

*Cosmic Flights of Fancy:
Mass and the Gravity that Moves It*

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University of Maryland*



*What gets us into trouble is not
what we don't know.*

*It's what we know for sure that just
aint so.*

- Mark Twain



A few things we know for sure...

$$\nabla^2\Phi = 4\pi G\rho$$

$$F = ma$$

which basically means

$$mV^2/R = GMm/R^2$$

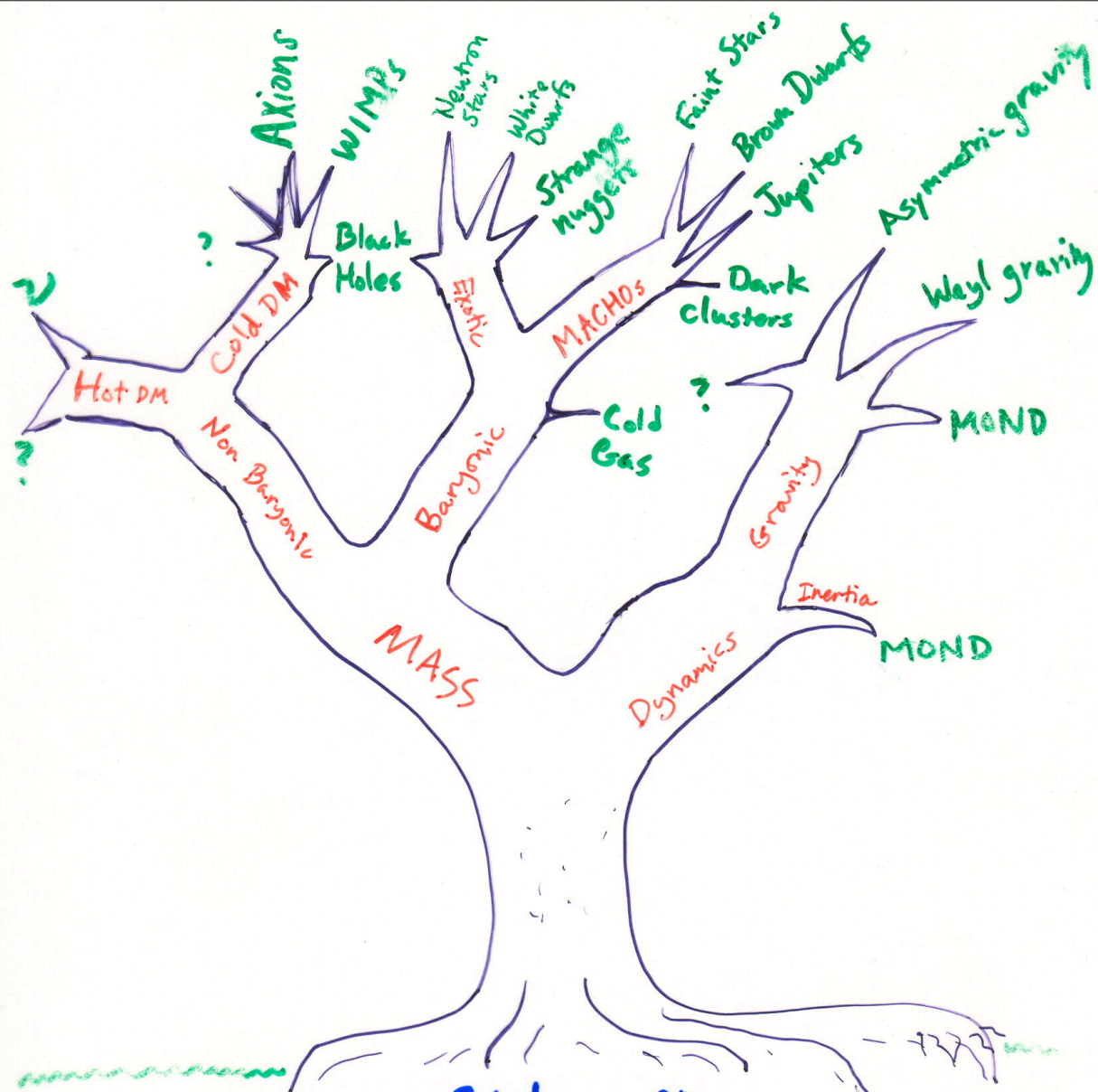
i.e.,

$$V^2 = GM/R$$

ergo...

The universe is filled with nonbaryonic cold dark matter.





Disk DM
Oort
discrepancy

Spiral
galaxy
flat
rotation
curves

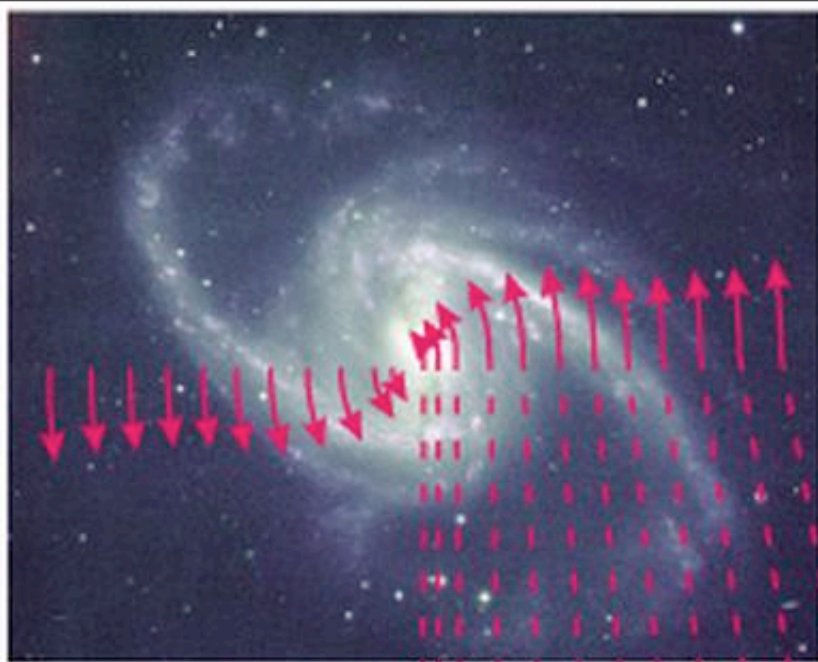
$$\frac{M_{HI}}{M_T} \approx 0.1$$

Cluster
Velocity
dispersions

$$\frac{M_L}{M_T} \approx 300$$

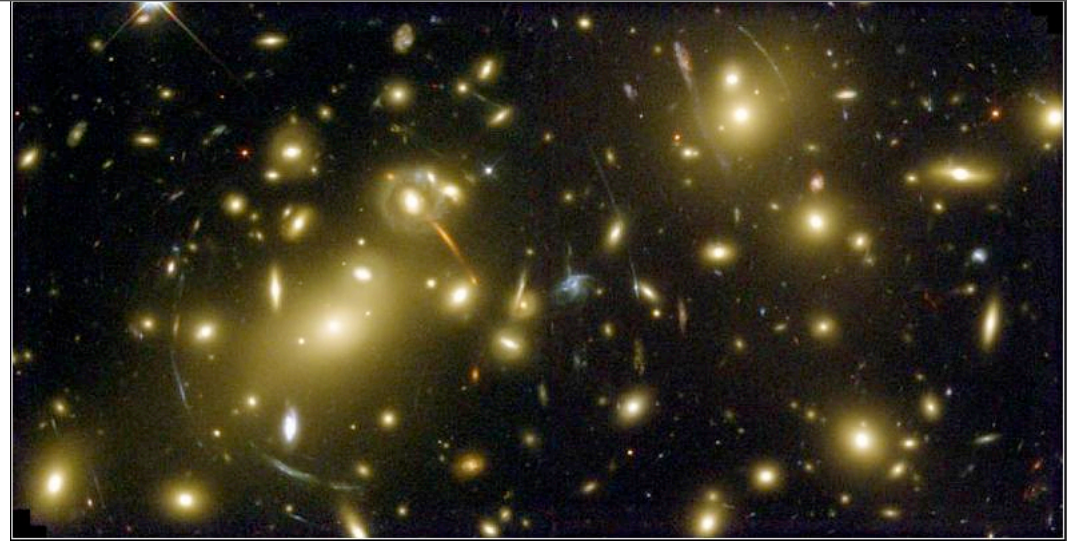
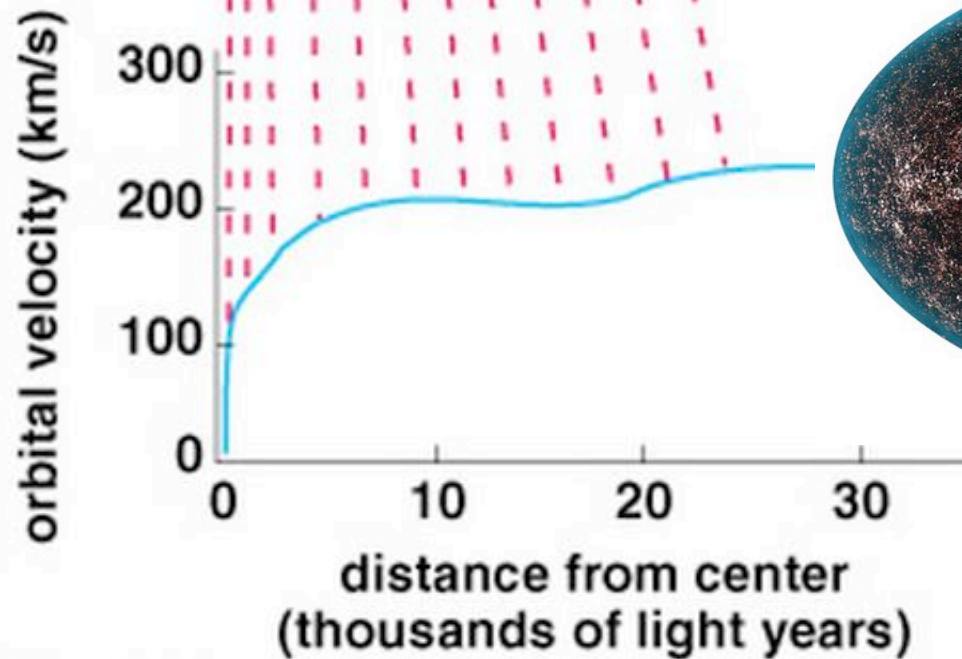
$$\frac{M_X}{M_T} \approx 0.2$$

$\Omega = 1$
Large
Scale
Structure
Bulk
flows



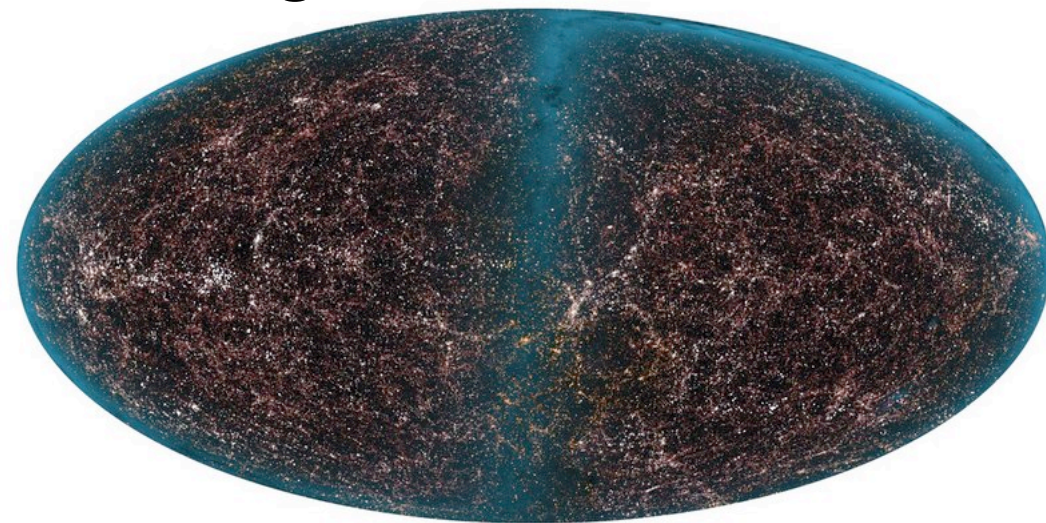
Longer arrows
represent larger
orbital velocities.

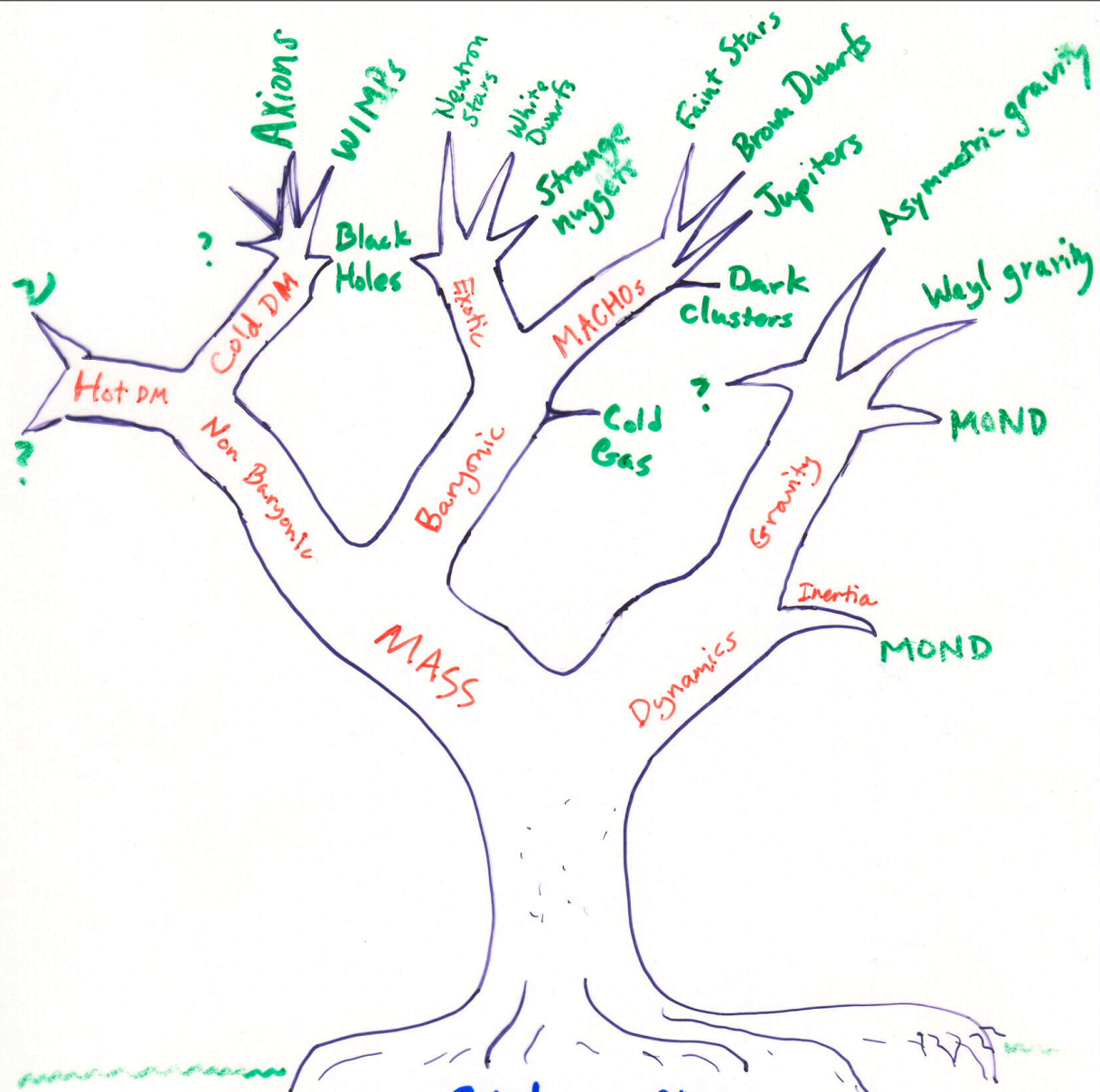
Galaxy Rotation Curves



Galaxy Clusters

Large Scale Structure





Disk DM
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Spiral
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$$\frac{M_{HI}}{M_T} \approx 0.1$$

Cluster
Velocity
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$$\frac{M_L}{M_T} \approx 300$$

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$\Omega = 1$
Large
Scale
Structure
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flows

What is the Dark Matter?

Baryonic Dark Matter

Normal things:

very faint stars, brown dwarfs
other hard-to-see objects (planets, gas)

Hot Dark Matter

neutrinos - got mass, but not enough



Cold Dark Matter

Some new fundamental particle

doesn't interact with light, so quite invisible.

Two big motivations:

- 1) total mass outweighs normal mass from BBN
- 2) needed to grow cosmic structure

(I)

Normal baryonic mass = 4% of total
from Primordial Nucleosynthesis

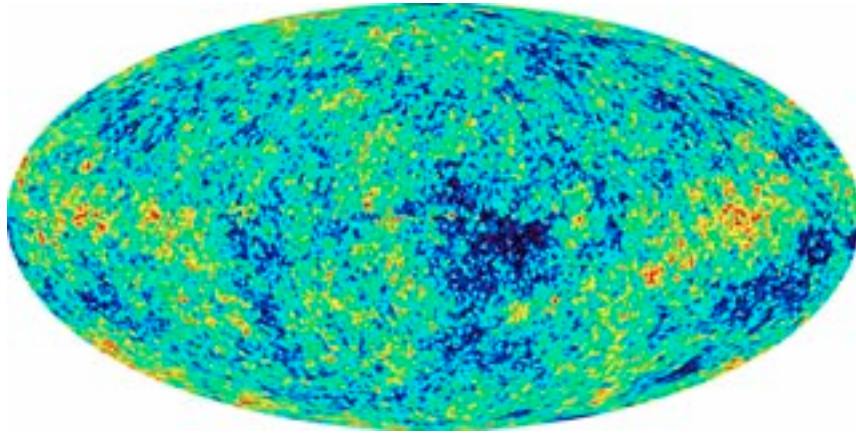
Total mass density = 27% of total
from gravity

gravitating mass \gg normal mass

Most of the mass needs to be
in some brand new form!

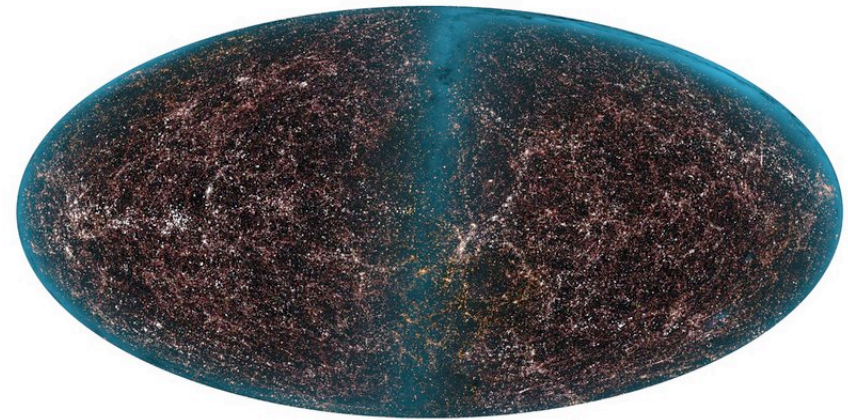
(2) There isn't enough time to form the observed cosmic structures from the smooth initial conditions unless there is a component of mass independent of photons.

$$t = 1.8 \times 10^5 \text{ yr}$$



very smooth: $\delta\rho/\rho \sim 10^{-5}$

$$t = 1.4 \times 10^{10} \text{ yr}$$



very lumpy: $\delta\rho/\rho \sim 1$

$$\delta\rho/\rho \propto t^{2/3}$$

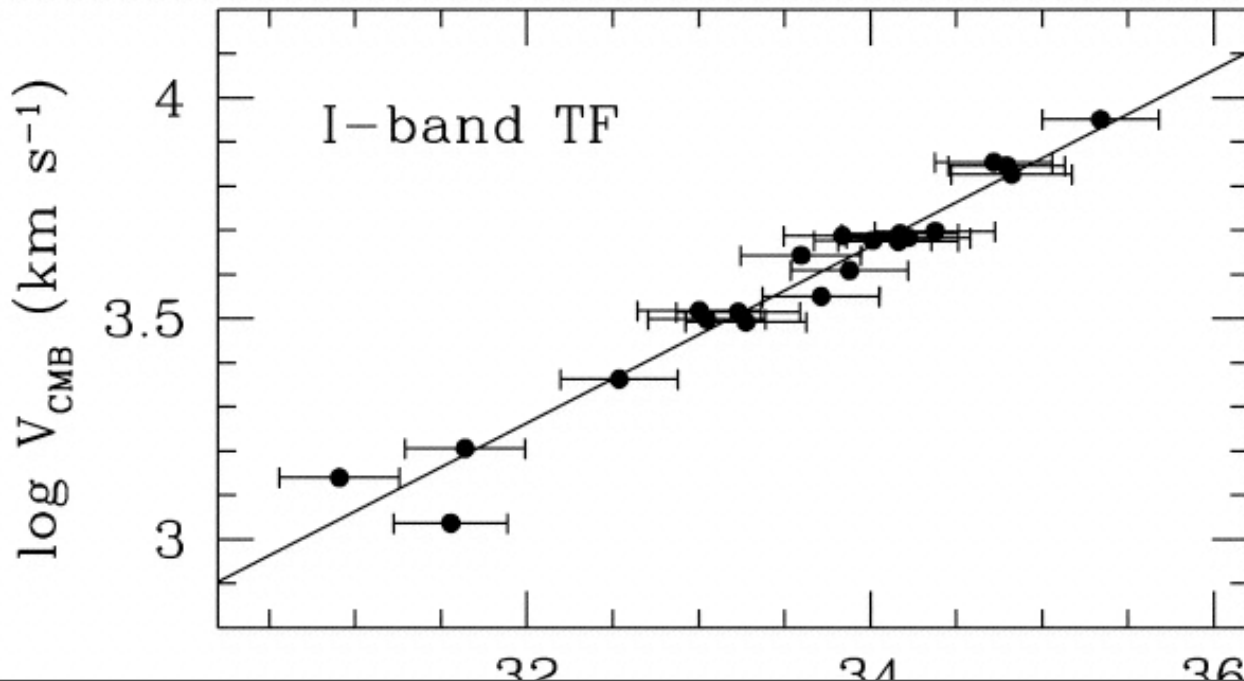
Both (1) and (2) hold only when gravity is normal.

