of the strong and V strongmental stream of the stream of the stream of the stream of the strong and V strong and

• No Dependence on Surface Brightness

radius and surface brightness

• Rotation Curve Shapes
• to exist.

- Surface Density ~ Surface Brightness
- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios



• The Tully-Fisher Relation

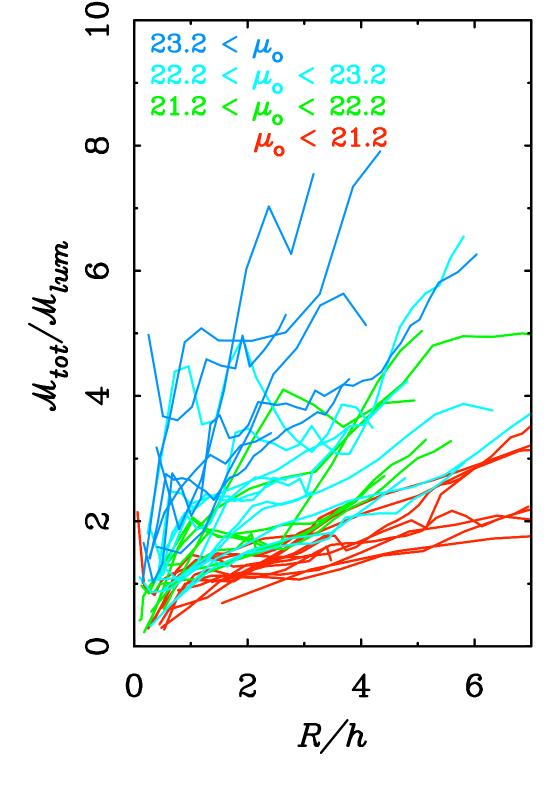
 $label{eq:Slope}$ Slope = 4

Normalization = $1/(a_0G)$

Fundamentally a relation between Disk Mass and V_{flat}

No Dependence on Surface Brightness

- Dependence of conventional M/L on radius and surface brightness
- Rotation Curve Shapes
- Surface Density ~ Surface Brightness
- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios



• The Tully-Fisher Relation

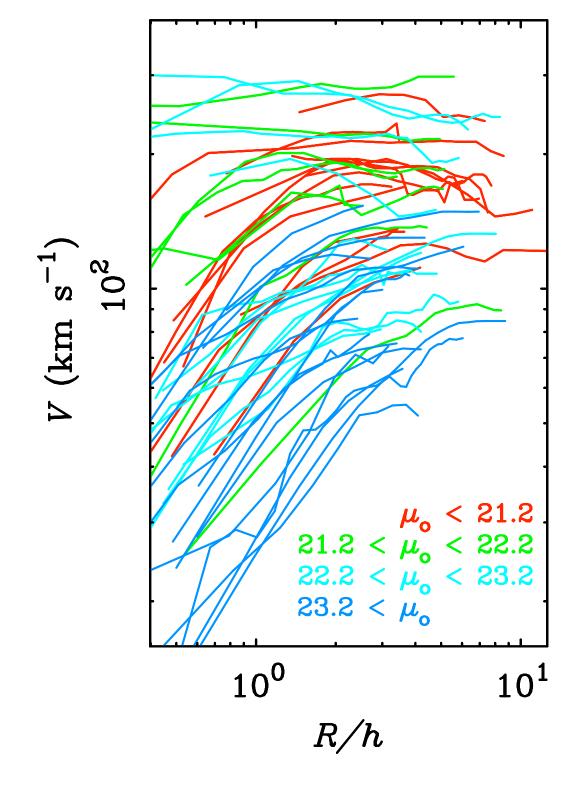
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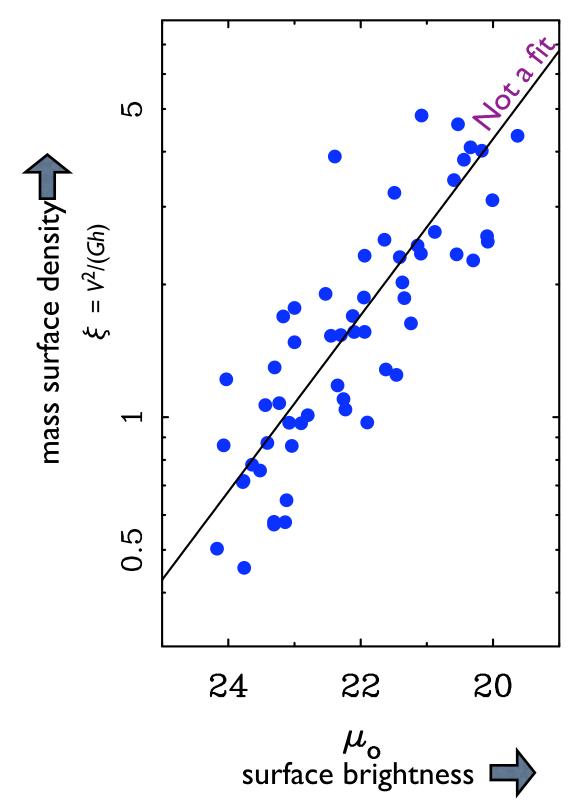
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No Dependence on Surface Brightness