Dark Energy

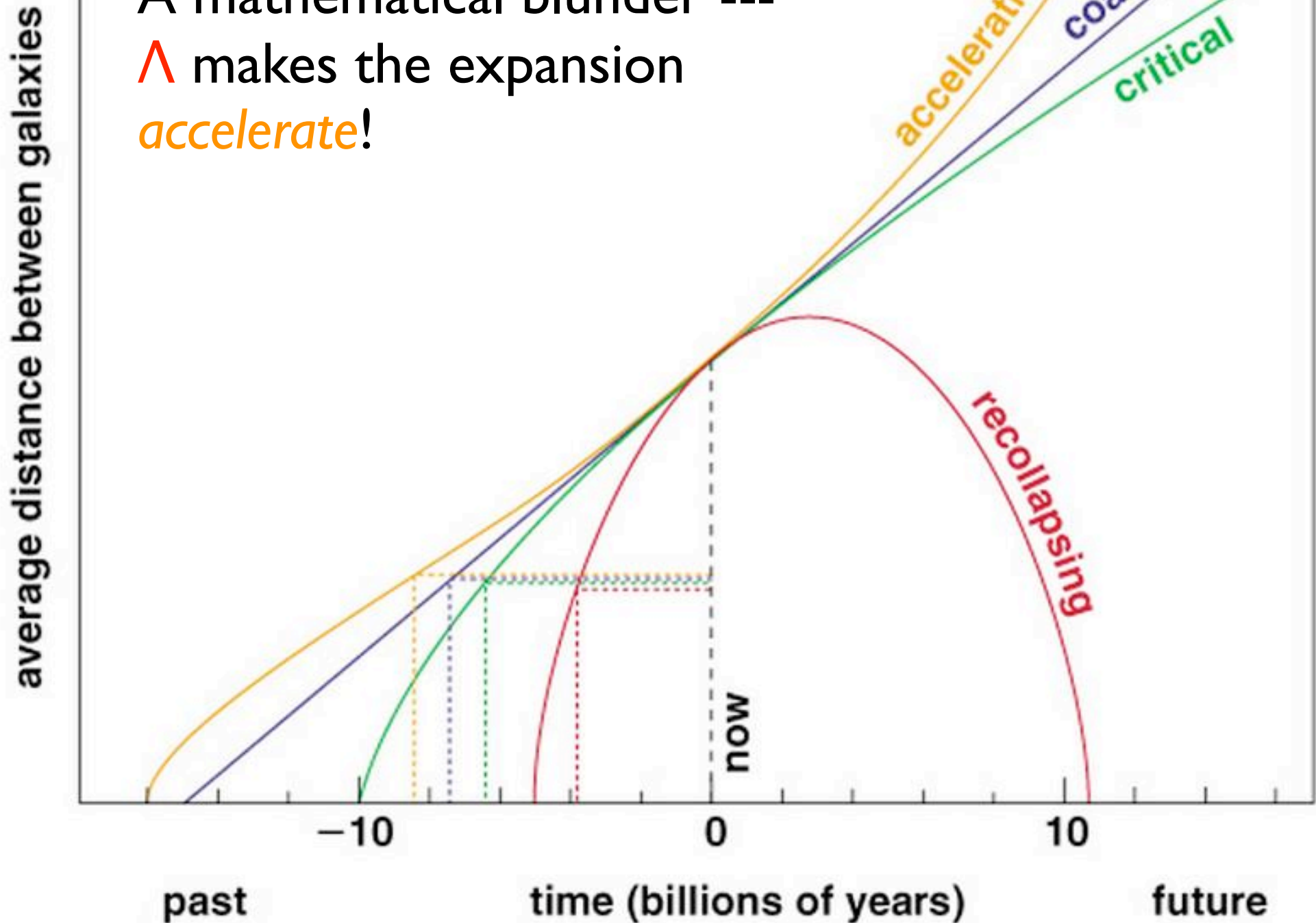
Einstein's greatest blunder?

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu} = 8\pi GT_{\mu\nu} + \Lambda g_{\mu\nu}$$

Einstein's intention was to keep the universe static. But it does expand!



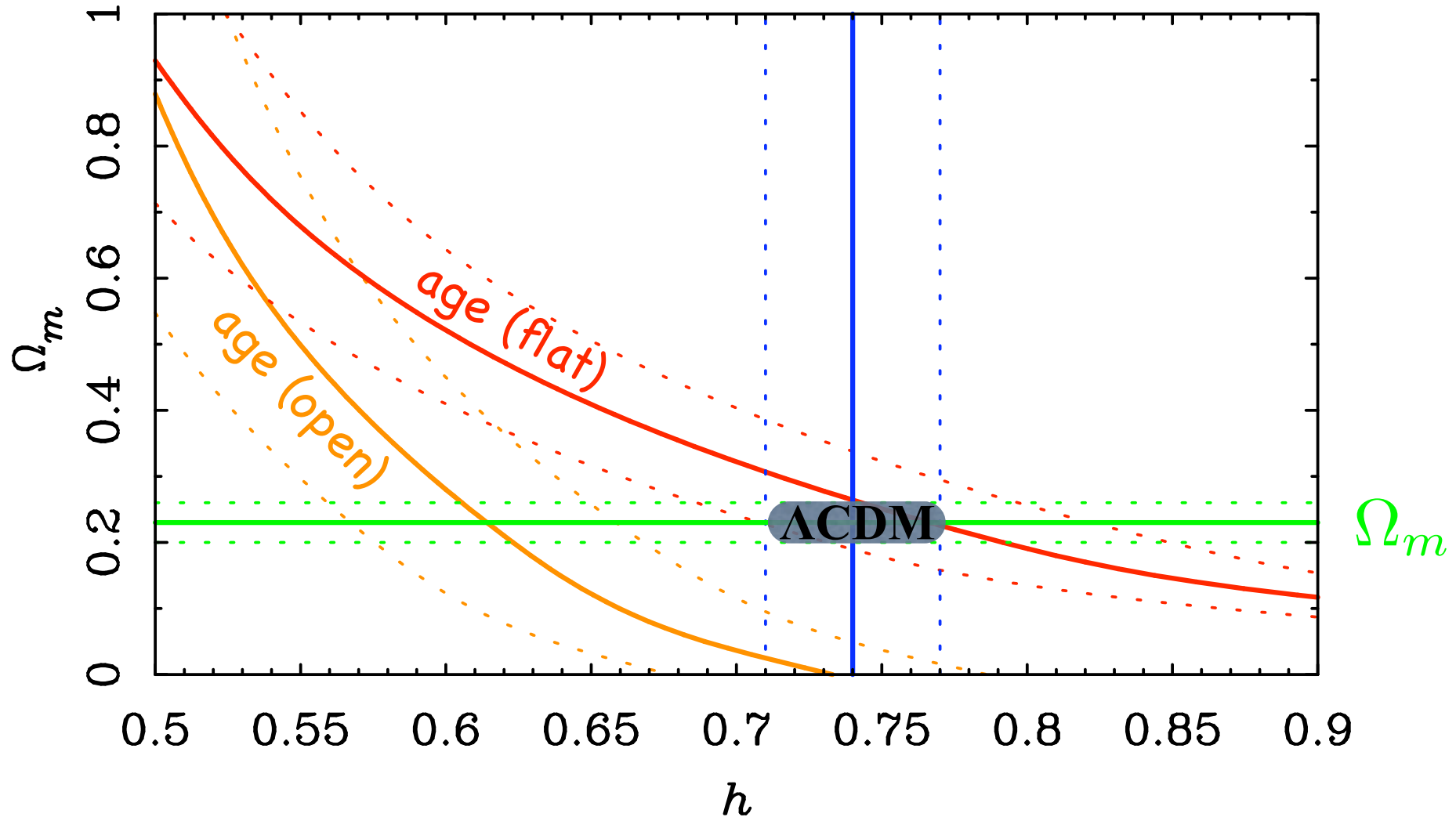
A mathematical blunder ---
Λ makes the expansion
accelerate!



The Search for Two (or Three) Numbers Expansion Rate & Mass Density (& Λ)

age = 13 ± 1 Gyr

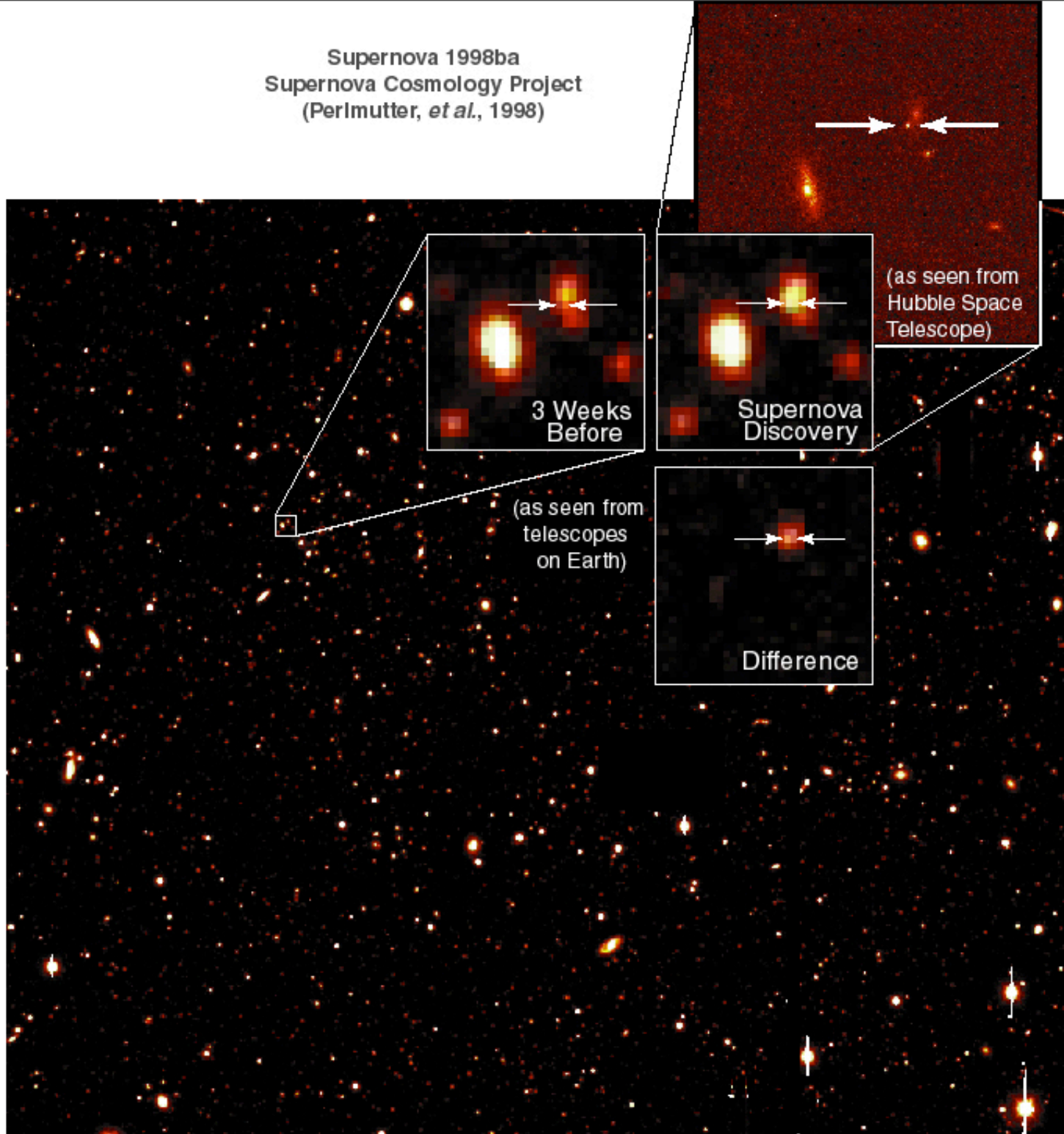
H_0



Other hints as well:

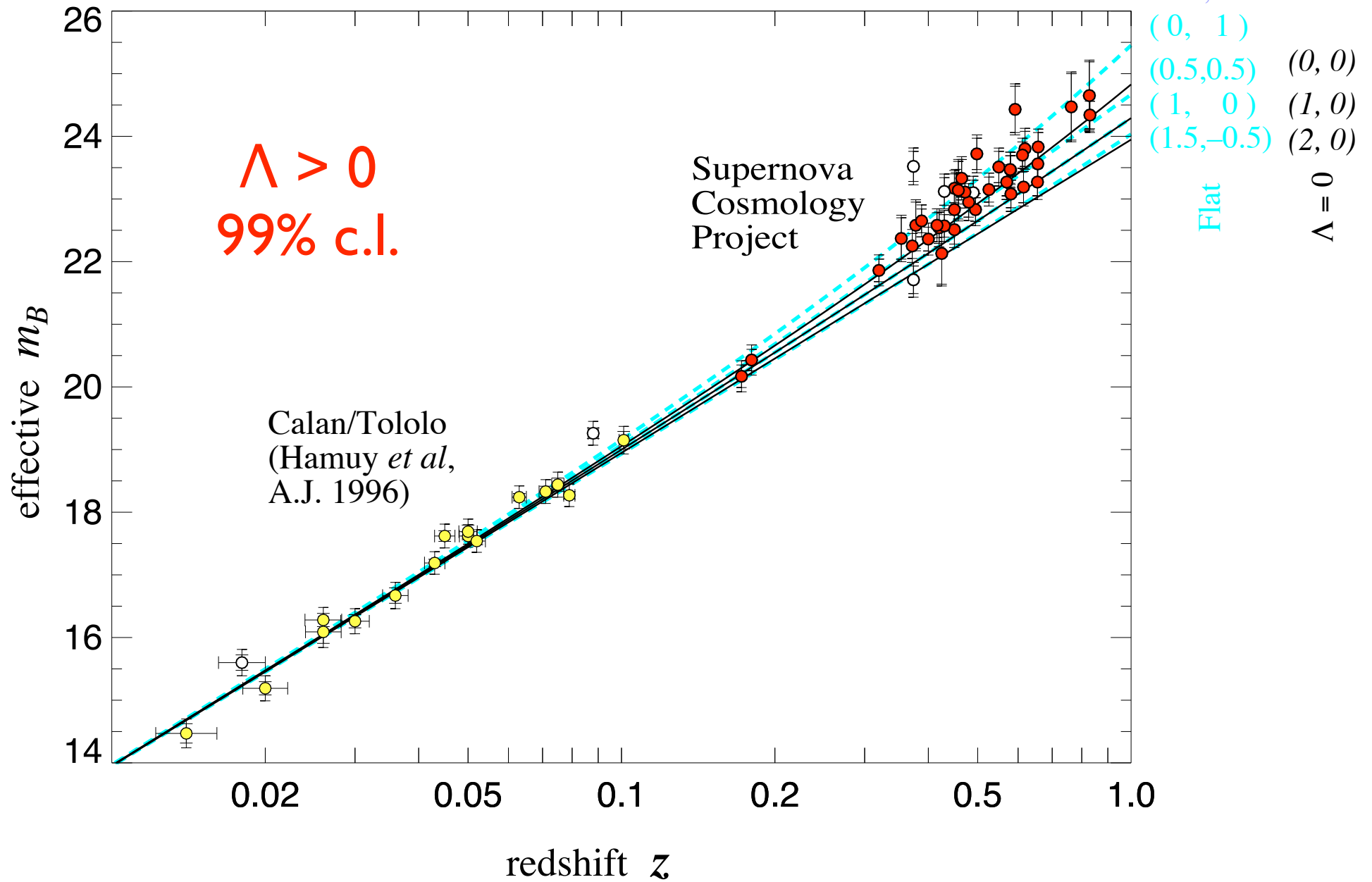
“...these number count data favor a flat, low-density $\Omega_m \approx 0.2$ universe with a nonzero **cosmological constant**.” (Yoshi & Peterson 1995, ApJ, 444, 15)

Several lines of evidence suggest that
the expansion of the universe is
indeed accelerating!

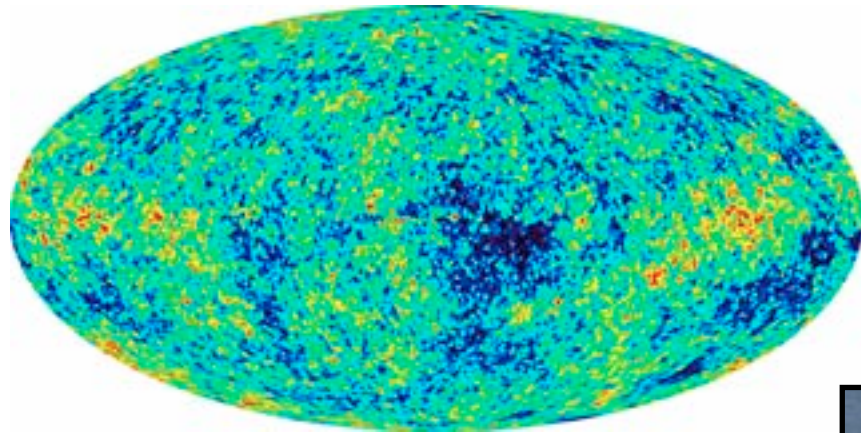


Supernova Cosmology Project

Perlmutter *et al.* (1998)



In flat universe: $\Omega_M = 0.28 [\pm 0.085 \text{ statistical}] [\pm 0.05 \text{ systematic}]$

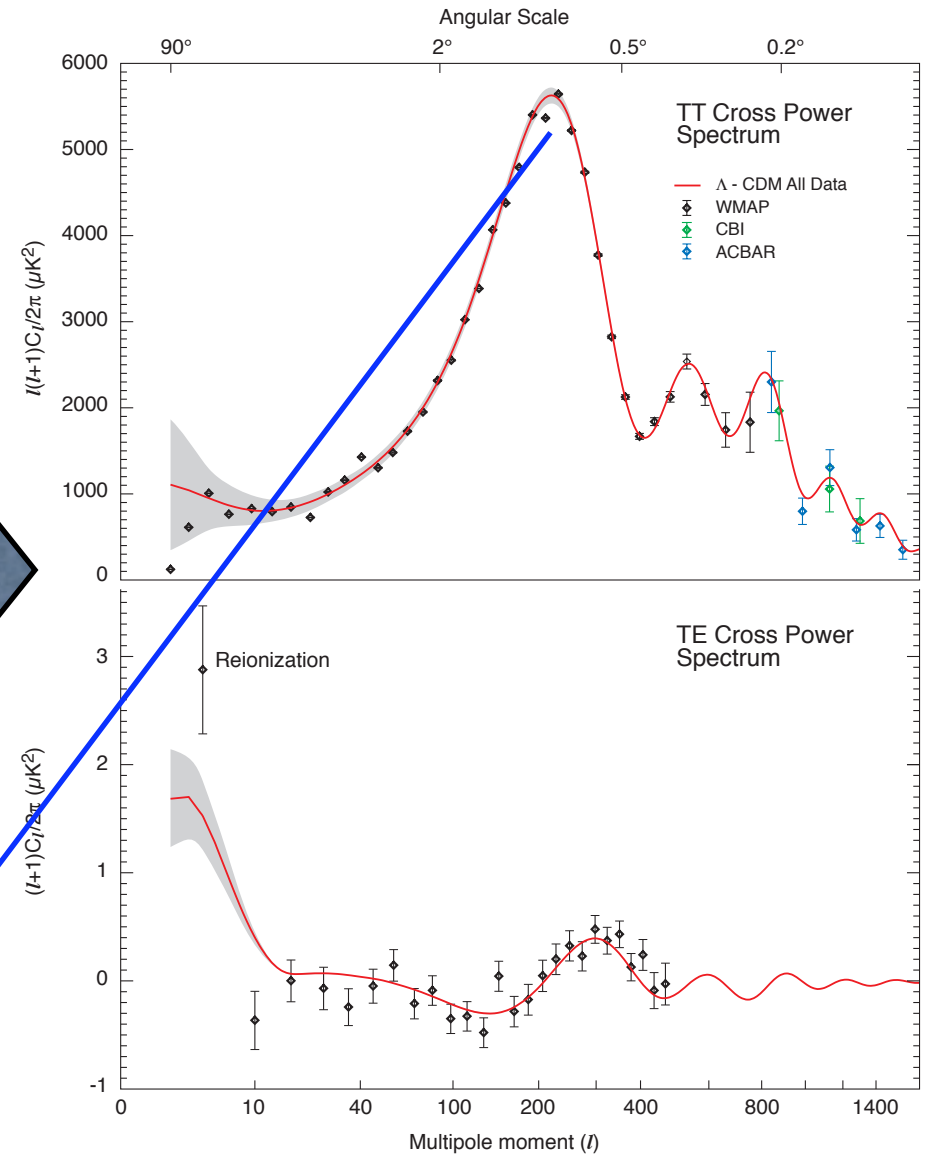
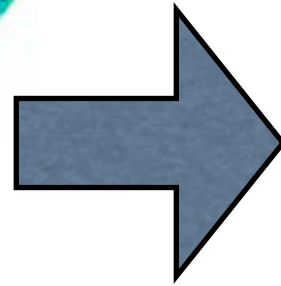


WMAP Cosmic Background Radiation

Position of first peak implies

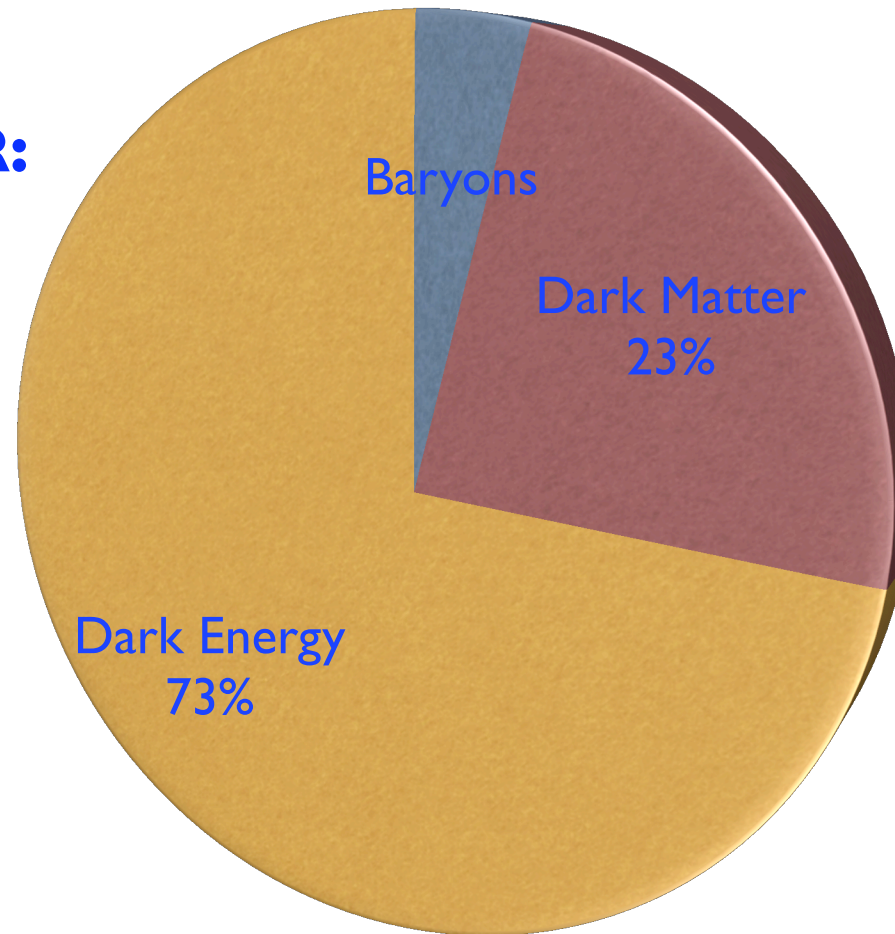
$$\Omega_m + \Omega_\Lambda \approx 1$$

$$\Omega_m < 1, \text{ so again, } \Lambda > 0$$



Cosmological Mass-Energy Budget

THE ANSWER:



Hot Big Bang
+
Dark Energy
+
Dark Matter

is now called

“ Λ CDM”

“The more important fundamental laws and facts of physical science have all been discovered, and these are now so firmly established that the possibility of their ever being supplanted in consequence of new discoveries is exceedingly remote.”
- L. Krauss (2004)

- A. Michelson (1894)

“Cosmologists are often wrong, but never in doubt”
- Lev Landau

Things we know **for sure** in cosmology:

1990:

$$\Omega_m = 1.00$$

$$\Omega_\Lambda = 0.00$$

$$\Omega_b h^2 = 0.0125$$

$$H_o = 50 \text{ km/s/Mpc}$$

Dark Matter = **C**old **D**ark **M**atter

“Cosmologists are often wrong, but never in doubt”
- Lev Landau

Things we know **for sure** in cosmology:

2003:

$$\Omega_m = 0.27$$

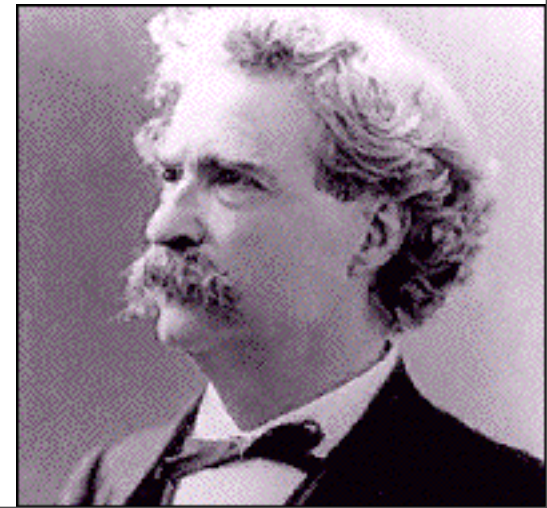
$$\Omega_\Lambda = 0.73$$

$$\Omega_b h^2 = 0.0224$$

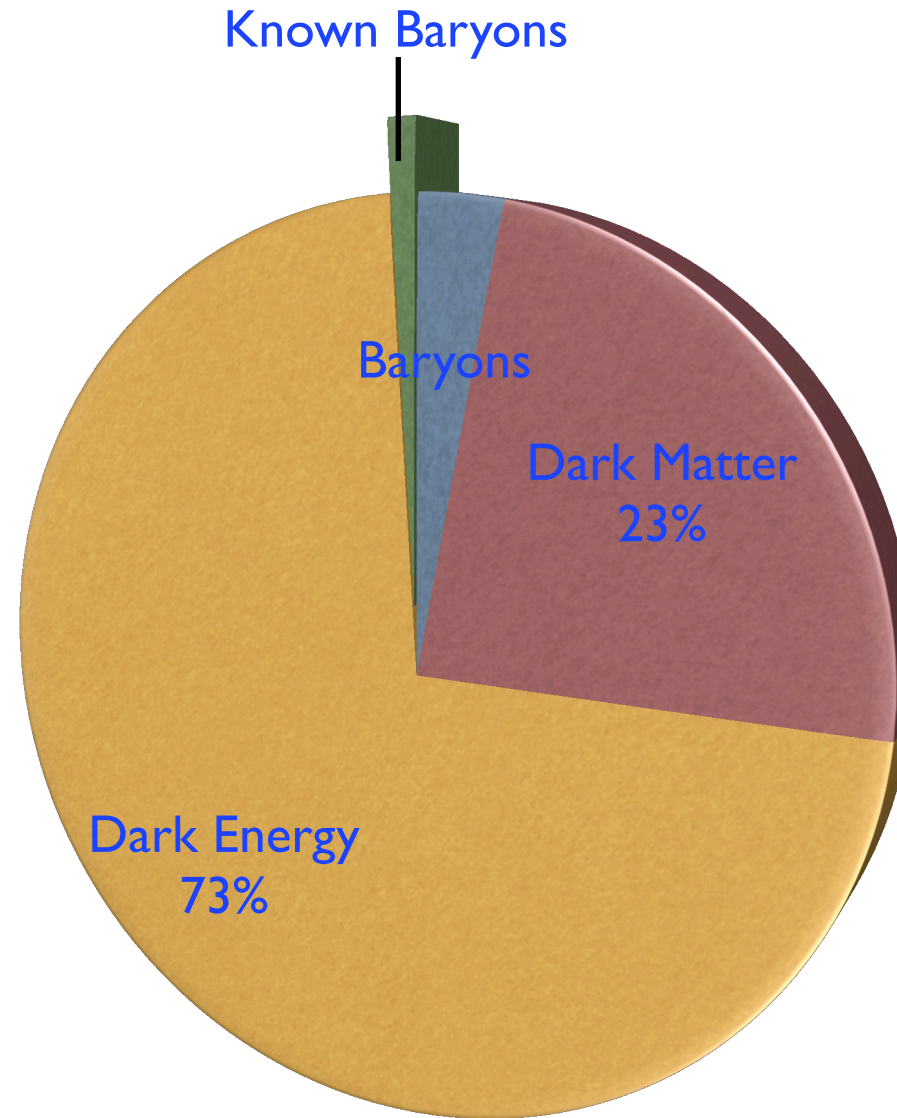
$$H_o = 72 \text{ km/s/Mpc}$$

Dark Matter = **C**old **D**ark **M**atter
... or maybe **W**arm **D**ark **M**atter
or **S**elf-Interacting **D**ark **M**atter

What did I say?



We have direct knowledge of < 1% of this stuff.



CDMS, LHC, & GLAST should all see something soon

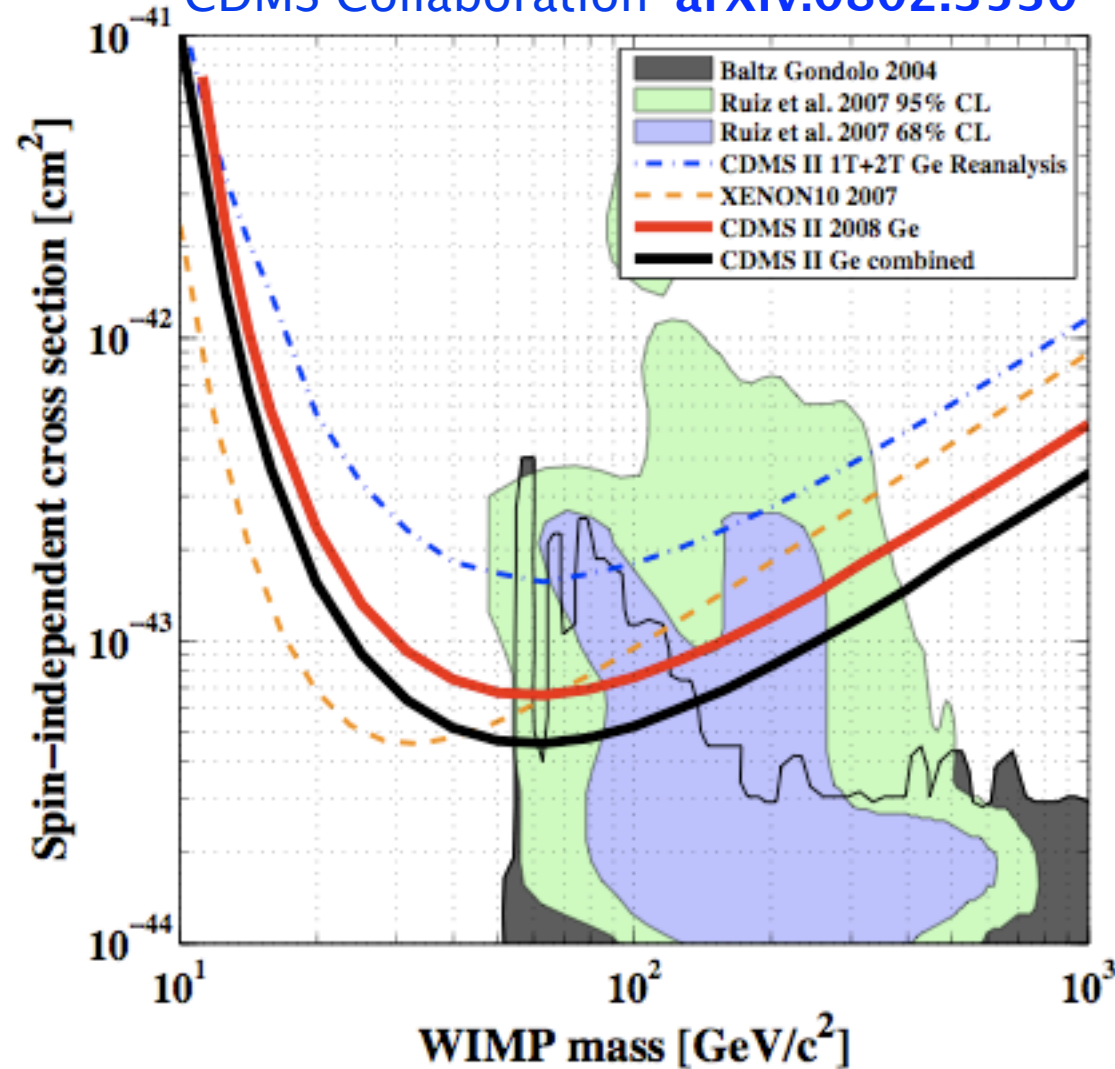


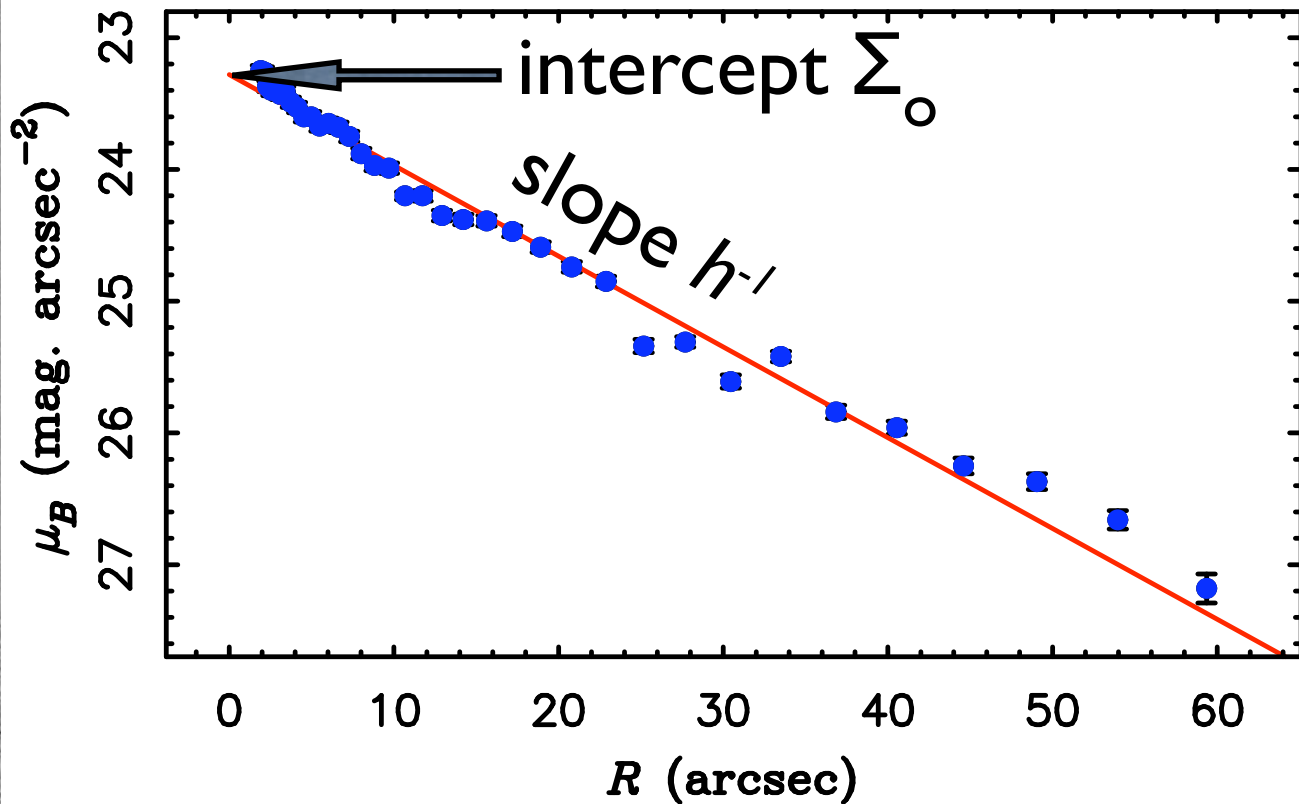
FIG. 4: Spin-independent WIMP-nucleon cross section upper limits (90% C.L.) versus WIMP mass. The upper curve (dash-dot) is the result of a re-analysis [17] of our previously published data. The upper solid line represents the limit derived from our new data set. The combined CDMS limit

On Galaxy Scales...

- Measure rotation velocity; find
- Properties depend systematically on
 - Total Baryonic Mass
 - Baryon Distribution
 - Acceleration



High Surface Brightness (HSB)



$$\Sigma(R) = \Sigma_o e^{-R/h}$$

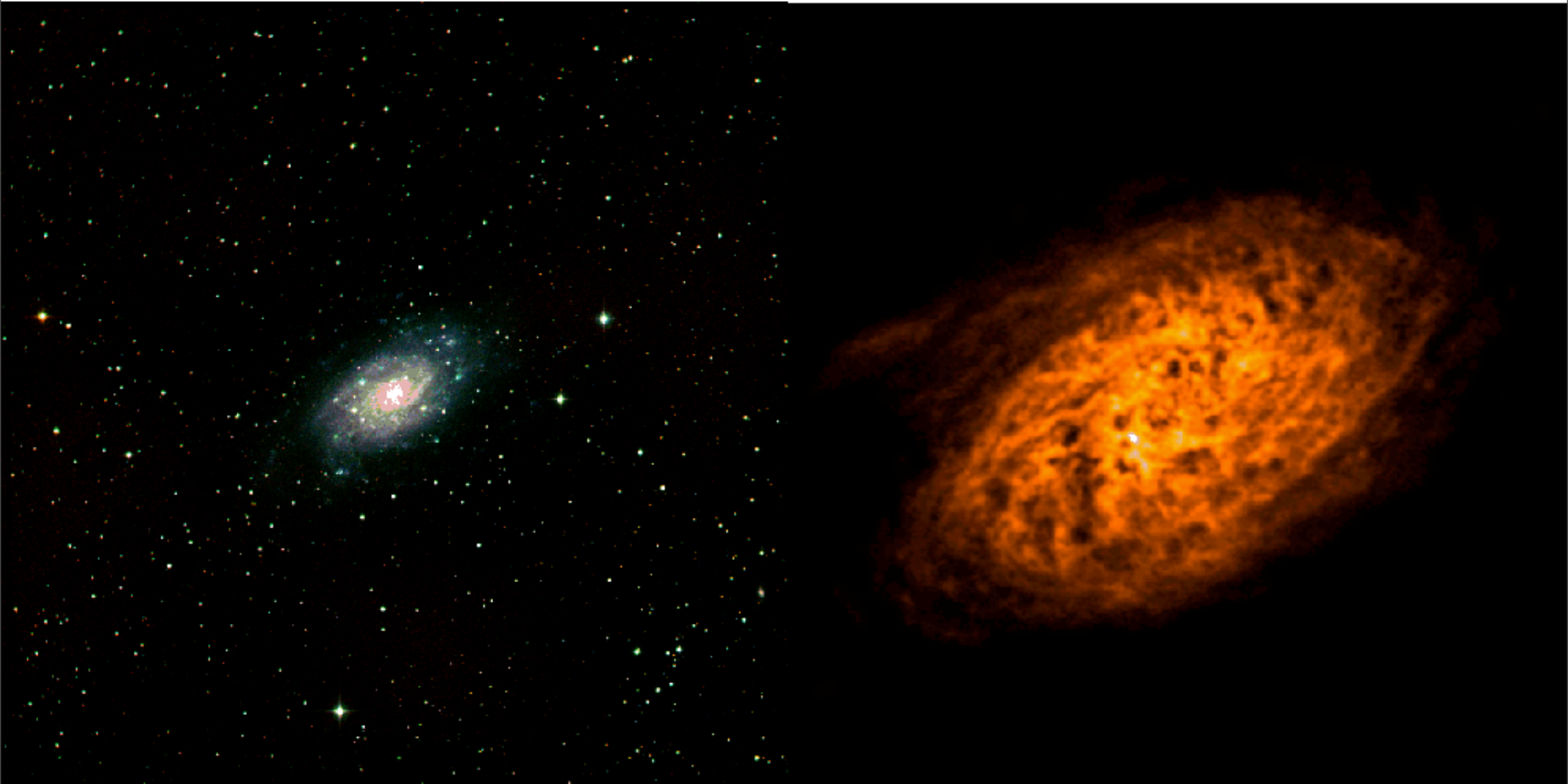
Azimuthally averaged light distribution
typically exponential for spiral disks.