Dependence of conventional M/L on radius and surface brightness

Rotation Curve Shapes

- Surface Density ~ Surface Brightness
- Detailed Rotation Curve Fits
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Normalization = $1/(a_0G)$

Fundamentally a relation between Disk Mass and V_{flat}

No Dependence on Surface Brightness

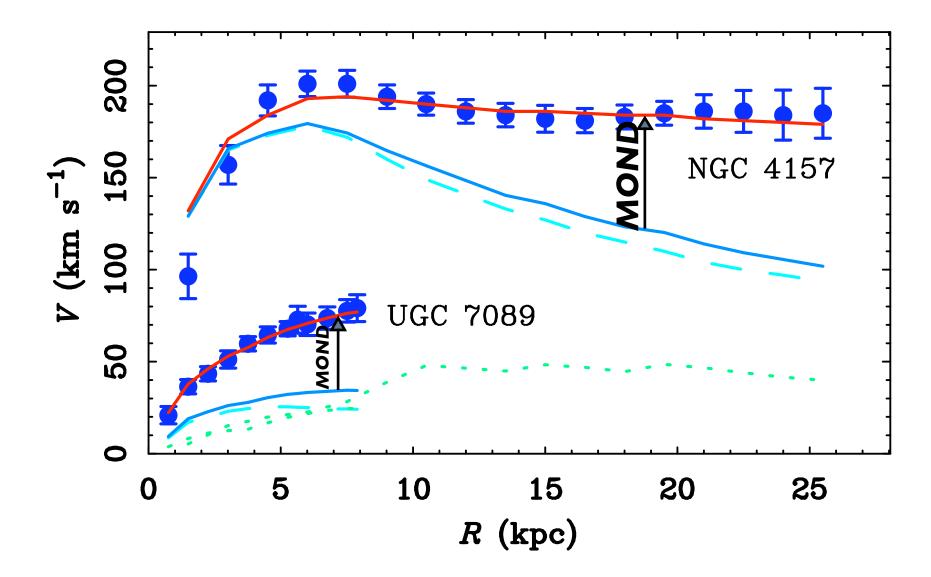
Dependence of conventional M/L on radius and surface brightness