Low Surface Brightness (LSB)

NGC 2403

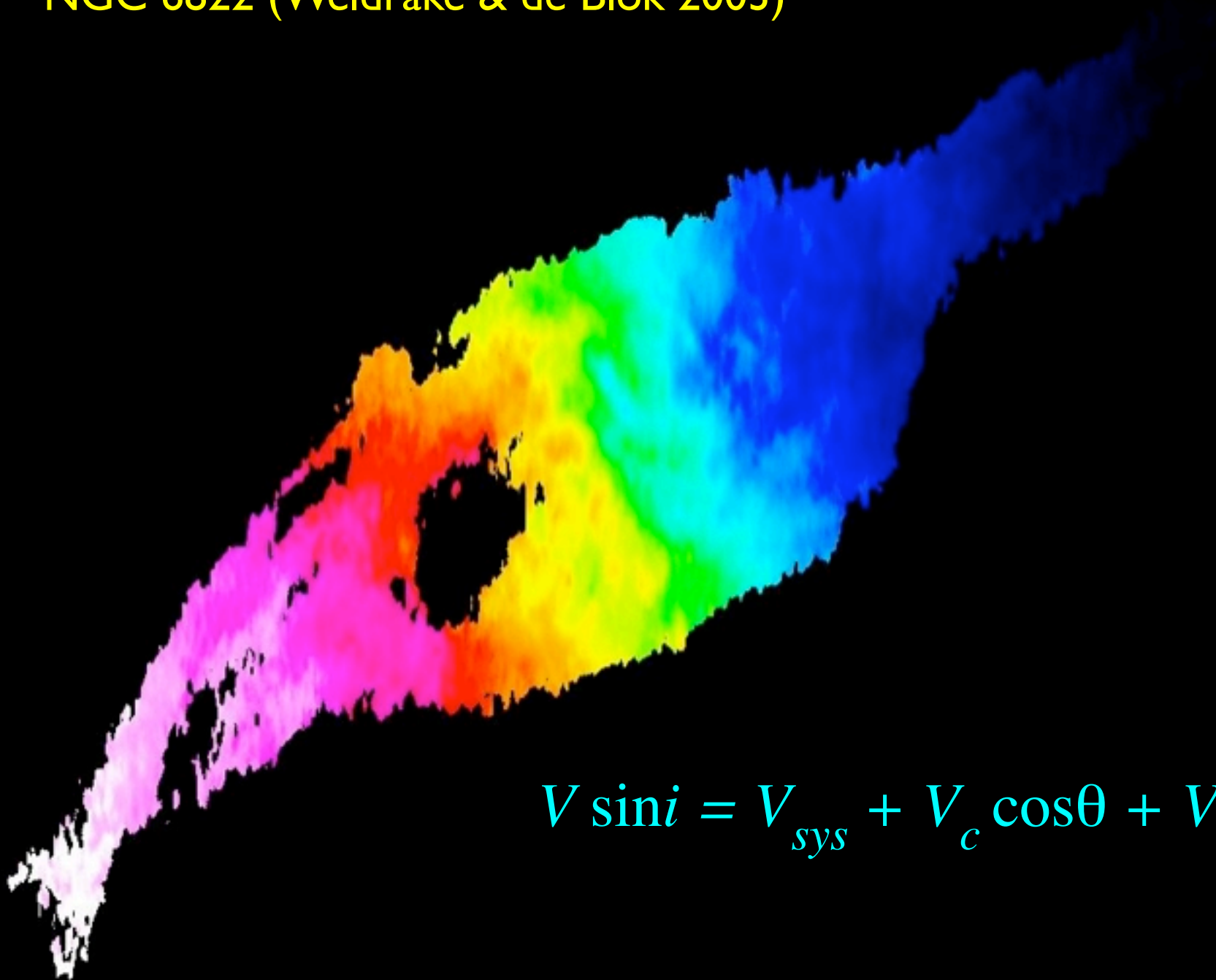
Stars

H I gas



Fraternali, Oosterloo, Sancisi, & van Moorsel 2001, ApJ, 562, L47

NGC 6822 (Weldrake & de Blok 2003)

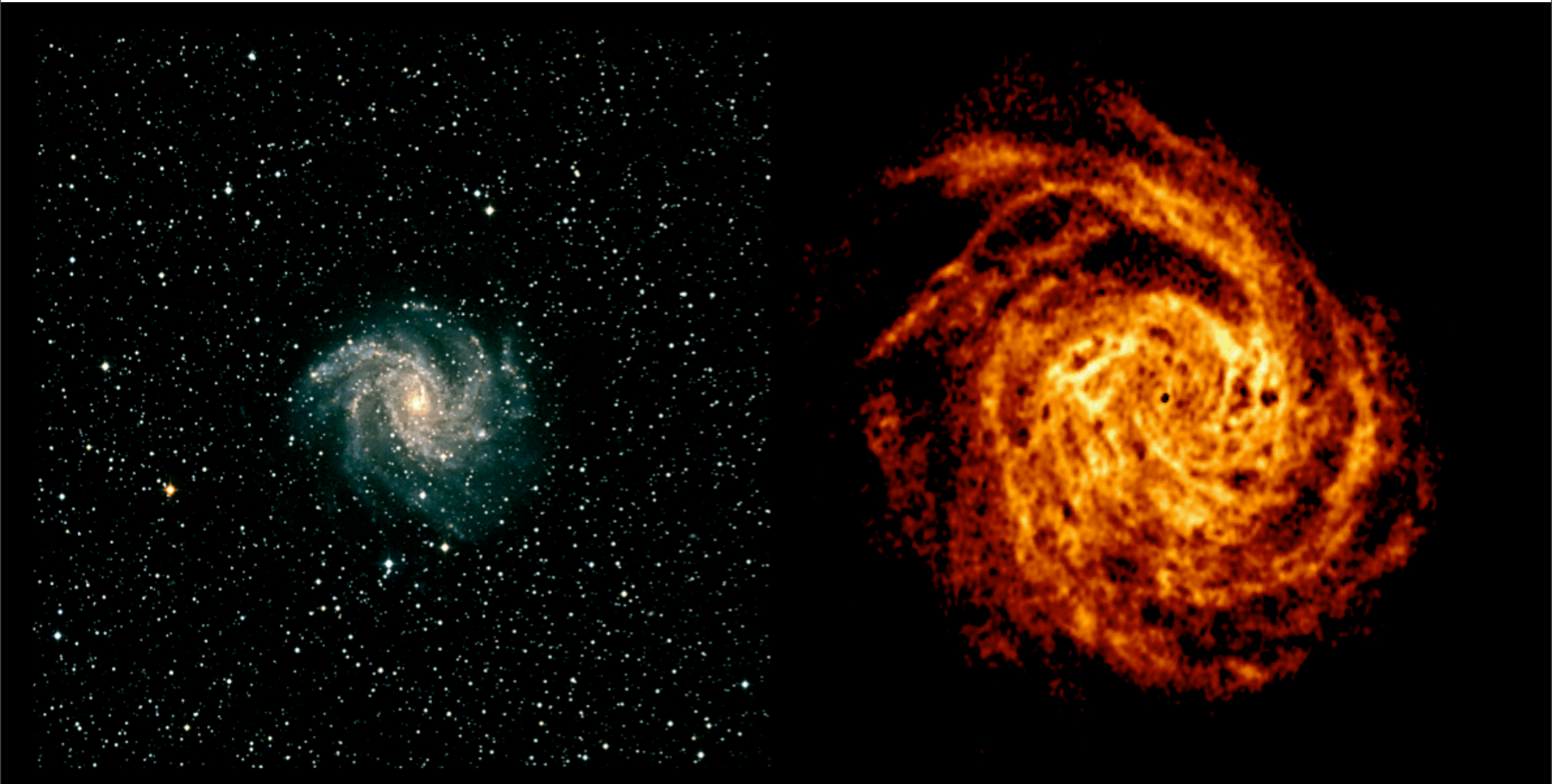


$$V \sin i = V_{sys} + V_c \cos \theta + V_r \sin \theta$$

NGC 6946

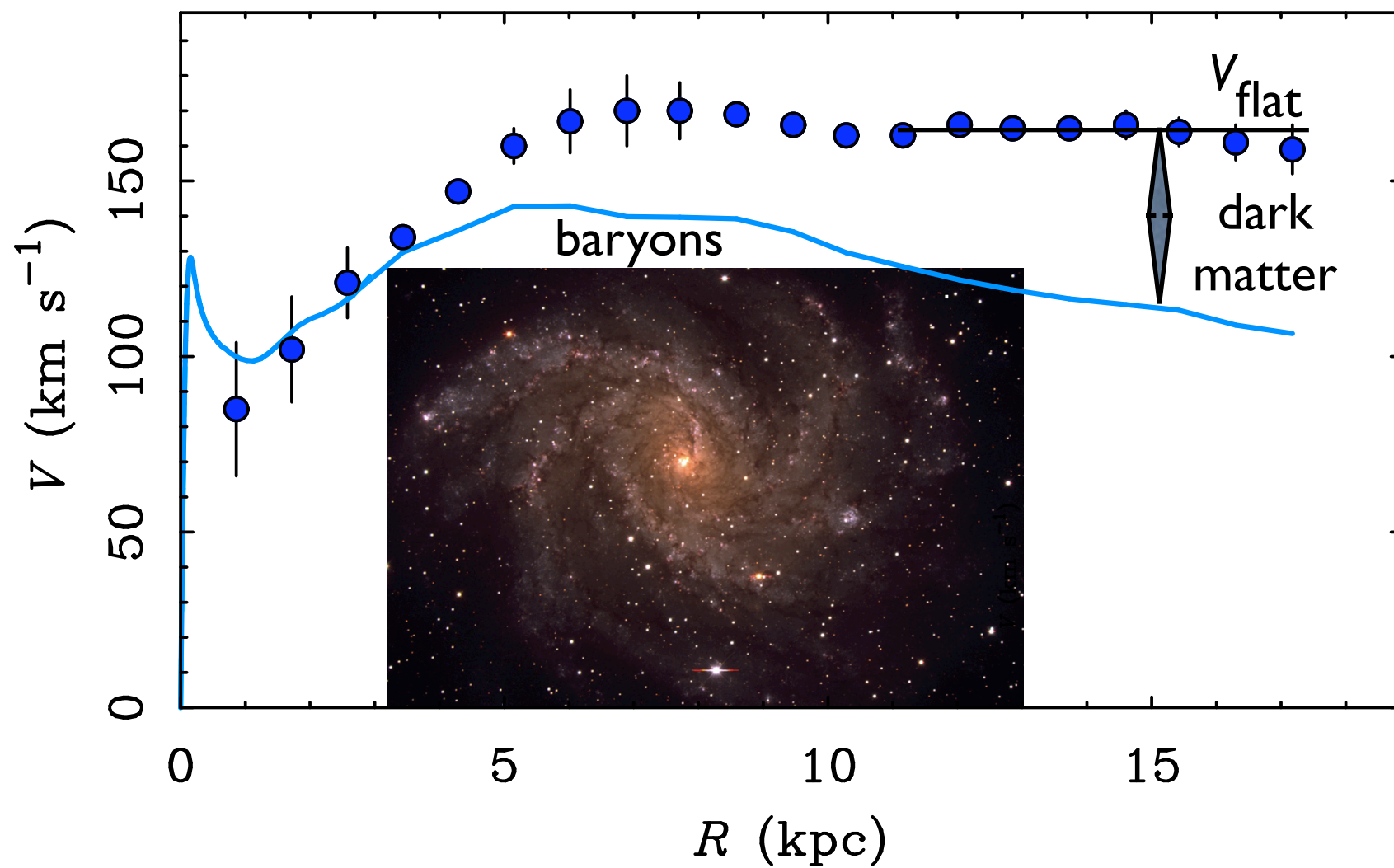
Stars

H I gas



Boomsma 2005

NGC 6946



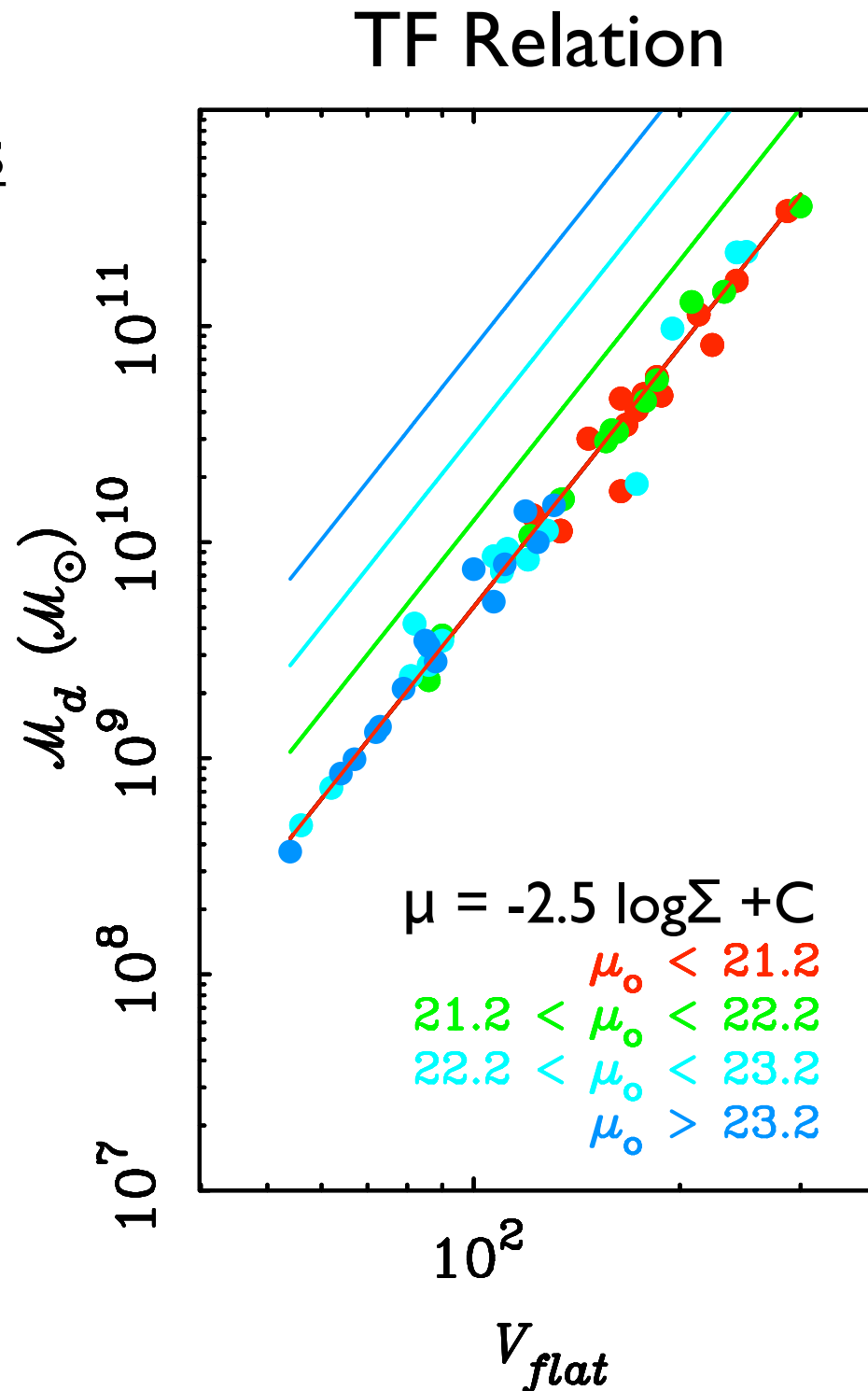
Newton says

$$V^2 = GM/R.$$

Equivalently,

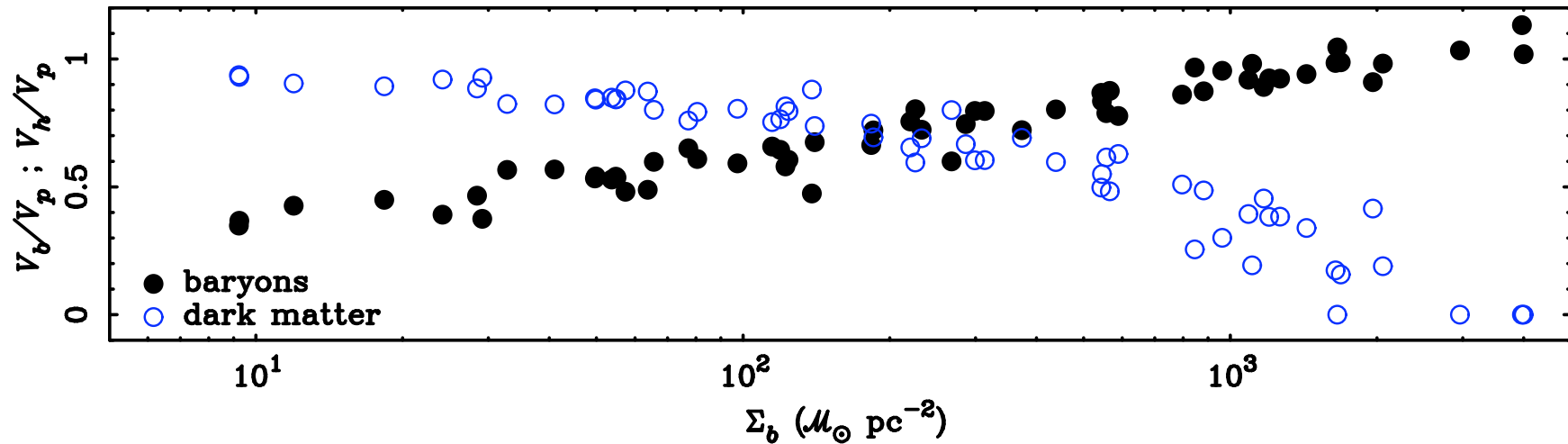
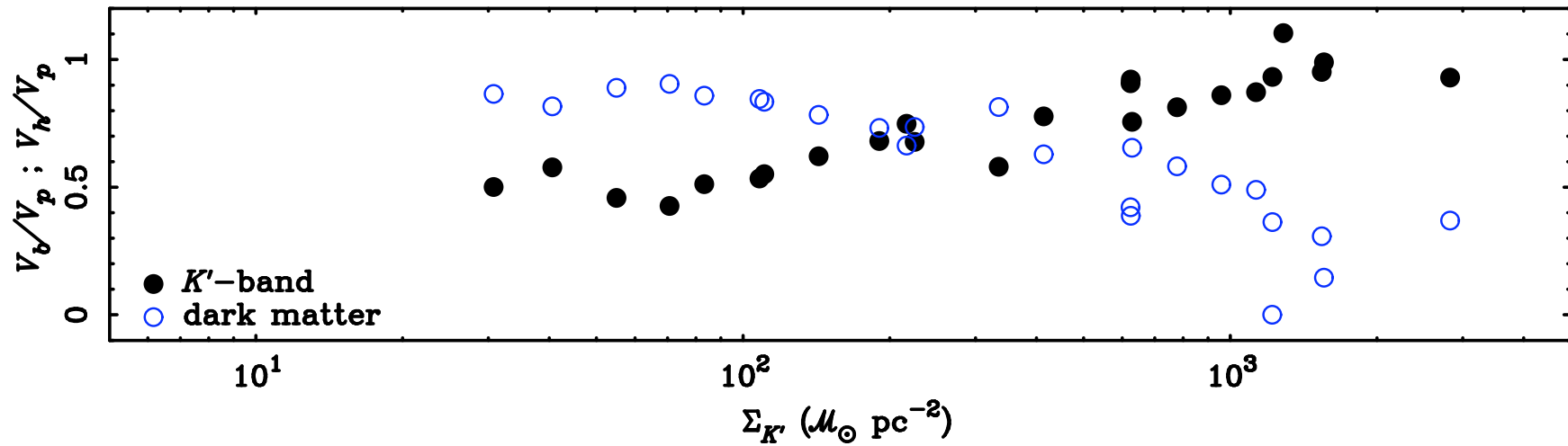
$$\Sigma = M/R^2$$

$$V^4 = G^2 M \Sigma$$

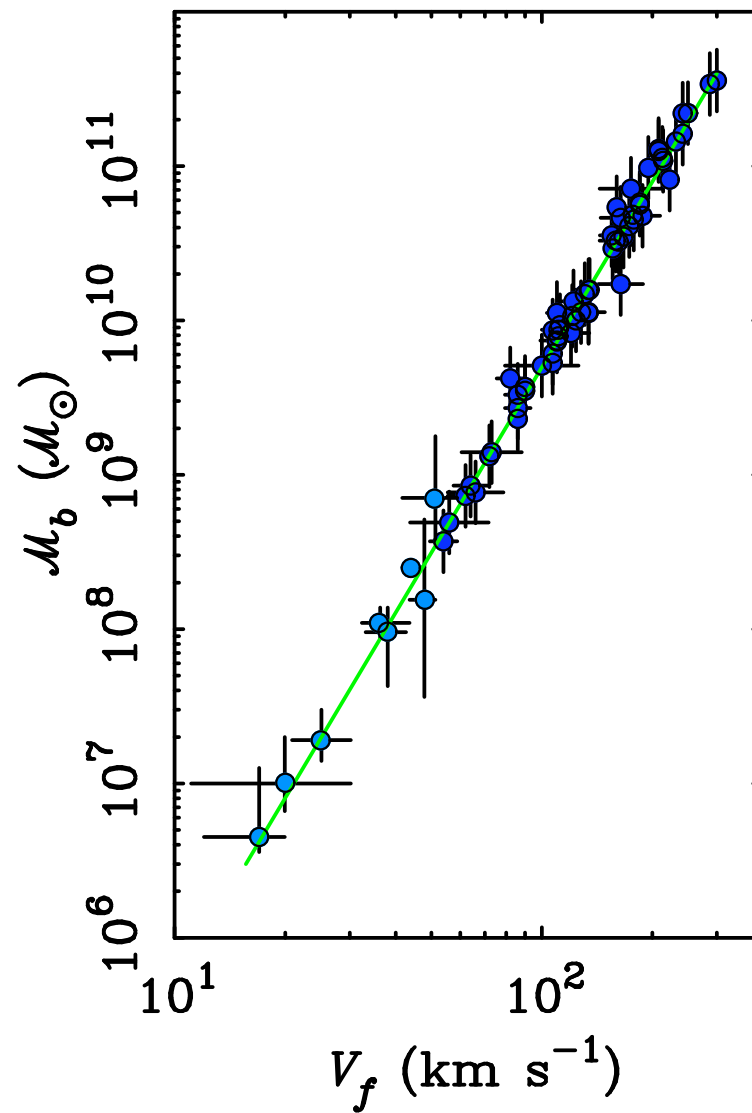
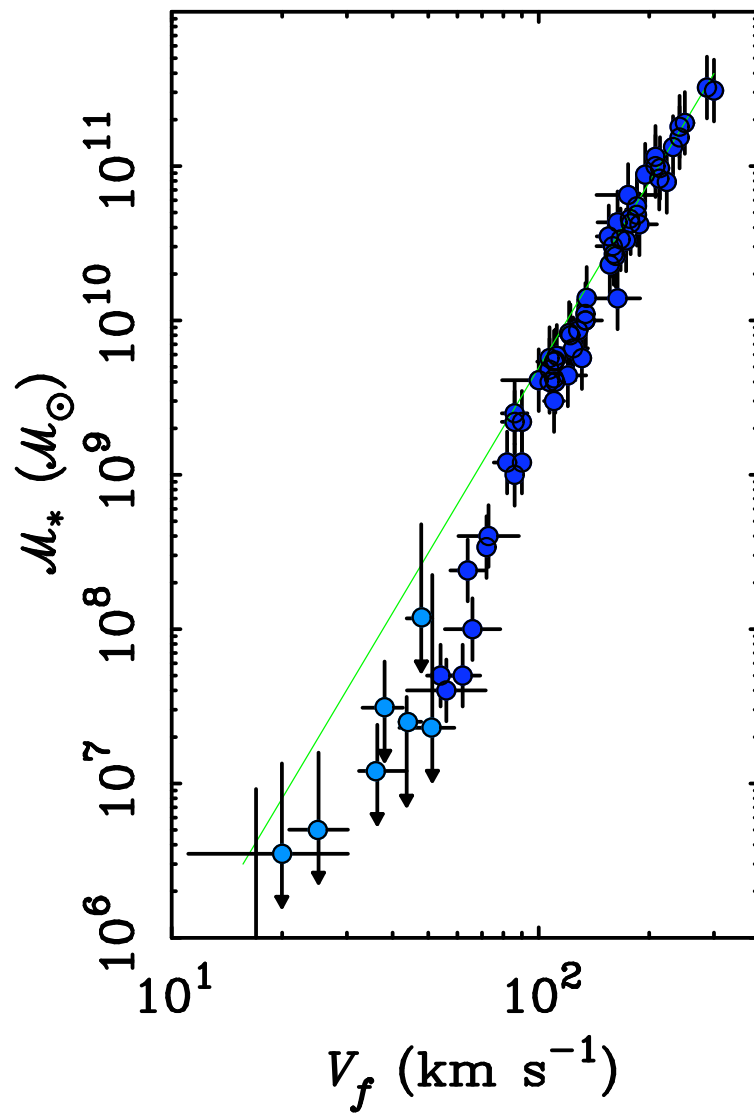


Therefore
Different Σ
should mean
different TF
normalization.