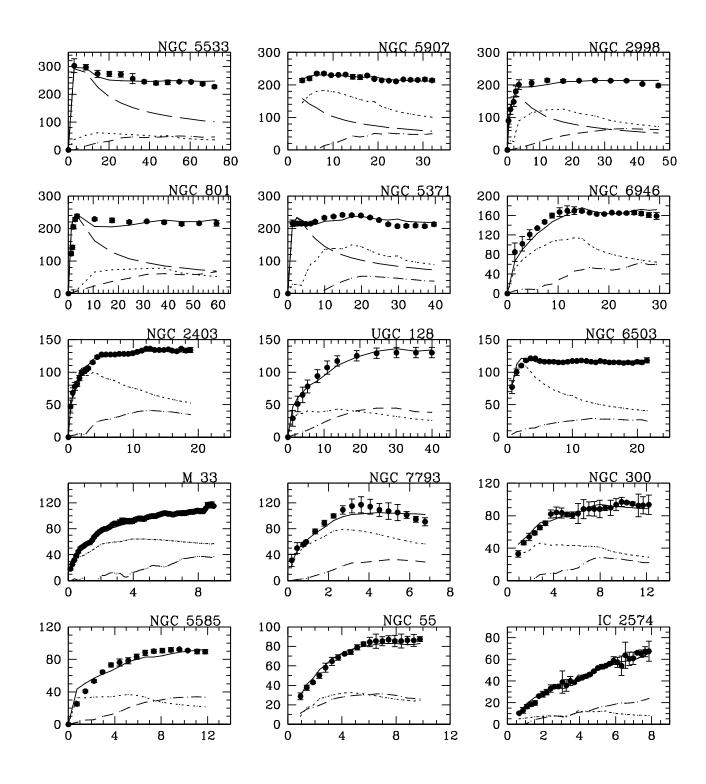
Rotation Curve Shapes

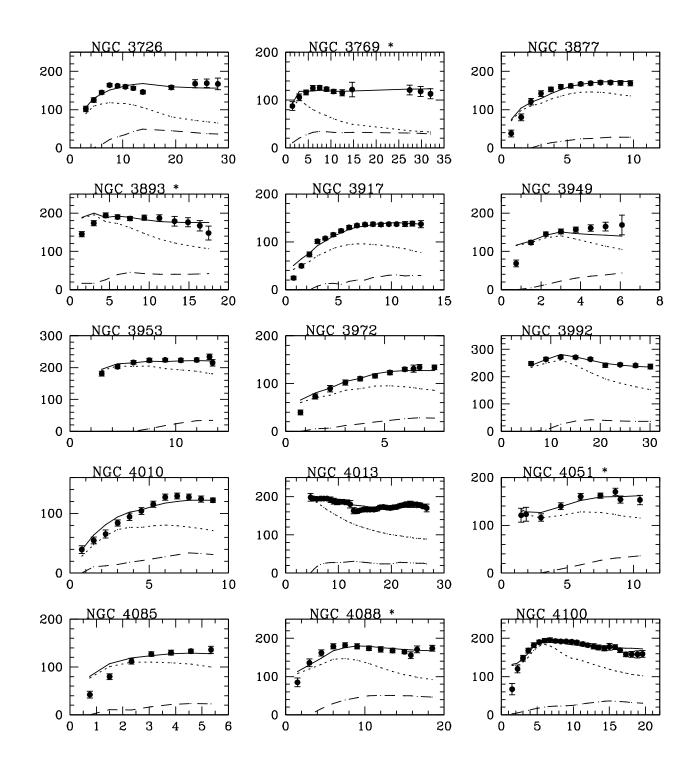
Surface Density ~ Surface Brightness

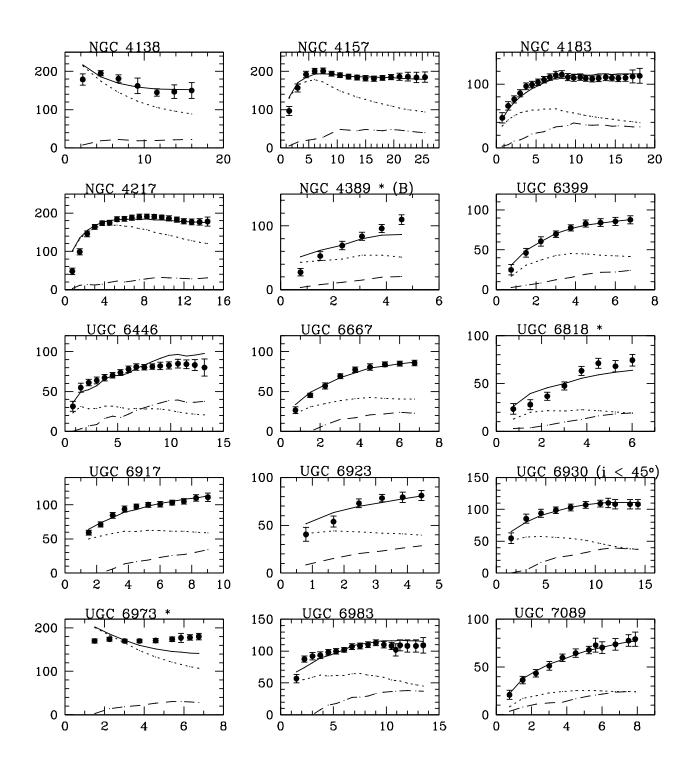
Detailed Rotation Curve Fits

• Stellar Population Mass-to-Light Ratios

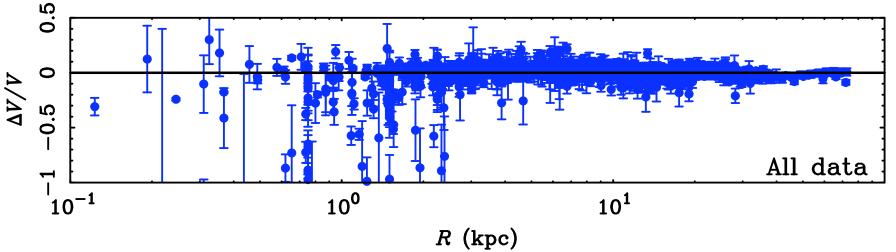




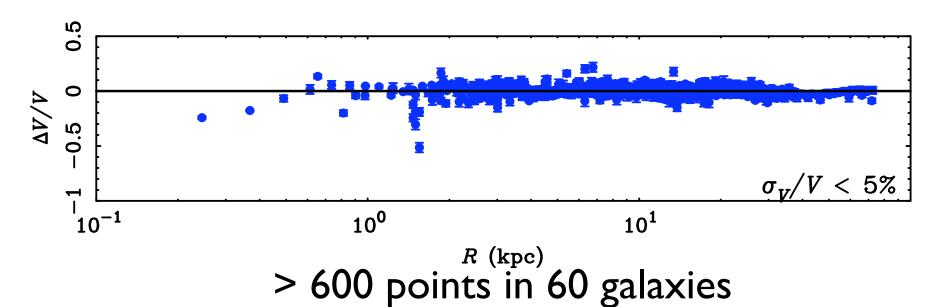




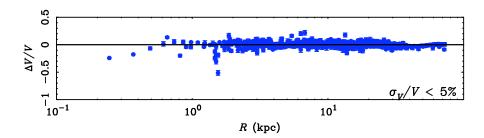
Residuals of MOND fits



> 1000 points in 74 galaxies



All data R (kpc)



MOND predictions

• The Tully-Fisher Relation

Slope = 4

Normalization = $1/(a_0G)$

Fundamentally a relation between Disk Mass and V_{flat}

No Dependence on Surface Brightness

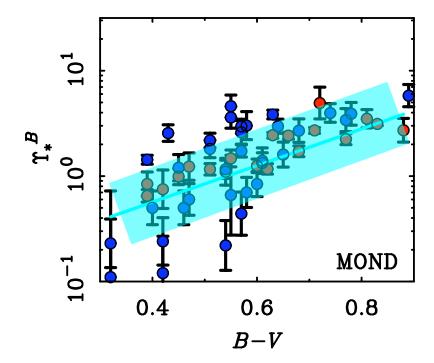
Dependence of conventional M/L on radius and surface brightness

Rotation Curve Shapes

Surface Density ~ Surface Brightness

Detailed Rotation Curve Fits

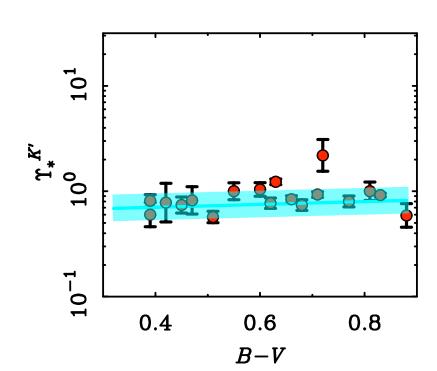
• Stellar Population Mass-to-Light Ratios



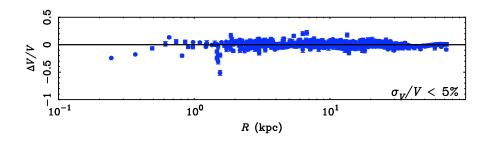


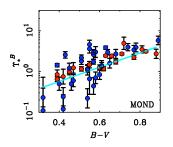