Requires fine balance between dark & baryonic mass

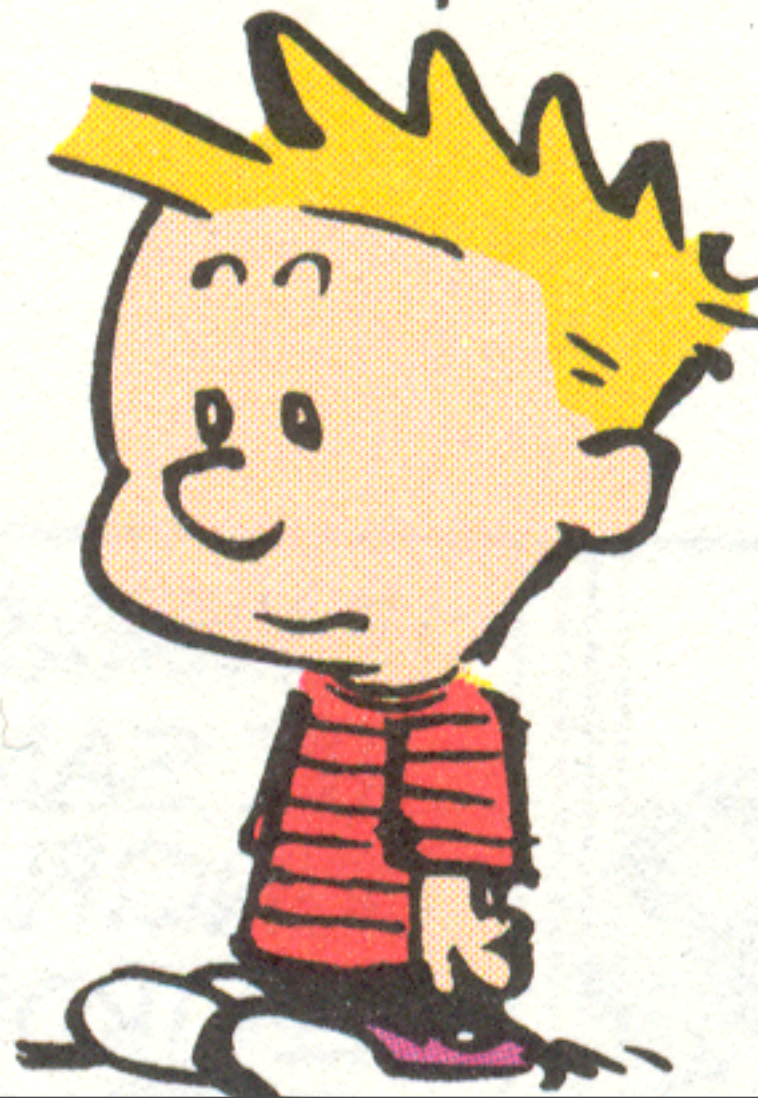
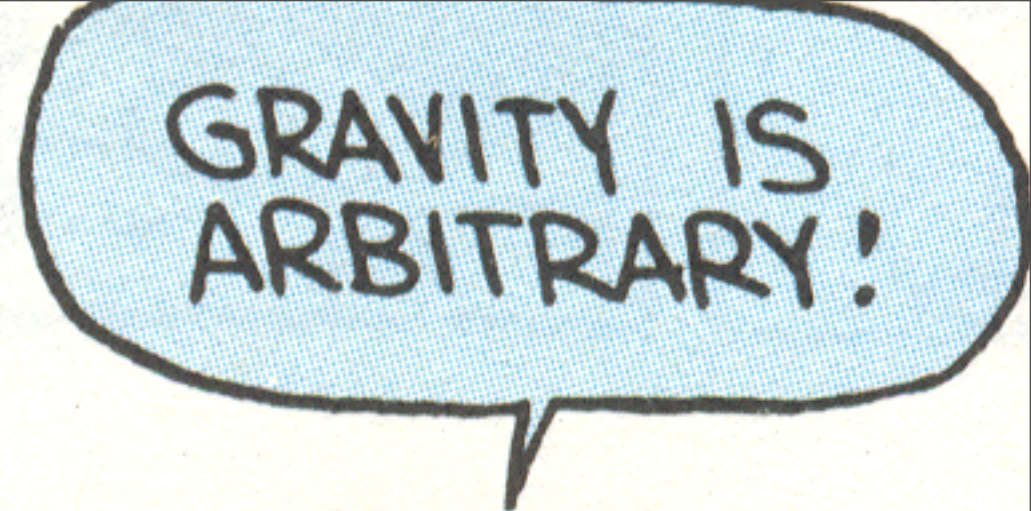


Baryonic Tully-Fisher Relation



Depends only on ordinary mass

One begins to worry that



MOND

MOdified Newtonian Dynamics

introduced by Moti Milgrom in 1983

instead of dark matter, suppose the force law changes such that

$$\text{for } a \gg a_o, \quad a \Rightarrow g_N$$

$$\text{for } a \ll a_o, \quad a \Rightarrow \sqrt{(g_N a_o)}$$

where

$$g_N = GM/R^2$$

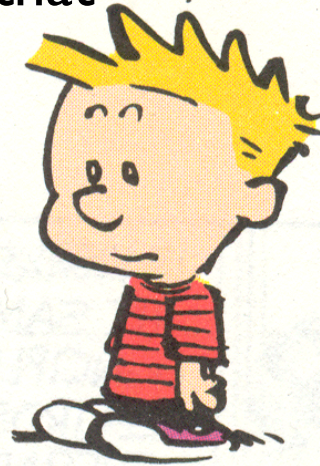
is the usual Newtonian acceleration.

More generally, these limits are connected by a smooth interpolation fcn $\mu(a/a_o)$ so that

$$\mu(a/a_o) a = g_N.$$

MOND can be interpreted as a modification of either **inertia** ($F = ma$) or **gravity** (the Poisson eqn).

GRAVITY IS
ARBITRARY.



Milgrom 1983

No. 2, 1983

MODIFICATION OF NEWTONIAN DYNAMICS

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A major step in understanding ellipticals can be made if we can identify them, at least approximately, with idealized structures such as the FRCL spheres discussed above. I have also studied isotropic and nonisotropic isothermal spheres, in the modified dynamics, as such possible structures. I found that they have properties which resemble those of ellipticals and galactic bulges. I discuss them in Milgrom (1983b).

VIII. PREDICTIONS

The main predictions concerning galaxies are as follows.

1. Velocity curves calculated with the modified dynamics on the basis of the observed mass in galaxies should agree with the observed curves. Elliptical and SO galaxies may be the best for this purpose since (a) practically no uncertainty due to obscuration is involved and (b) there is not much uncertainty due to the possible presence of molecular hydrogen.

2. The relation between the asymptotic velocity (V_∞) and the mass of the galaxy (M) ($V_\infty^4 = MG a_0$) is an absolute one.

3. Analysis of the π -dynamics in disk galaxies using the modified dynamics should yield surface densities which agree with the observed ones. In principle, the same analysis using conventional dynamics should yield a discrepancy which increases with radius in a predictable manner.

4. Effects of the modified dynamics are predicted to be particularly strong in dwarf elliptical galaxies (for review of properties see, e.g., Hodge 1971 and Zinn 1980). For example, those dwarfs believed to be bound to our Galaxy would have internal accelerations typically of order $a_{in} \sim a_0/30$. Their (modified) acceleration, g , in the field of the Galaxy is larger than the internal ones but still much smaller than a_0 , $g \approx (8 \text{ kpc}/d) a_0$, based on a value of $V_\infty = 220 \text{ km s}^{-1}$ for the Galaxy, and where d is the distance from the dwarf galaxy to the center of the Milky Way ($d \sim 70\text{--}220 \text{ kpc}$). Whichever way the external acceleration turns out to affect the internal dynamics (see the discussion at the end of § II, the section on small groups in Paper III, and Paper I), we predict that when velocity dispersion data is available for the dwarfs, a large mass discrepancy will result when the conventional dynamics is used to determine the masses. The dynamically determined mass is predicted to be larger by a factor of order 10 or more than that which can be accounted for by stars. In case the internal dynamics is determined by the external acceleration, we predict this factor to increase with d and be of order $(d/8 \text{ kpc})$ (as long as $a_{in} \ll g$, $h_{50} = 1$).

Prediction 1 is a very general one. It is worthwhile listing some of its consequences as separate predictions, numbered 5–7 below (note that, in fact, even prediction 2 is already contained in prediction 1).

5. Measuring local M/L values in disk galaxies (assuming conventional dynamics) should give the following results: In regions of the galaxy where $V^2/r \gg a_0$ the local M/L values should show no indication of hidden mass. At a certain transition radius, local M/L should start to increase rapidly. The transition radius should occur where $V^2/r \approx a_0$. This was the first of the two stages. (a) The first stage requires an absolute calibration of M/L as we are concerned only with variations of this quantity; (b) Effects of the modified dynamics manifest themselves more clearly in local mass determinations than in integrated masses and (c) in many cases (b) requires information on local behavior in the disk only while the spheroid can be neglected. This makes the determination of mass from velocity more certain.

6. Disk galaxies with low surface brightness provide particularly strong tests (a study of a sample of such galaxies is described by Strom 1982 and by Romanishin *et al.* 1982). As low surface brightness means small accelerations, the effects of the modification should be more noticeable in such galaxies. We predict, for example, that the proportionality factor in the $M \propto V_\infty^4$ relation for these galaxies is the same as for the high surface density galaxies. In contrast, if one wants to obtain a relation $M \propto V_\infty^2$ in the conventional dynamics with modification as in prediction 1, one has to let the relation be $M \propto \Sigma^{-1} V_\infty^4$ (see, for example, Aaronson, Huchra, and Mould 1979), where Σ is the average surface brightness. This implies that low surface density galaxies, of a given total mass, have mass higher than predicted by the $M \propto V_\infty^4$ relation derived for normal surface density galaxies.

We also predict that the lower the average surface density of a galaxy is, the smaller is the transition radius. The predicted transition radius of the galaxy scales as $1/g$. Since the average surface density is very small we may have a galaxy in which $V^2/r < a_0$ everywhere, and analysis with conventional dynamics should yield local M/L values starting to increase from very small radii.

7. As the study of model rotation curves shows, we predict a correlation between the value of the average surface density (or brightness) of a galaxy and the steepness with which the rotational velocity rises to its asymptotic value (as measured, for example, by the radius at which $V = V_\infty/2$ in units of the scale length of the disk). Small surface densities imply slow rise of V .

IX. DISCUSSION

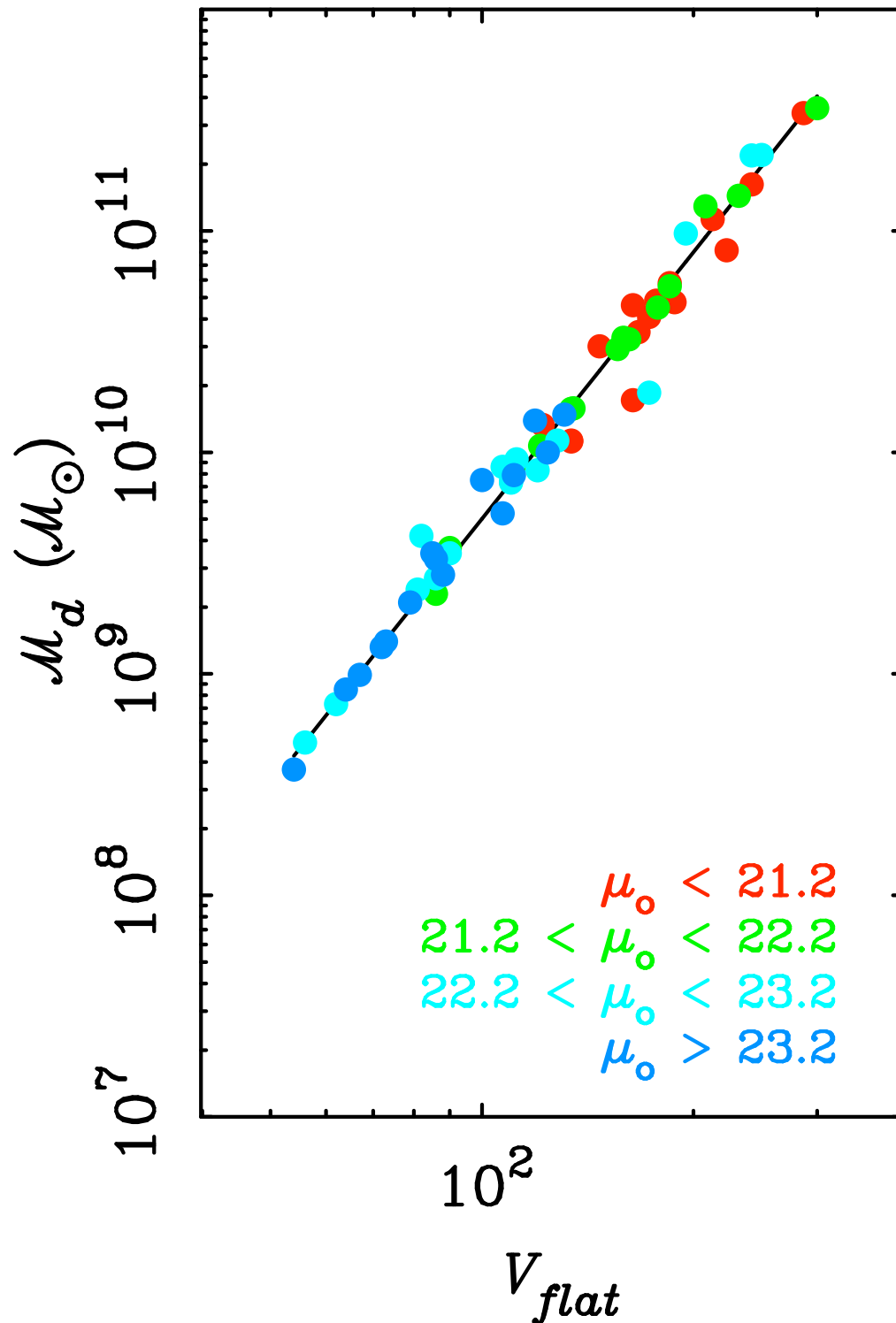
The main results of this paper can be summarized by the statement that the modified dynamics eliminates the need to assume hidden mass in galaxies. The effects in galaxies which I have considered, and which are commonly attributed to such hidden mass, are readily explained by the modification. More specifically:

MOND predictions

- The Tully-Fisher Relation
- Slope = 4
- Normalization = $V_\infty^4/(a_0 G)$
- Fundamental correlation between Disk Mass and V_{flat}
- No Dependence on Surface Brightness
- Dependence of conventional V/V_∞ on radius and surface brightness
- Rotation Curve Shapes
- Surface Density \sim Surface Brightness
- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios

“Disk Galaxies with low surface brightness provide particularly strong tests”

None of the following data existed in 1983. At that time, LSB galaxies were widely thought not to exist.



MOND predictions

- The Tully-Fisher Relation
 - ✓ • Slope = 4
 - ✓ • Normalization = $1/(a_0 G)$
 - ✓ • Fundamentally a relation between Disk Mass and V_{flat}
 - ✓ • No Dependence on Surface Brightness !
- Dependence of conventional M/L on radius and surface brightness
- Rotation Curve Shapes
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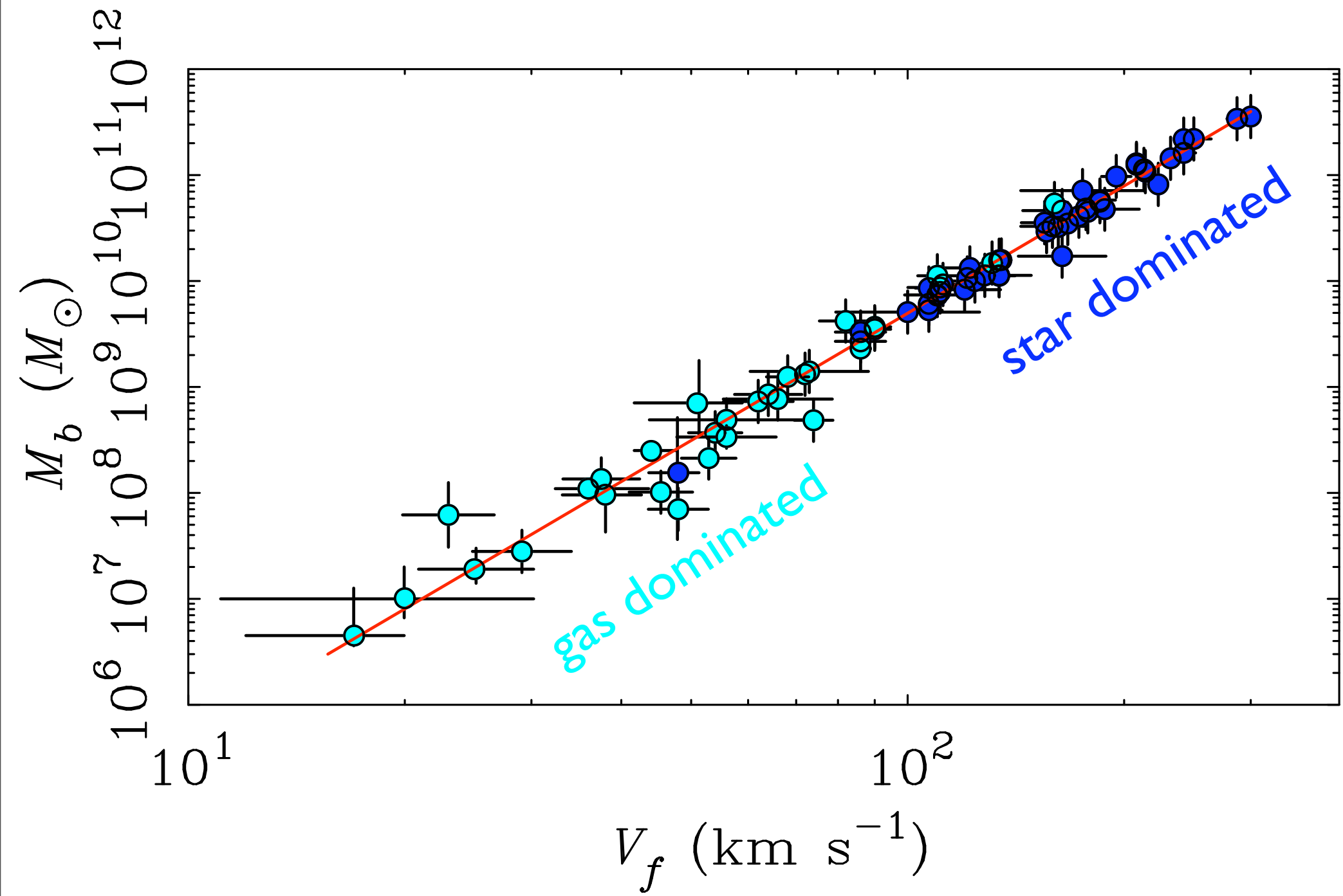
In MOND limit of low acceleration

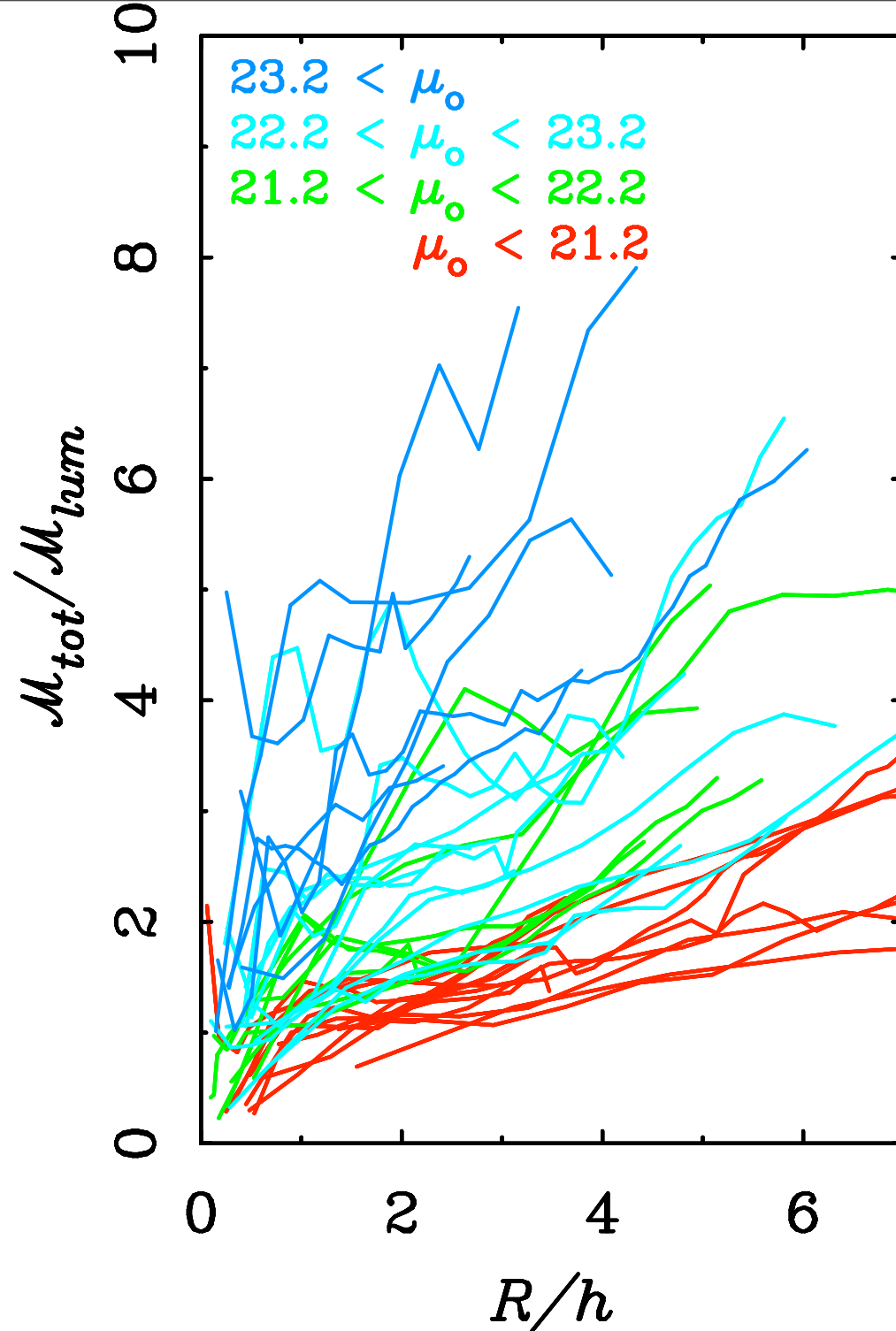
$$a = \sqrt{g_N a_0}$$

$$\frac{V^2}{\cancel{R}} = \sqrt{\frac{GM}{\cancel{R^2}}} a_0$$

$$V^4 = a_0 G M$$

observed TF!