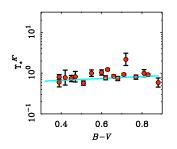
- mean normalization
 - trend with color
 - expected scatter



All data R (kpc)







MOND predictions

• The Tully-Fisher Relation

Slope = 4

Normalization = $1/(a_0G)$

Fundamentally a relation between Disk Mass and V_{flat}

No Dependence on Surface Brightness

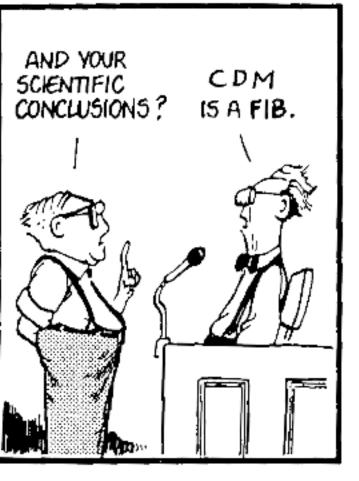
Dependence of conventional M/L on radius and surface brightness

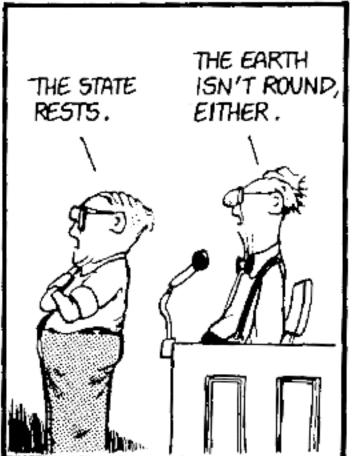
Rotation Curve Shapes

Surface Density ~ Surface Brightness

Detailed Rotation Curve Fits

Stellar Population Mass-to-Light Ratios







Everything happens... as if the force between two bodies is directly proportional to the product of their masses and inversely proportional to the square of the distance between them.

In spiral galaxies, everything happens as if

$$\mu (a/a_o) a = g_N$$
.



Observed behavior is simple - seems unlikely to originate fomr complicated baryonic gastrophysics





Are you suggesting there is no dark matter?

Distinguishing between dark matter and MOND is a nightmare

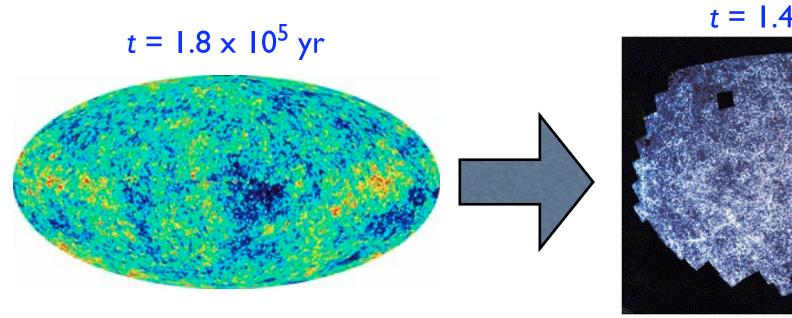
- CDM is the ultimate accommodator
 - can (and does) explain just about anything
 - not explicitly falsifiable
- MOND makes very specific predictions (delta-fcn priors)
 - about some things
 - is completely mute about others

What does MOND do to cosmology?

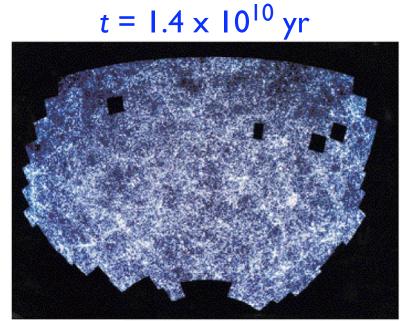


In MOND

$$\Omega_b \approx \Omega_m$$

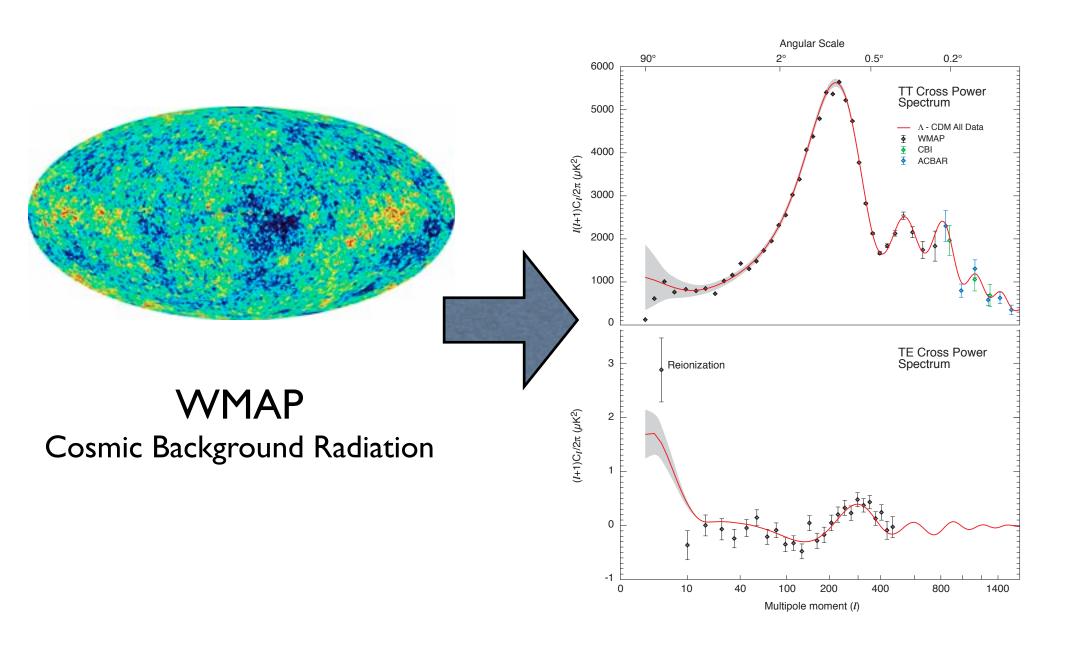


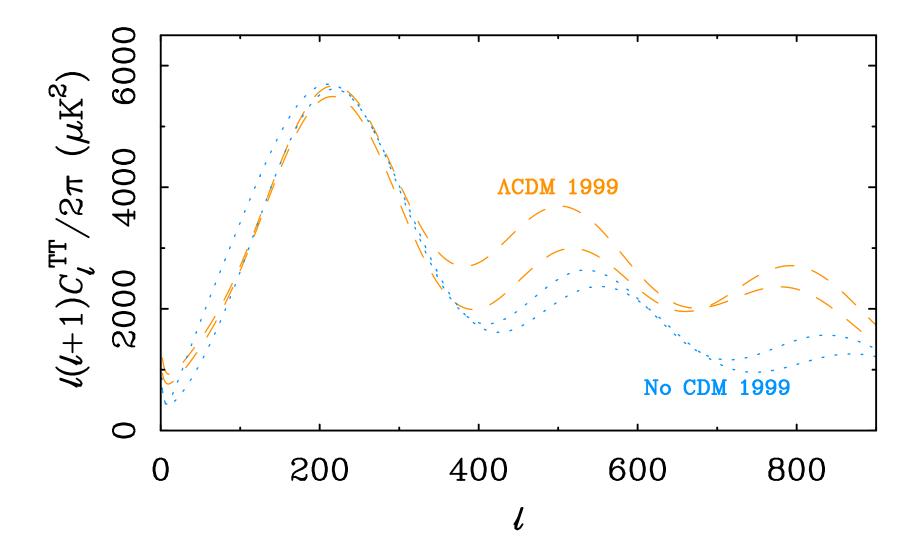
very smooth: $\delta \rho / \rho \sim 10^{-5}$