MOND predictions

- The Tully-Fisher Relation



Slope = 4



Normalization = $1/(a_0 G)$



Fundamentally a relation between
Disk Mass and V_{flat}



No Dependence on Surface
Brightness



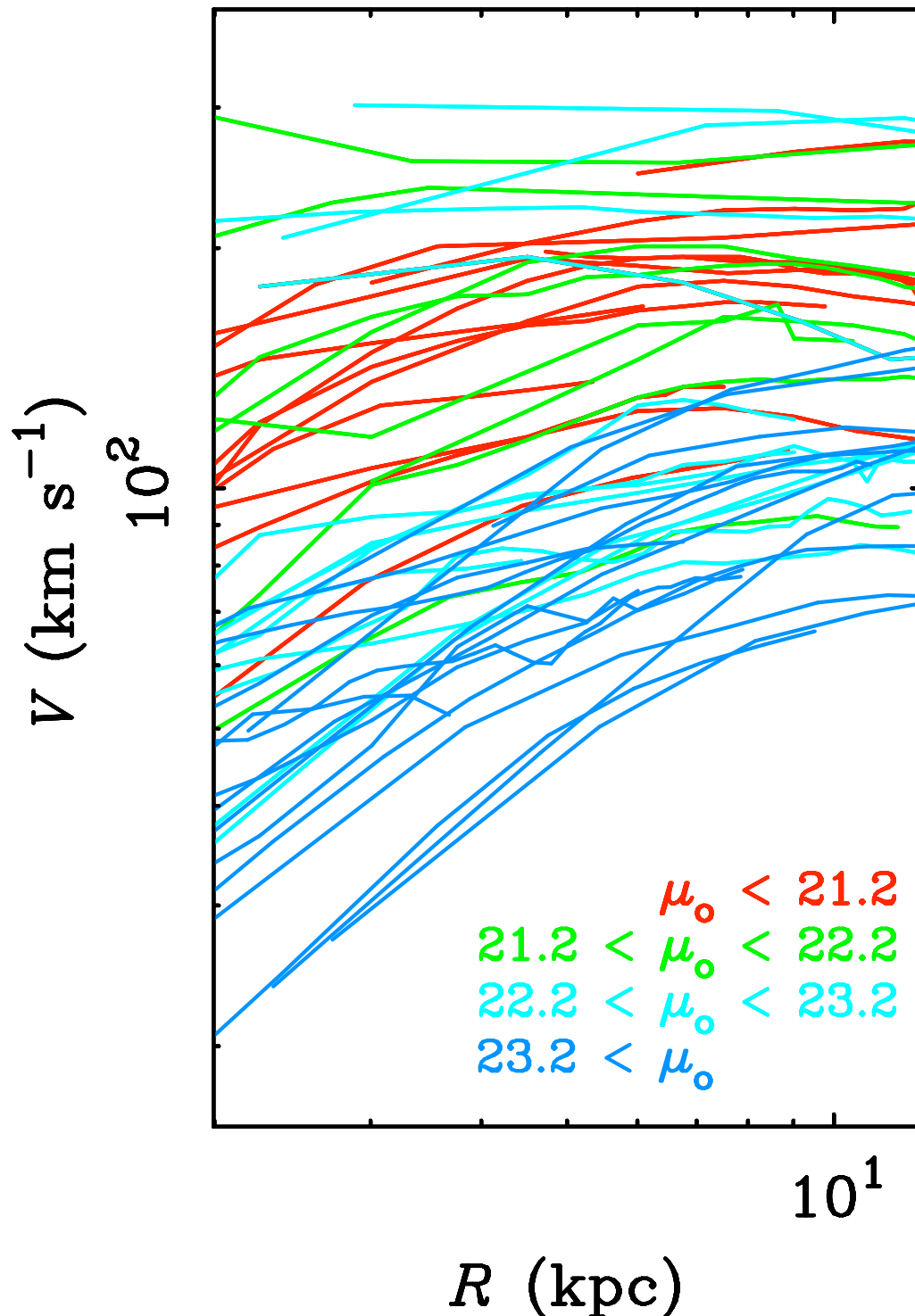
- Dependence of conventional M/L on
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Fundamentally a relation between Disk Mass and V_{flat}