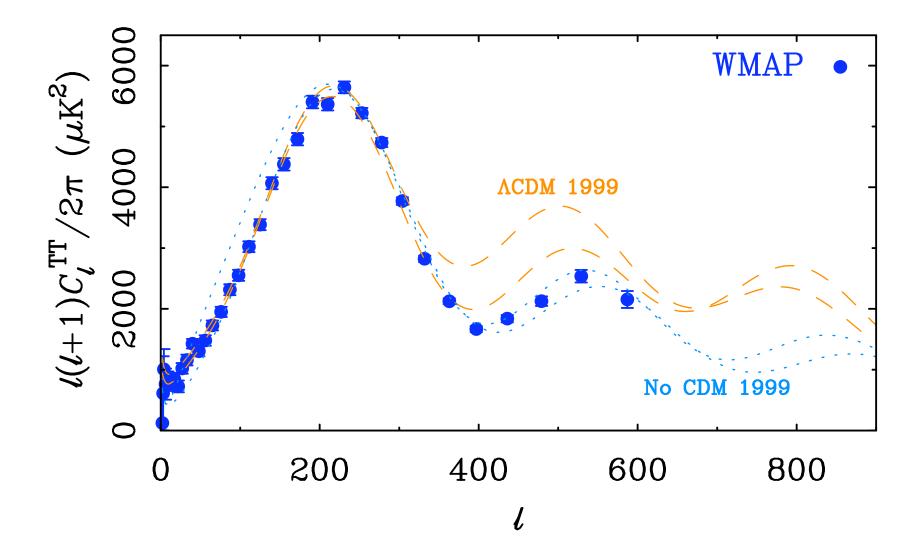


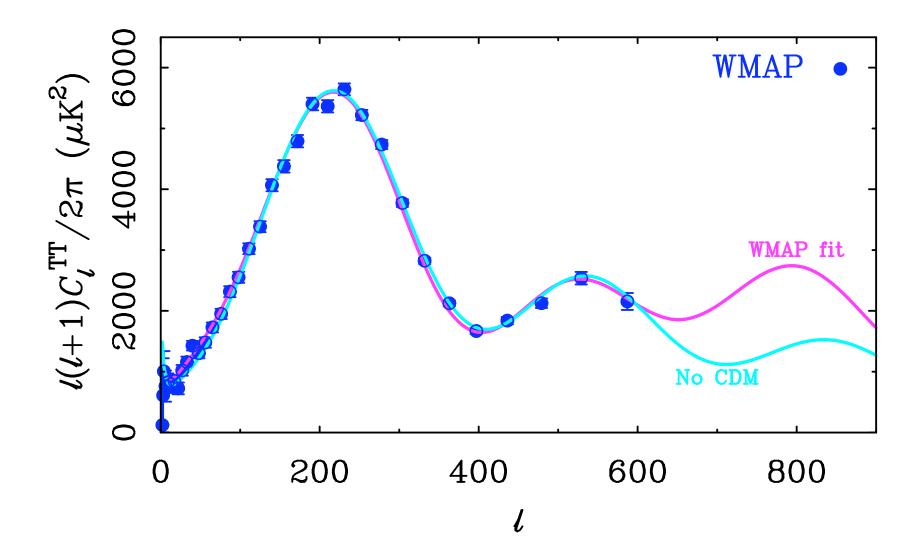
very lumpy: $\delta \rho / \rho \sim 1$

 $\delta \rho / \rho \propto \text{faster than } t^{2/3}$

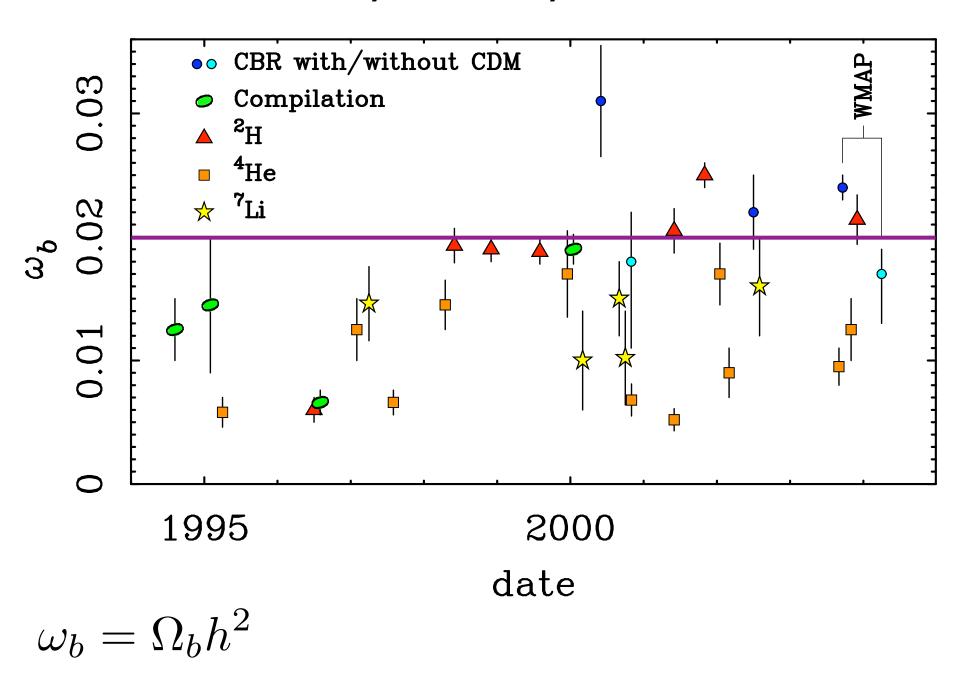


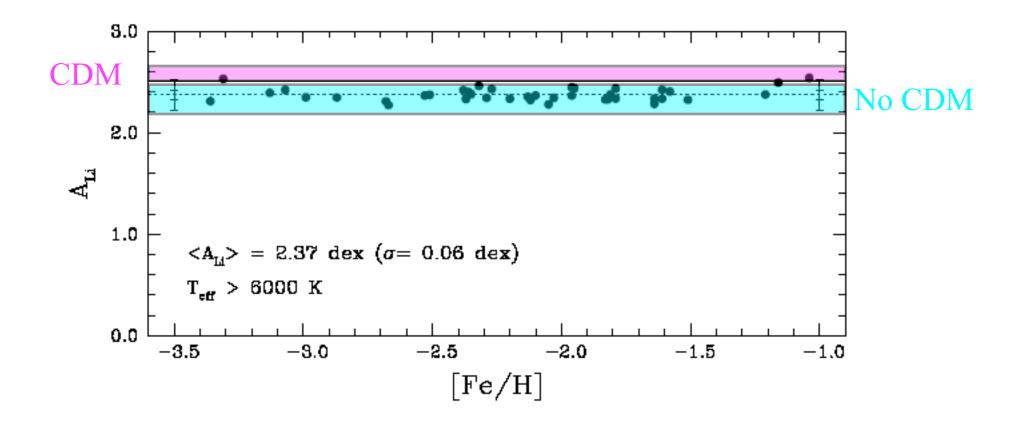




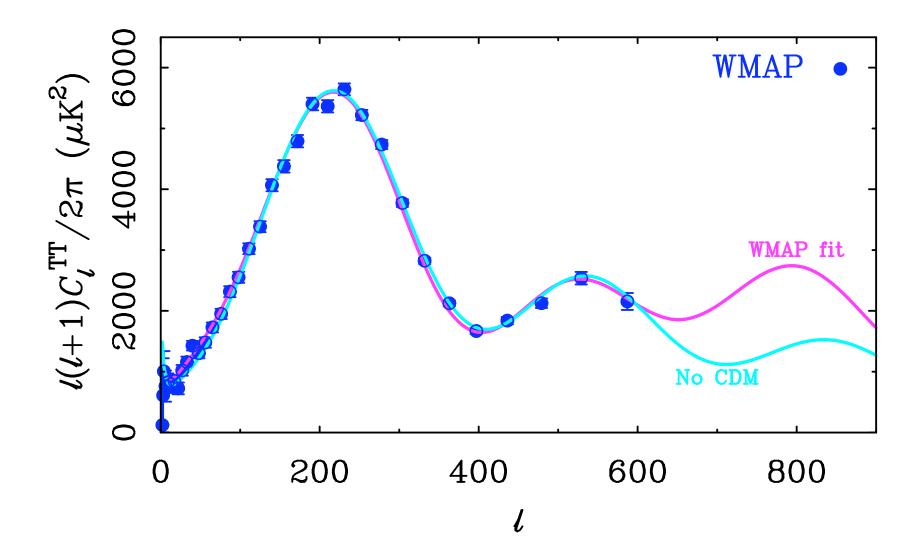


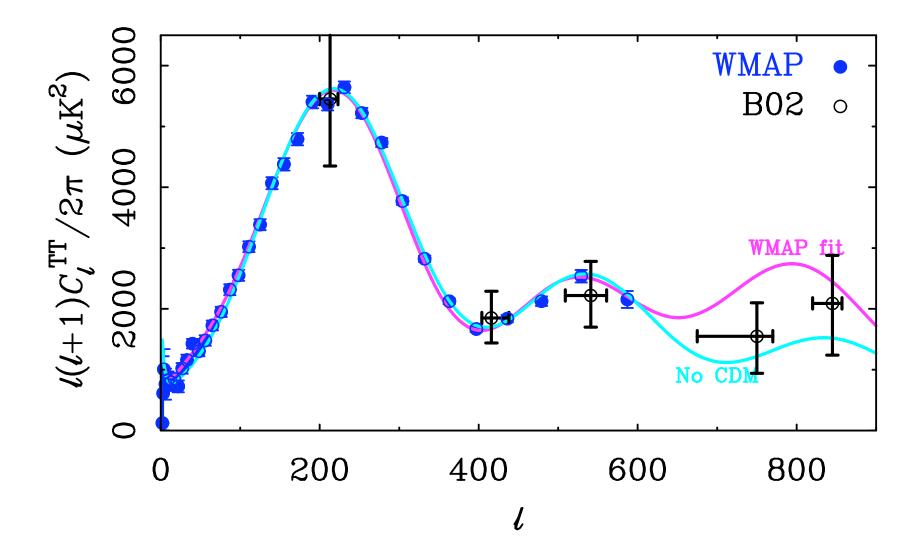
Baryon density from BBN

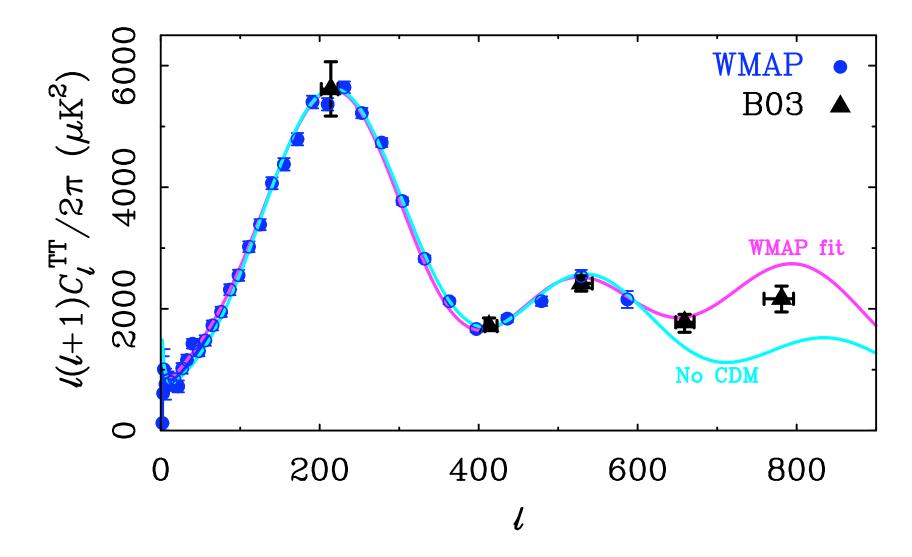




Melendez & Ramirez Astrophys.J. 615 (2004) L33







Ways Out

- Detect Dark Matter in the laboratory
- Third peak of CBR power spectrum
- Neutrino mass (more/less than ~0.5 eV)
- Find systems MOND can't fit (clusters?)

TeVeS: Relativistic MOND?

"I think Bekenstein's March 04
paper changed the intellectual landscape. The
classic objection to McGaugh etc, that MOND is not a proper theory, is now
void, and it is imperative that we discover whether Milgrom's theory is
compatible with (a) CMB and (b) LSS. In view of the non-zero value of
Lambda, which plainly states that minimal GR won't do, I consider that MOND
now has the intellectual edge & should be taken very seriously indeed.
The only problem is that I have yet to encounter someone who has taken up the
challenge that Bekenstein has posed. This I feel reflects worse on the field than on
MOND."

James Binney, August 3, 2004