No Dependence on Surface Brightness



Dependence of conventional M/L on radius and surface brightness



Rotation Curve Shapes

- Surface Density \sim Surface Brightness

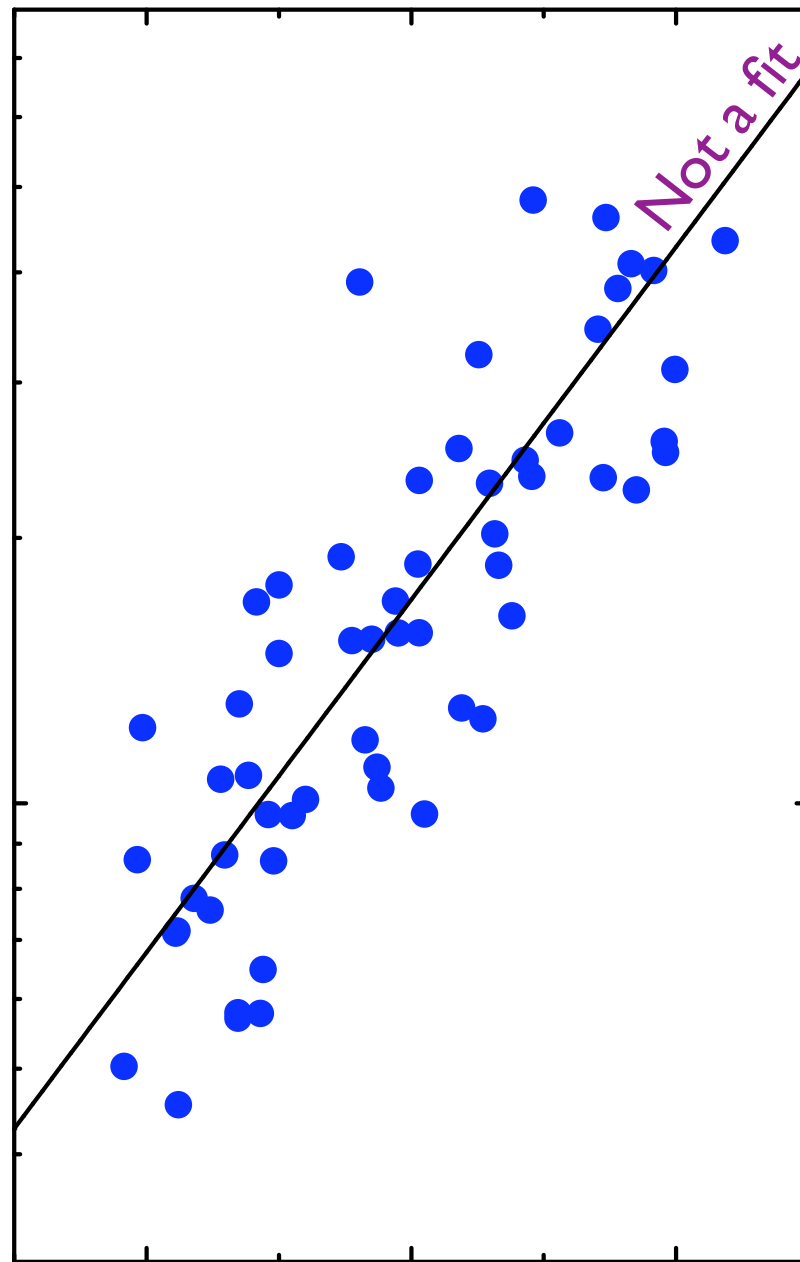
- Detailed Rotation Curve Fits

- Stellar Population Mass-to-Light Ratios

mass surface density ↑

$$\xi = V^2/(Gh)$$

5
1
0.5



24

22

20

μ_0
surface brightness →

MOND predictions

- The Tully-Fisher Relation



Slope = 4



Normalization = $1/(a_0 G)$



Fundamentally a relation between
Disk Mass and V_{flat}



No Dependence on Surface
Brightness



Dependence of conventional M/L on
radius and surface brightness



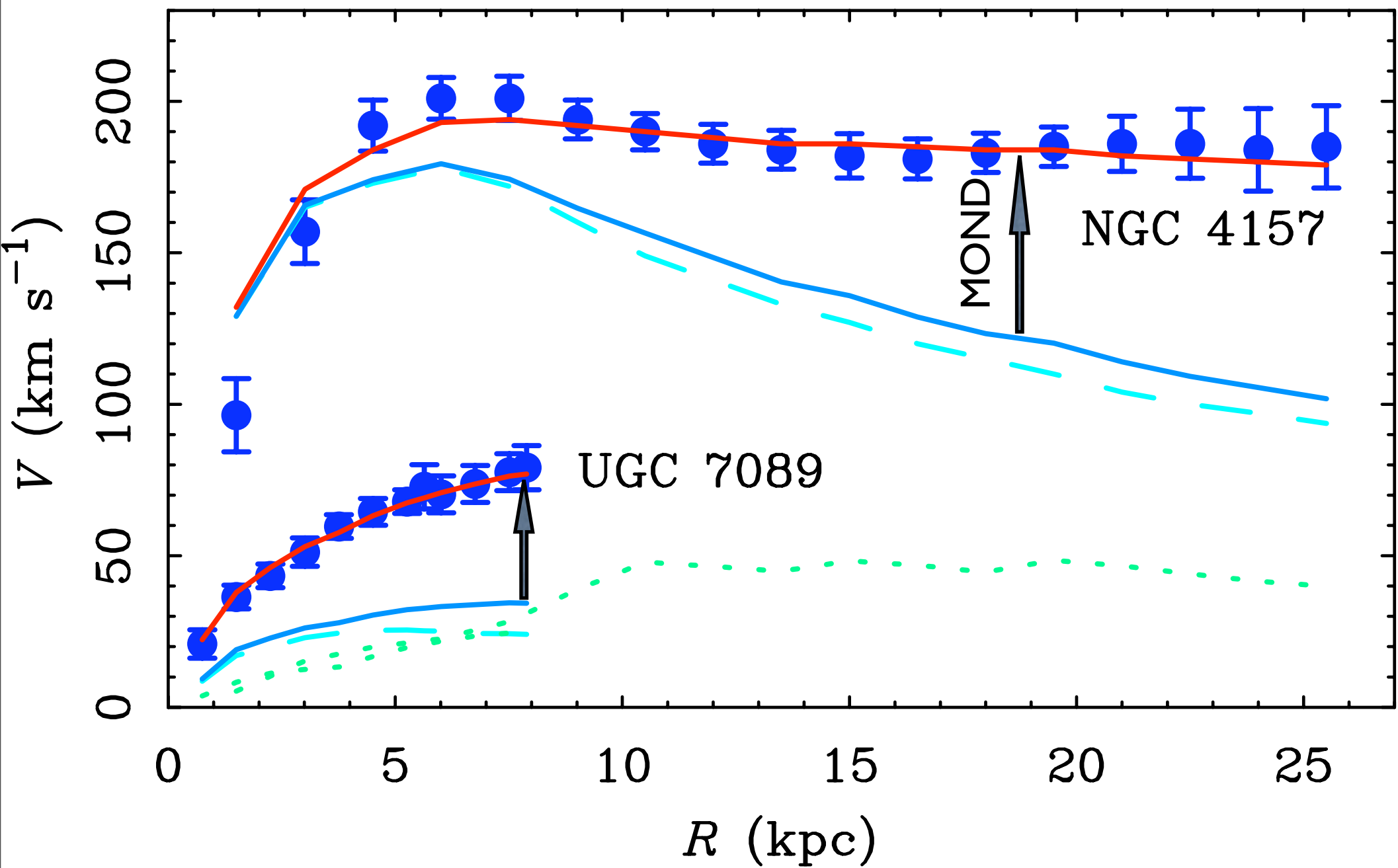
Rotation Curve Shapes

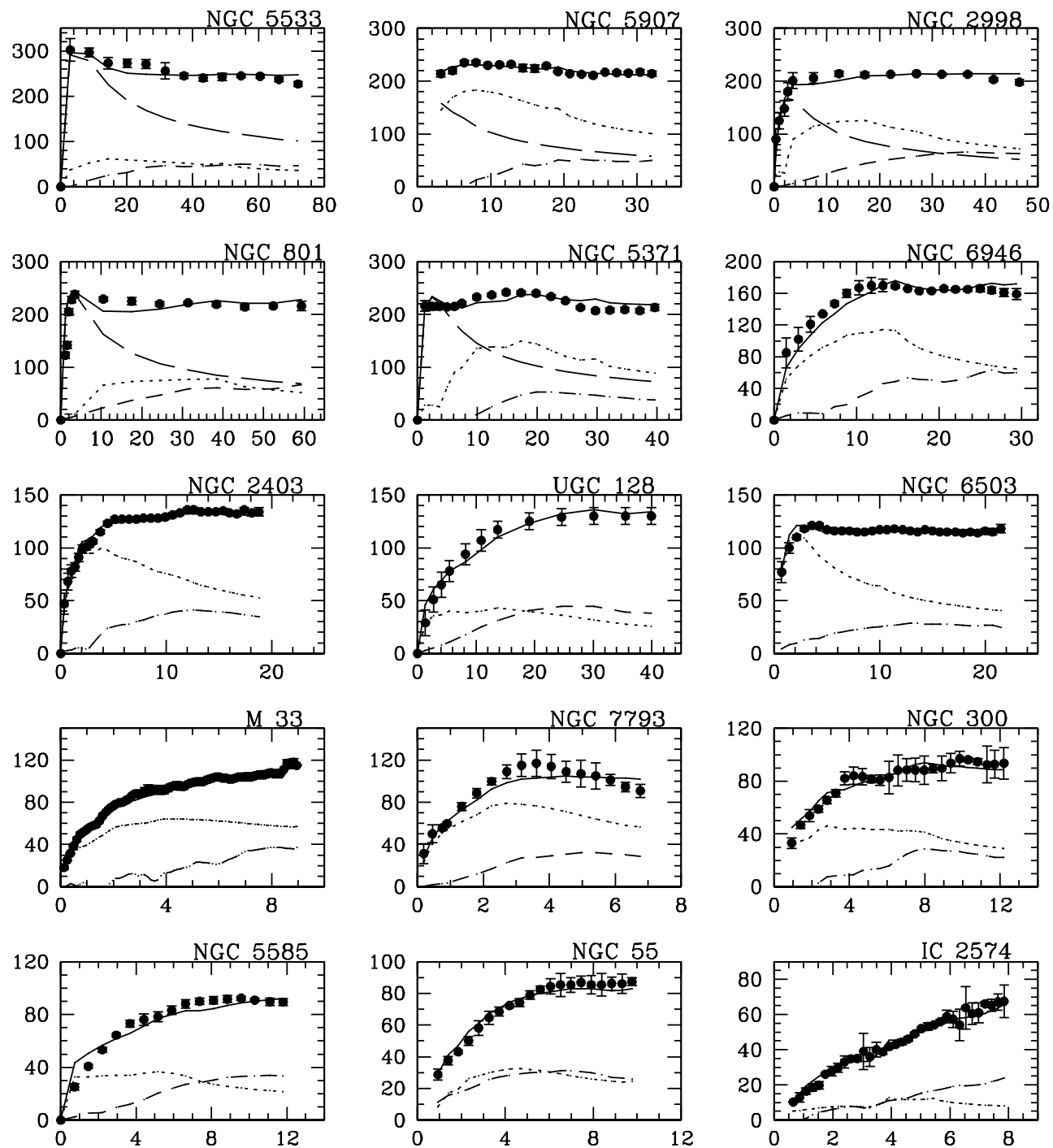


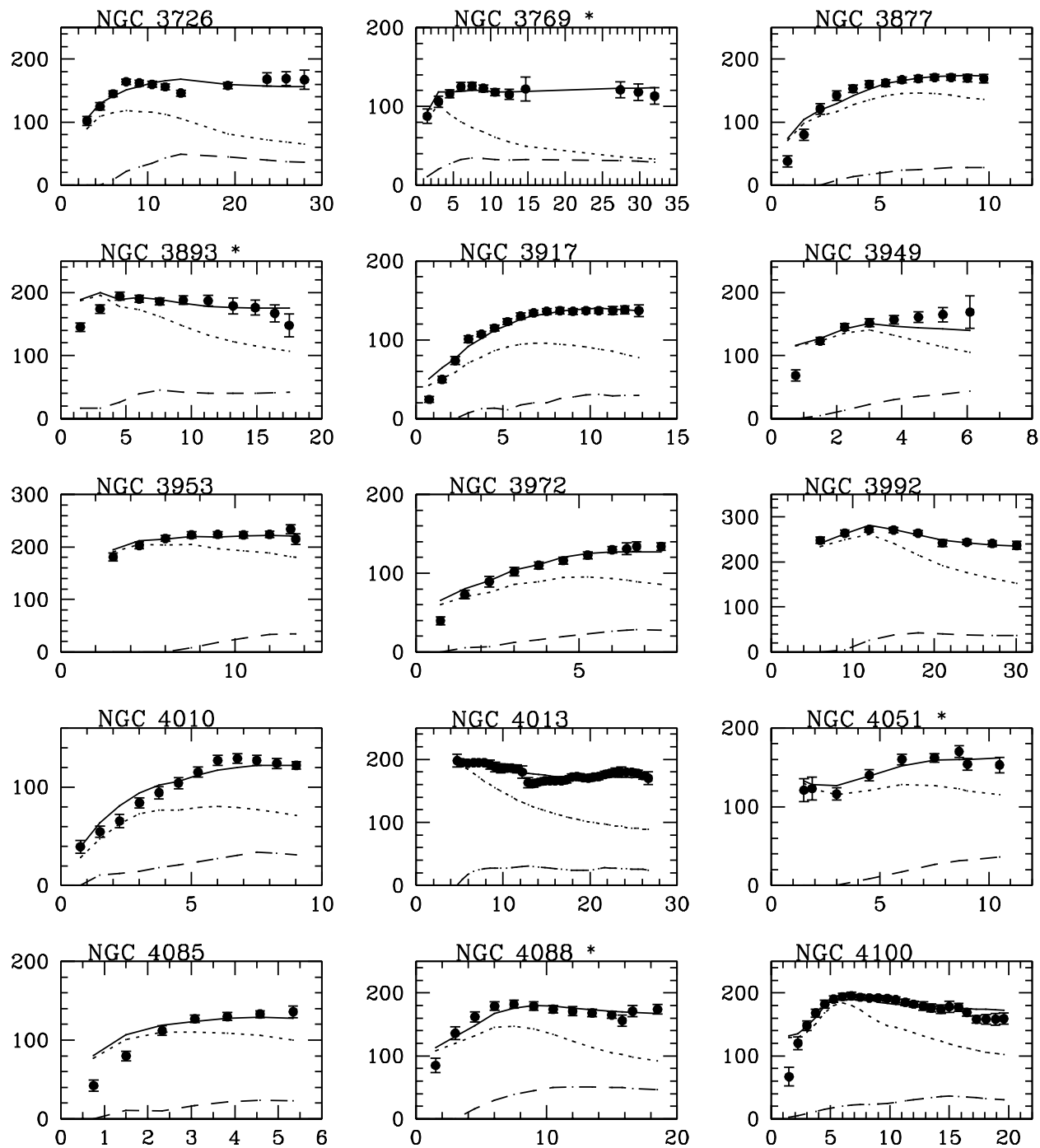
Surface Density \sim Surface Brightness

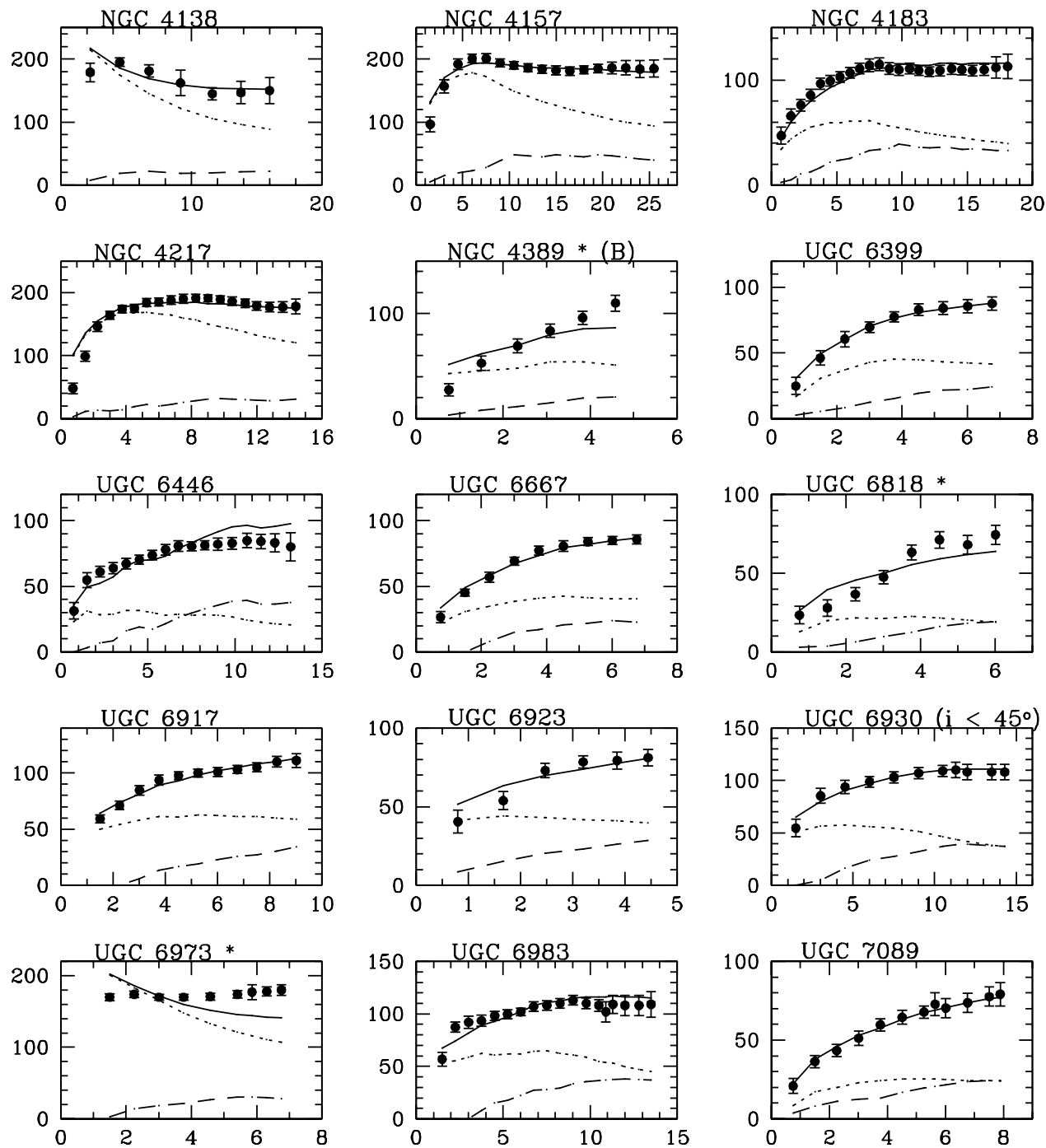
- Detailed Rotation Curve Fits

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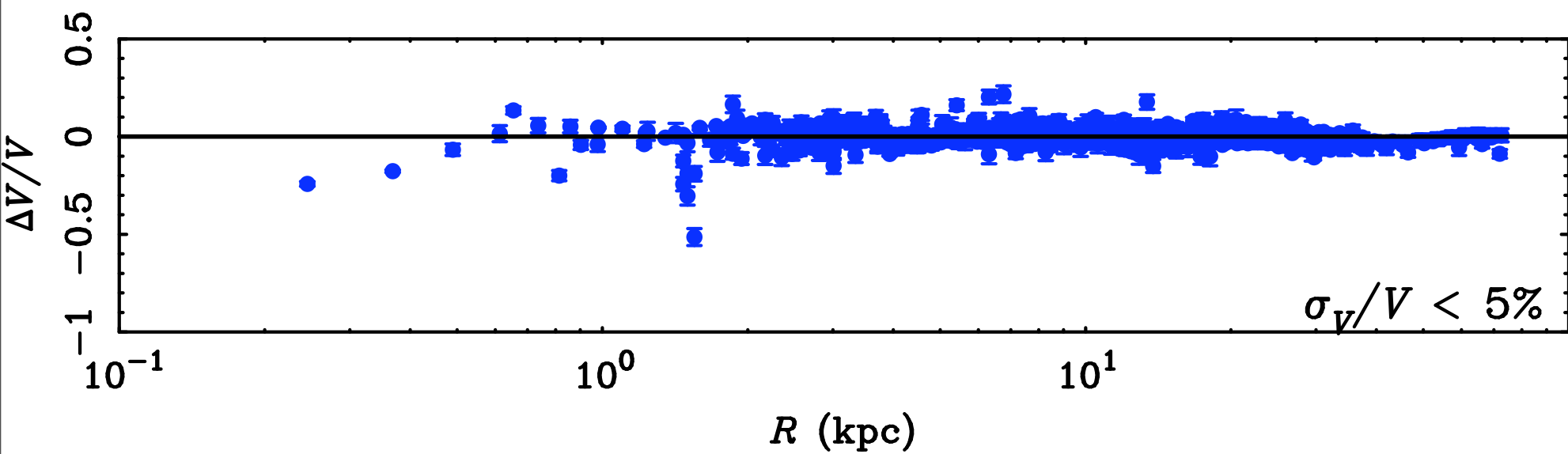
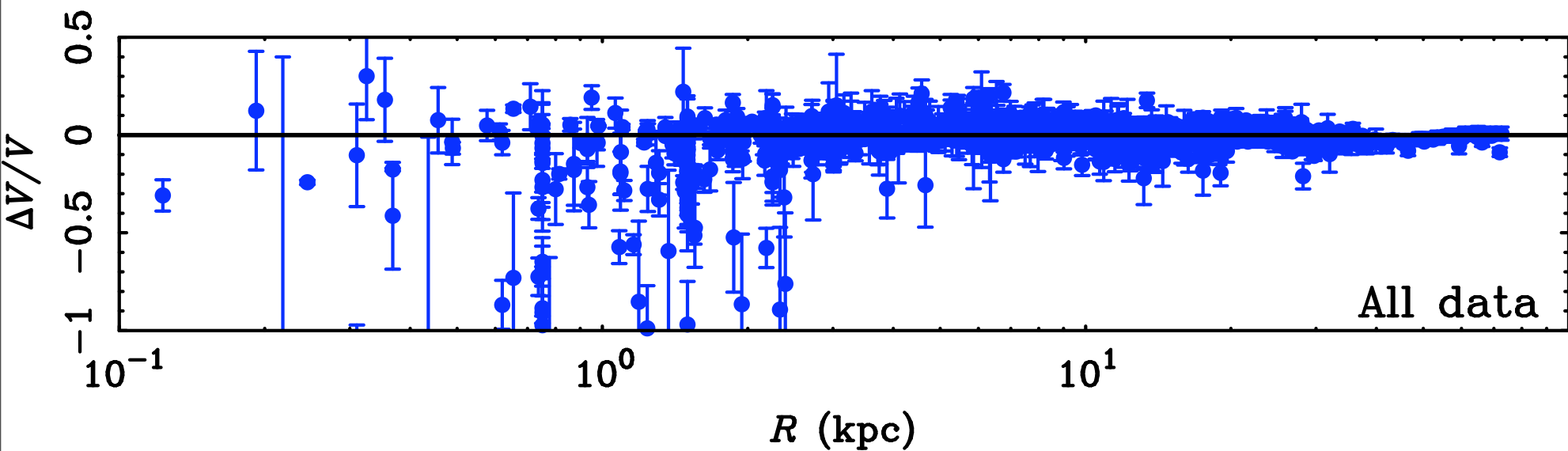








Residuals of MOND fits



MOND predictions

- The Tully-Fisher Relation

- ✓ Slope = 4
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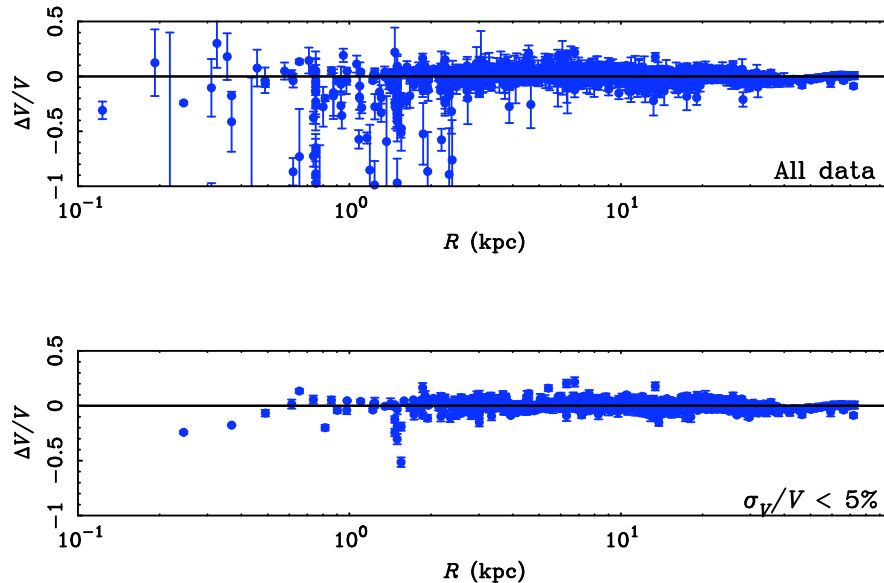
- ✓ Dependence of conventional M/L on radius and surface brightness

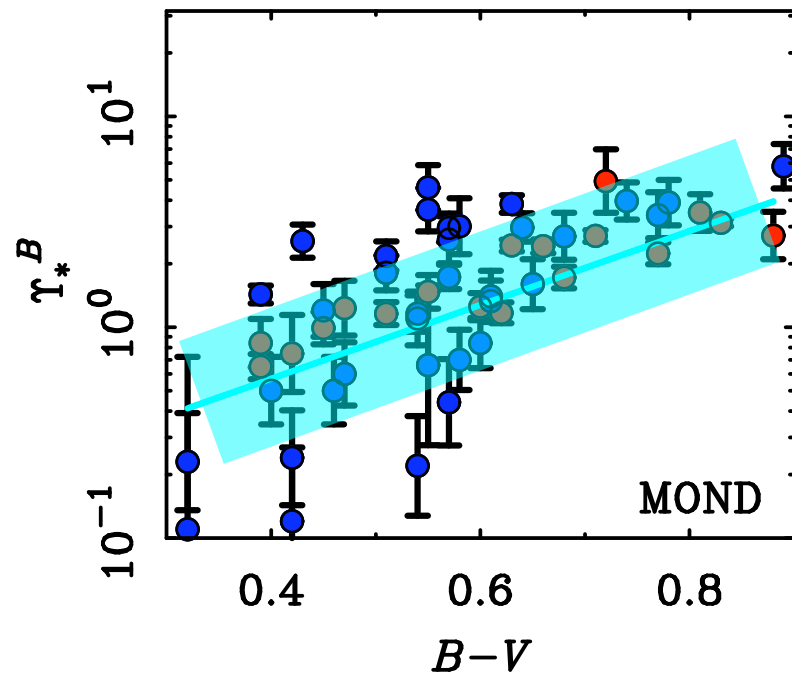
- ✓ Rotation Curve Shapes

- ✓ Surface Density \sim Surface Brightness

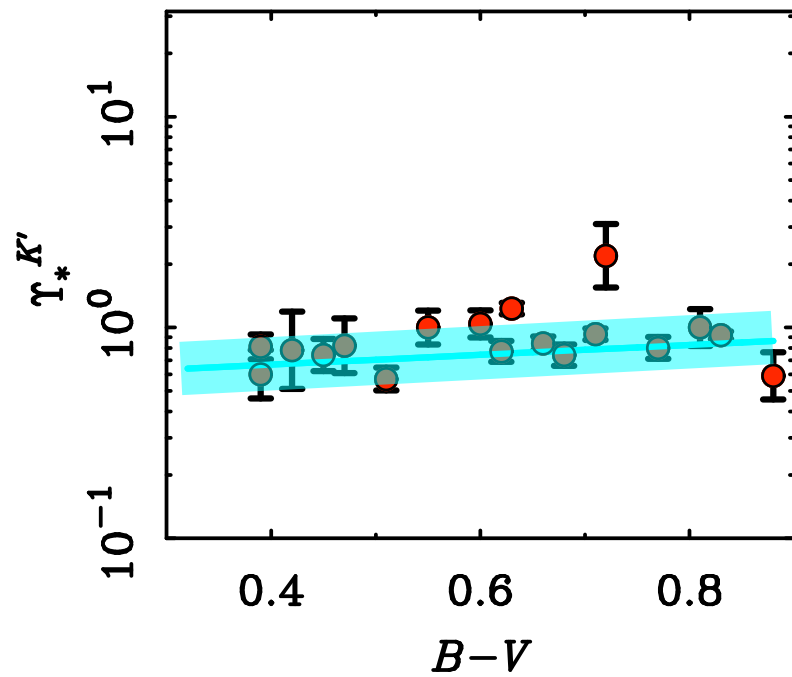
- ✓ Detailed Rotation Curve Fits

- Stellar Population Mass-to-Light Ratios





Line: stellar population model
(mean expectation)



MOND predictions

- The Tully-Fisher Relation

- ✓ Slope = 4
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- ✓ Fundamentally a relation between Disk Mass and V_{flat}
- ✓ No Dependence on Surface Brightness

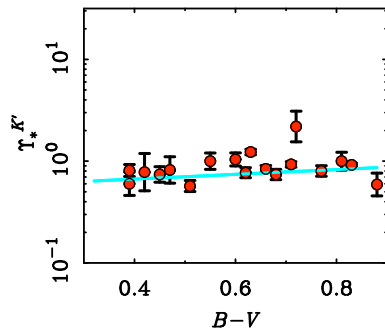
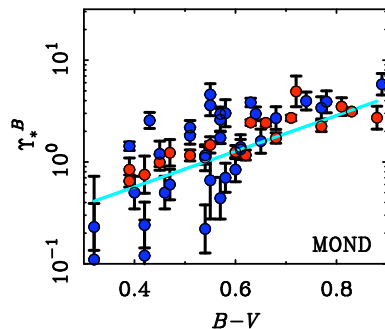
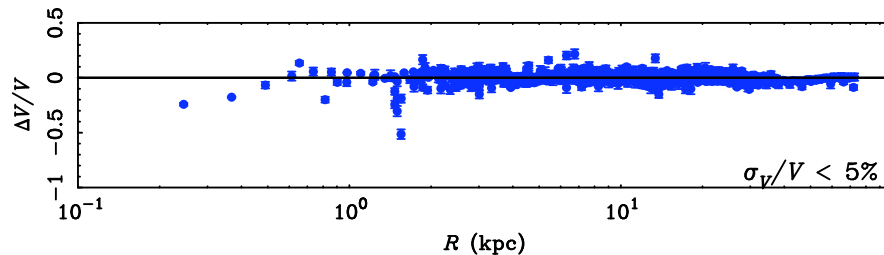
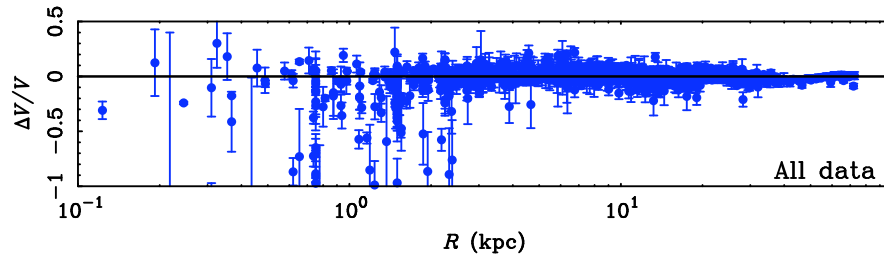
- ✓ Dependence of conventional M/L on radius and surface brightness

- ✓ Rotation Curve Shapes

- ✓ Surface Density \sim Surface Brightness

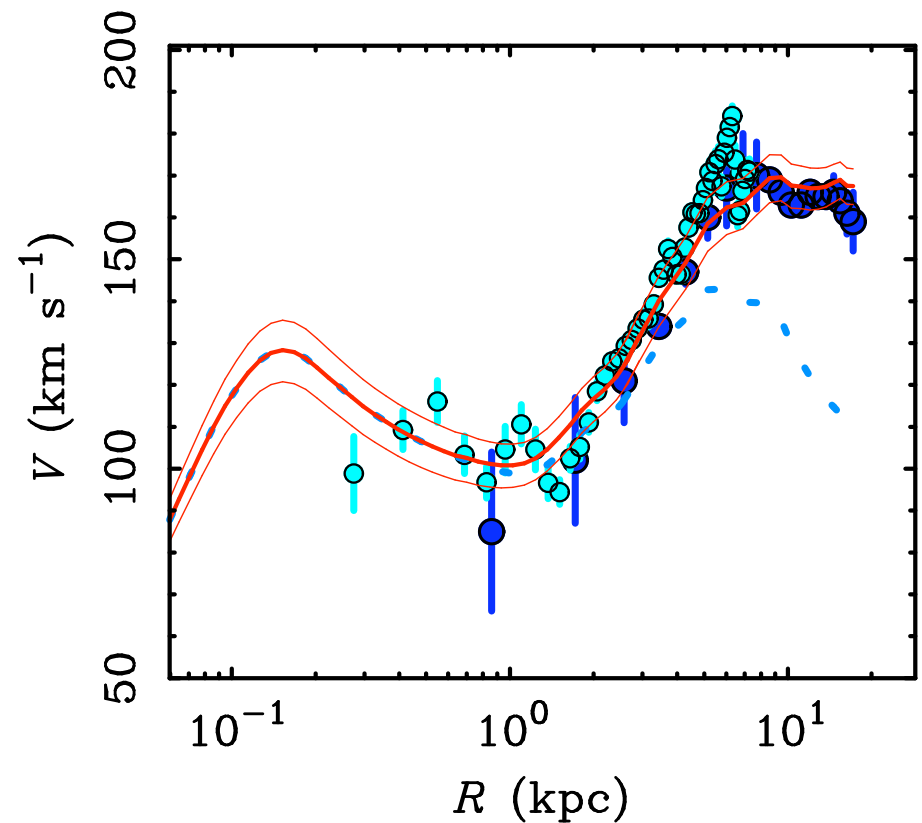
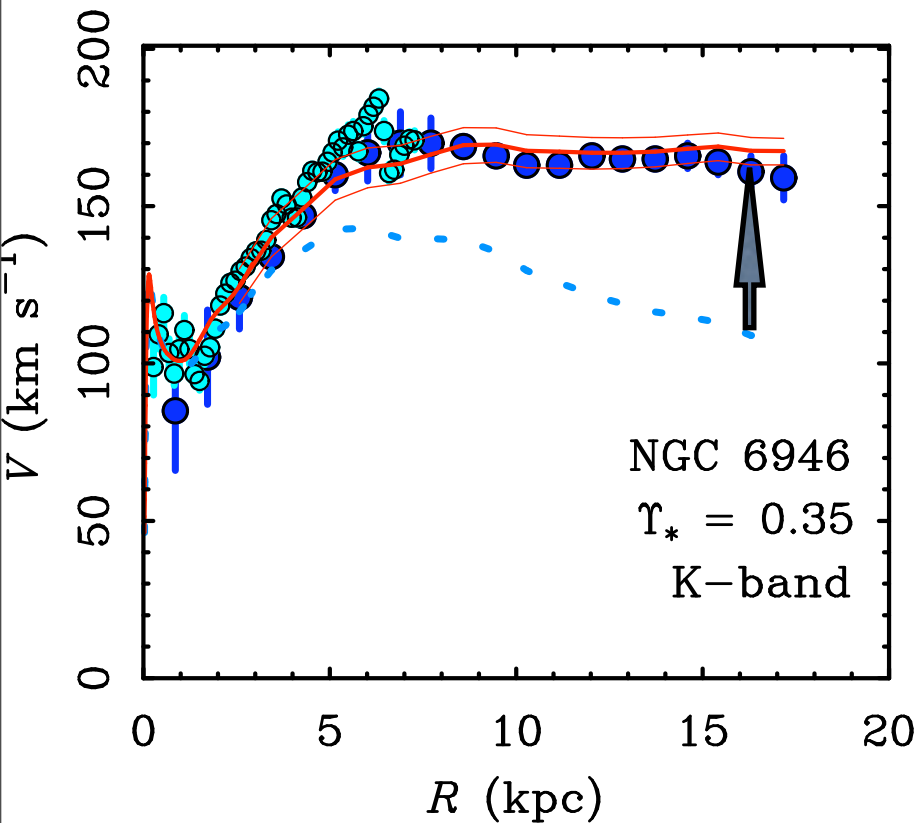
- ✓ Detailed Rotation Curve Fits

- ✓ Stellar Population Mass-to-Light Ratios



Renzo's Rule:

“When you see a feature in the light, you see a corresponding feature in the rotation curve.”

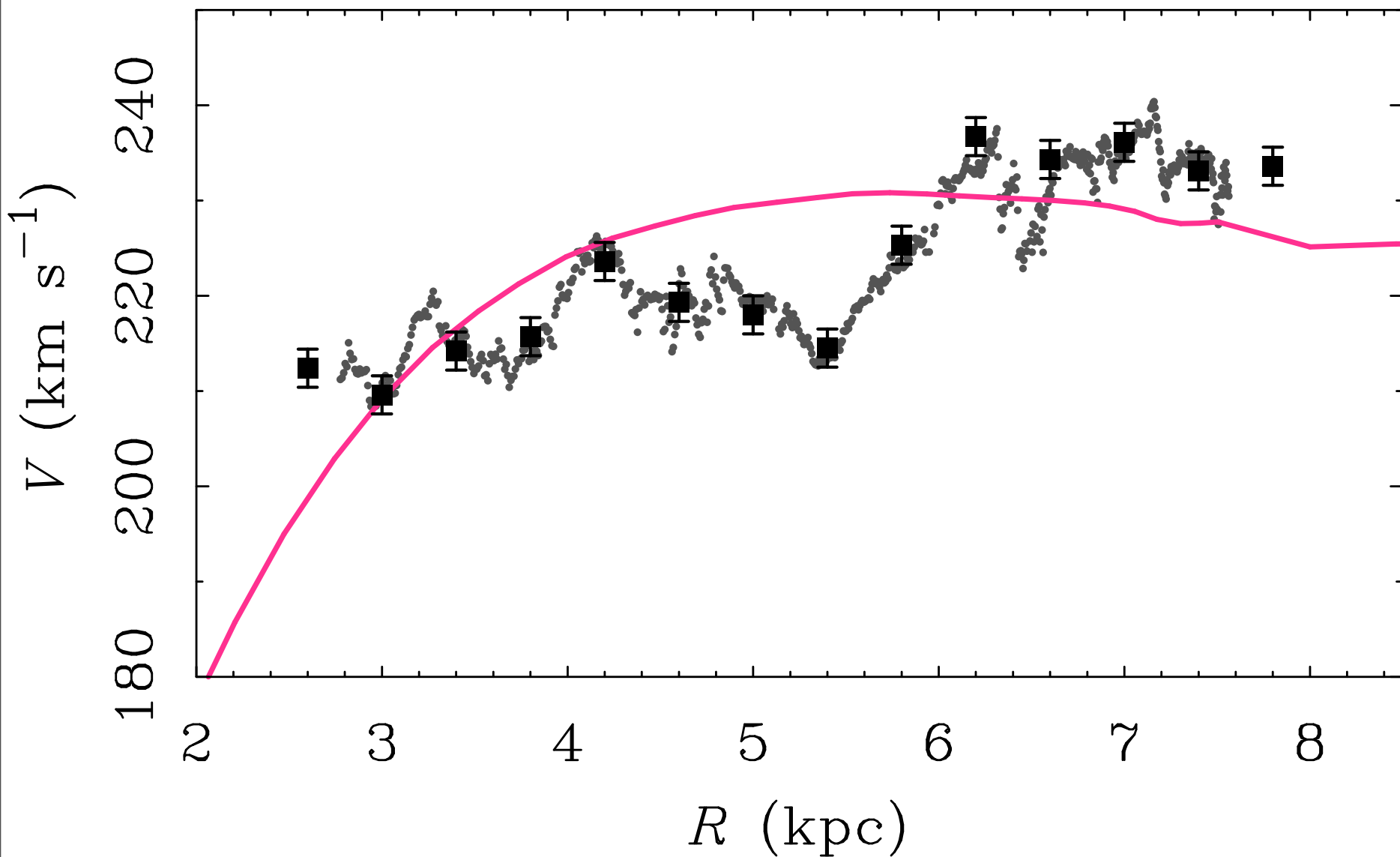


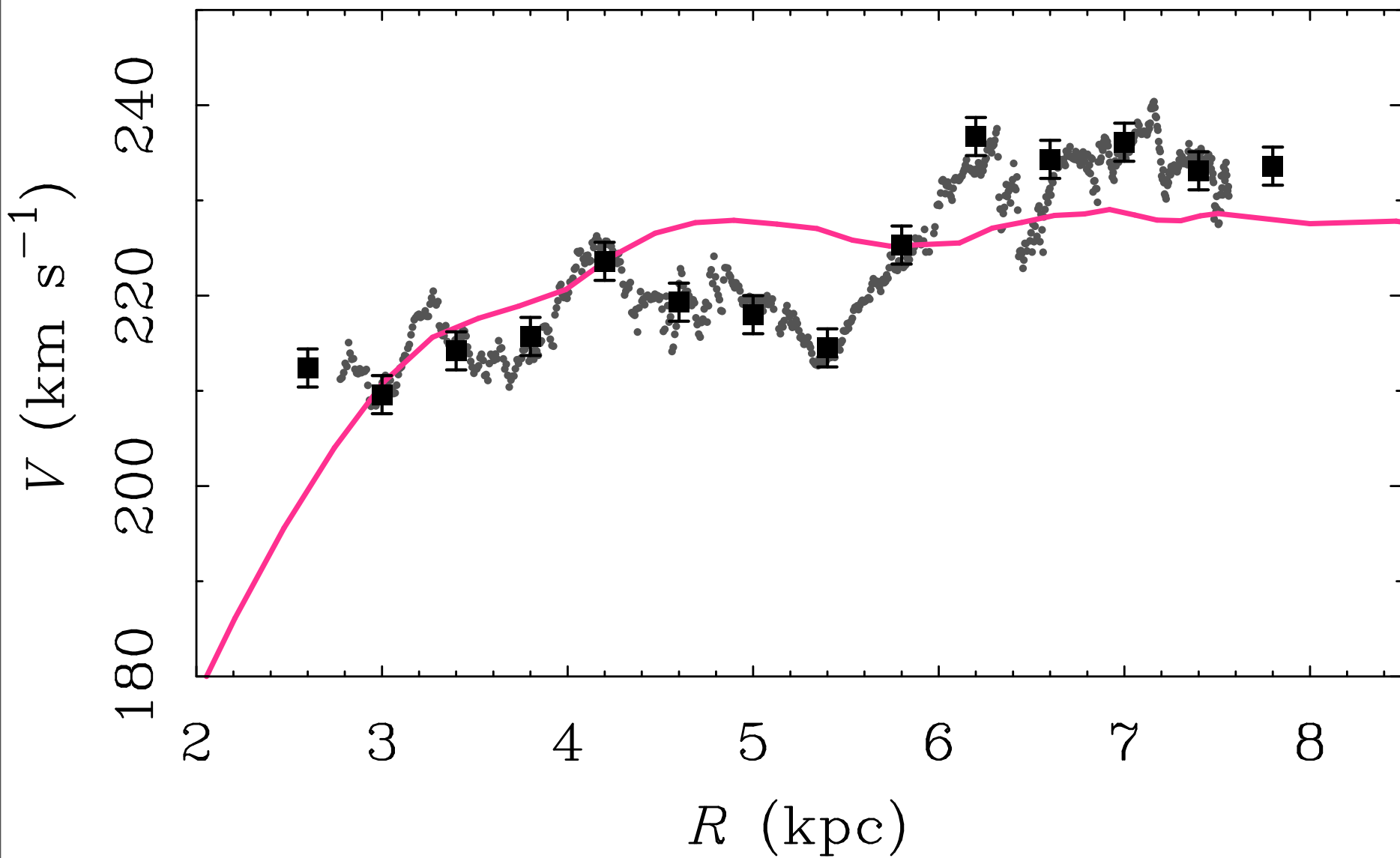
$$\Upsilon_*^K = 0.35 M_\odot / L_\odot$$

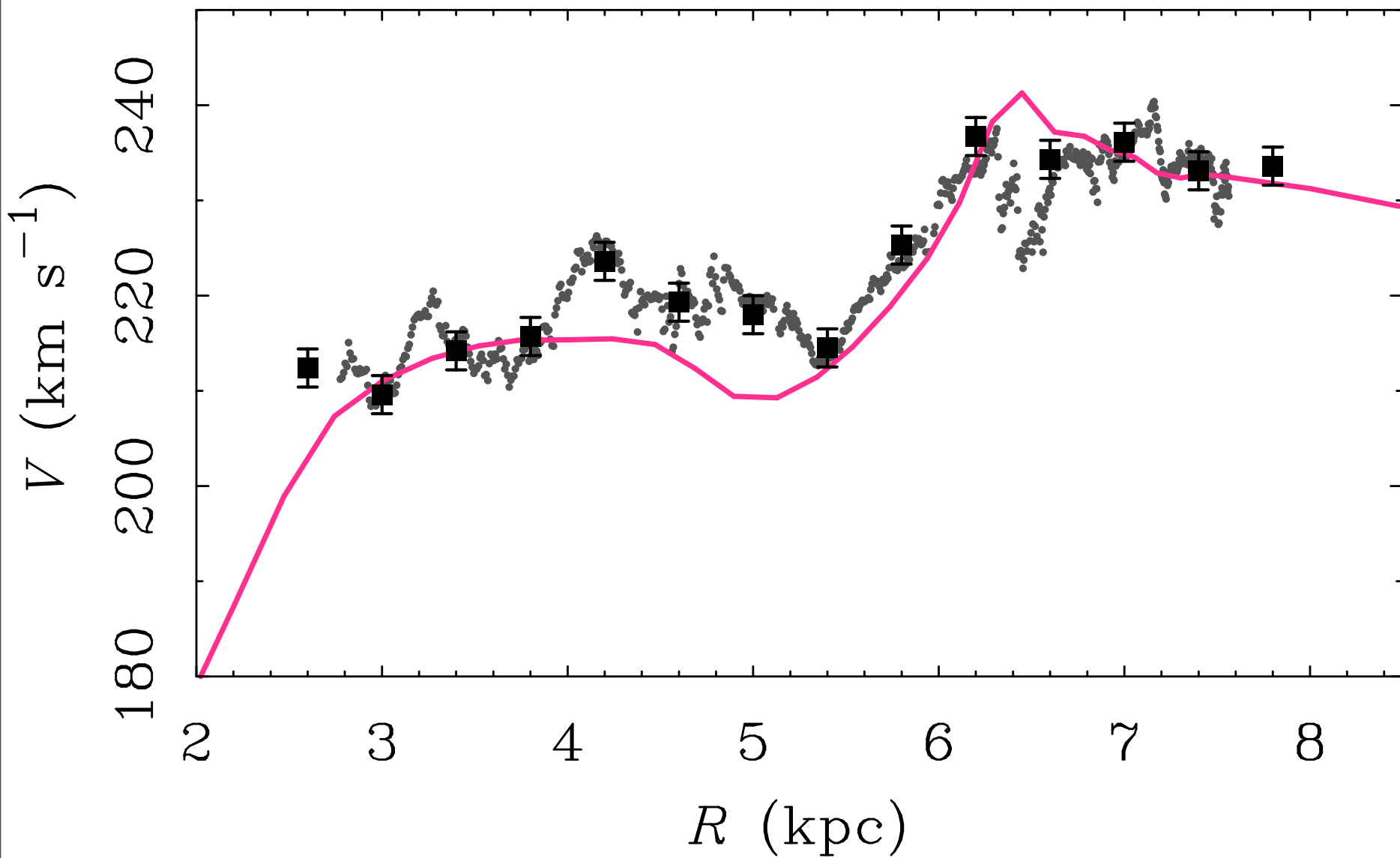
based on ppsynth models of Portinari et al. (2004)

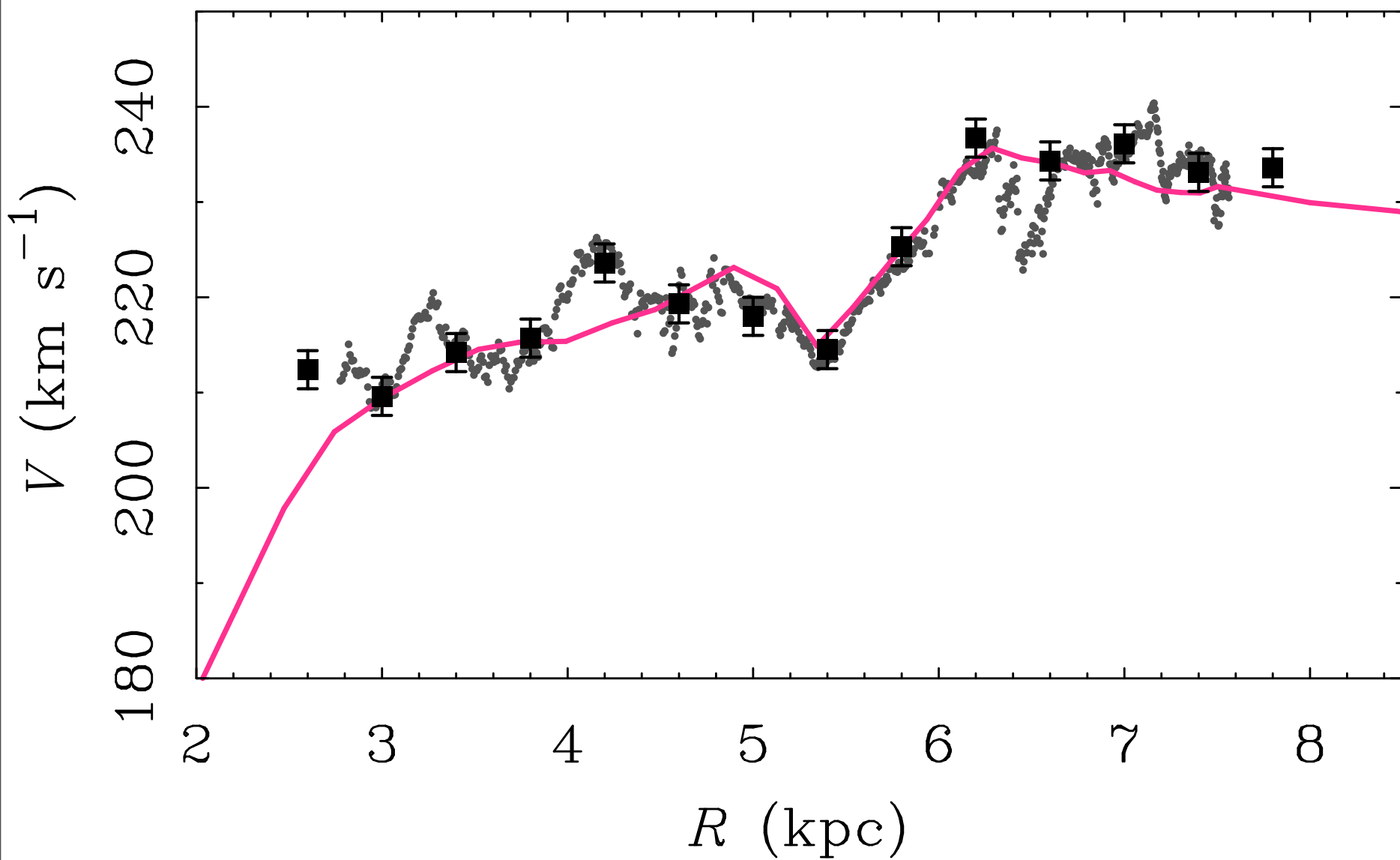
What about our own Galaxy?

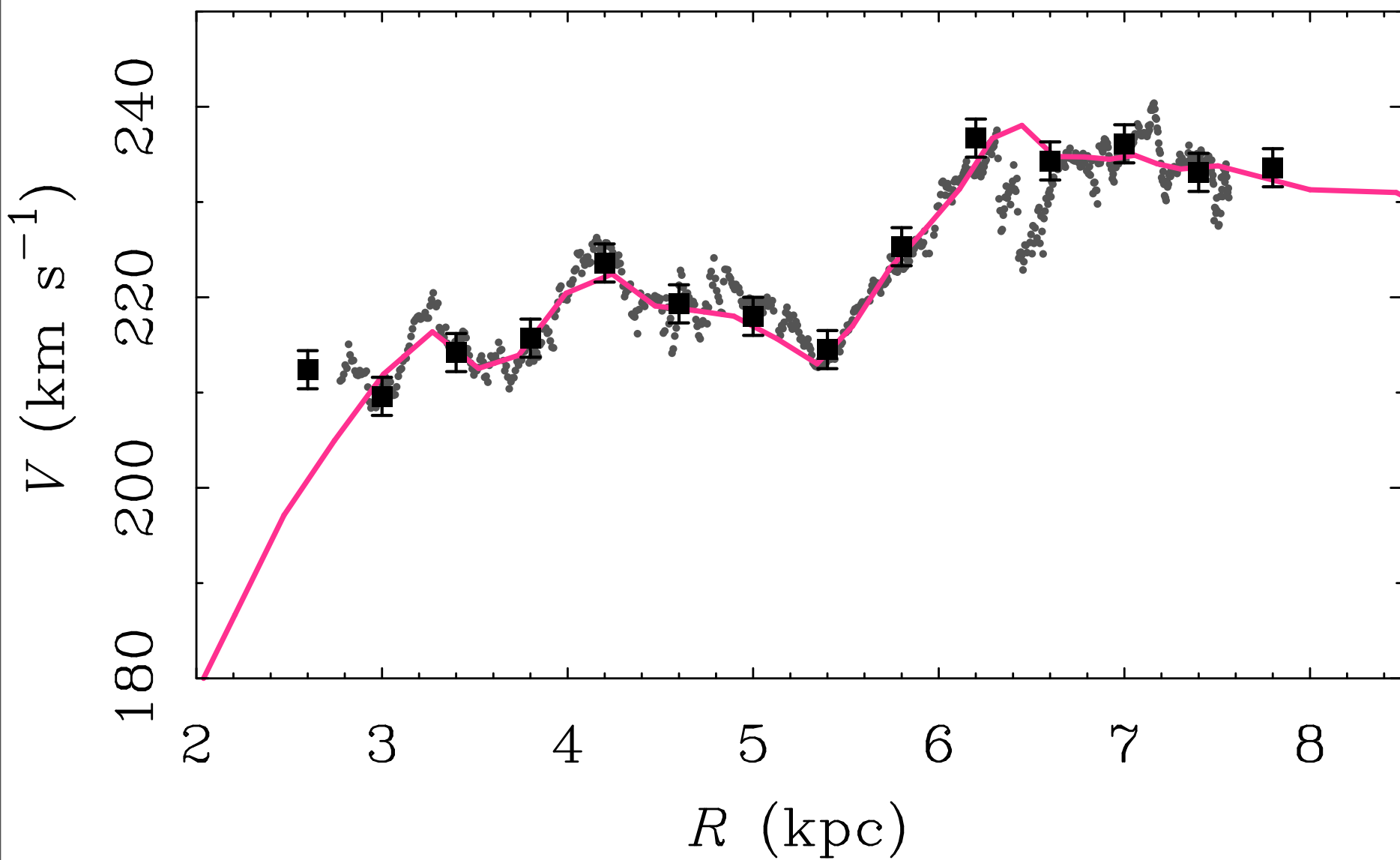
Luna et al. (2006: **CO**); McClure-Griffiths & Dickey (2007: **HI**)



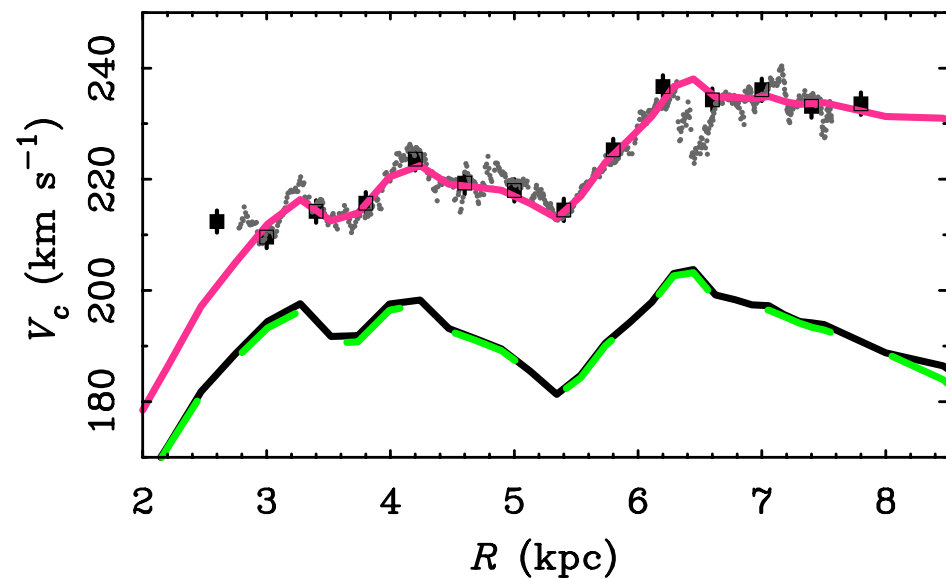
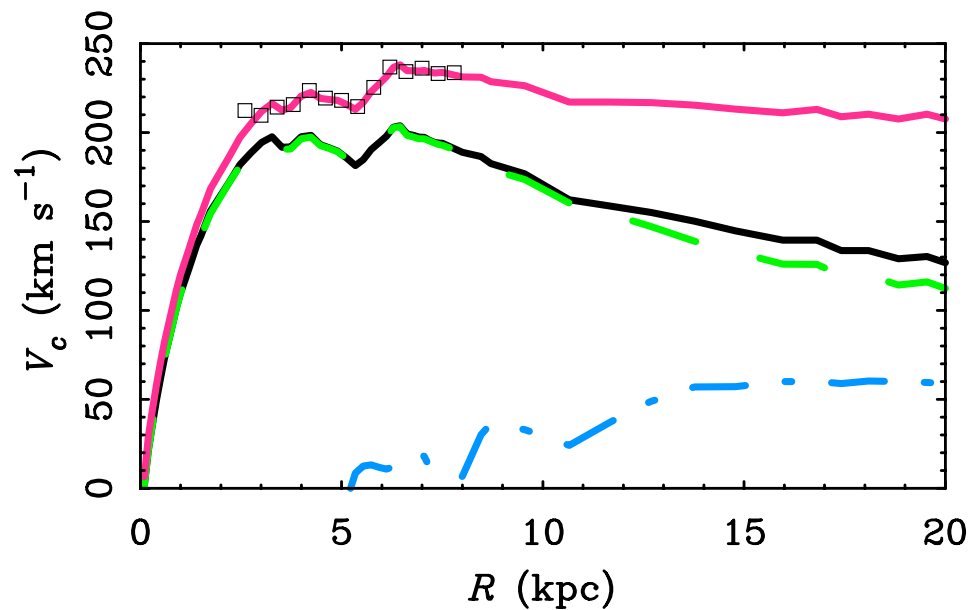
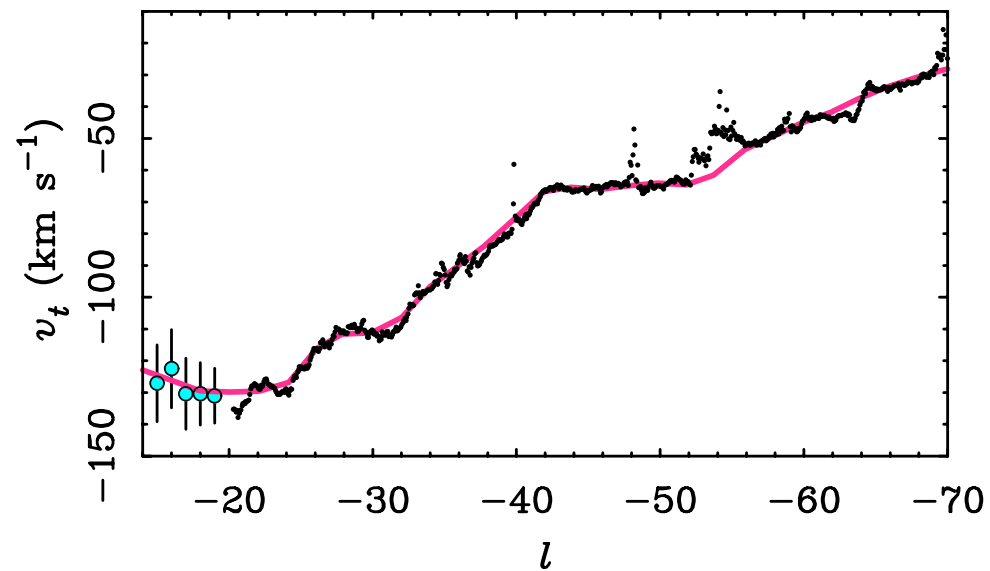
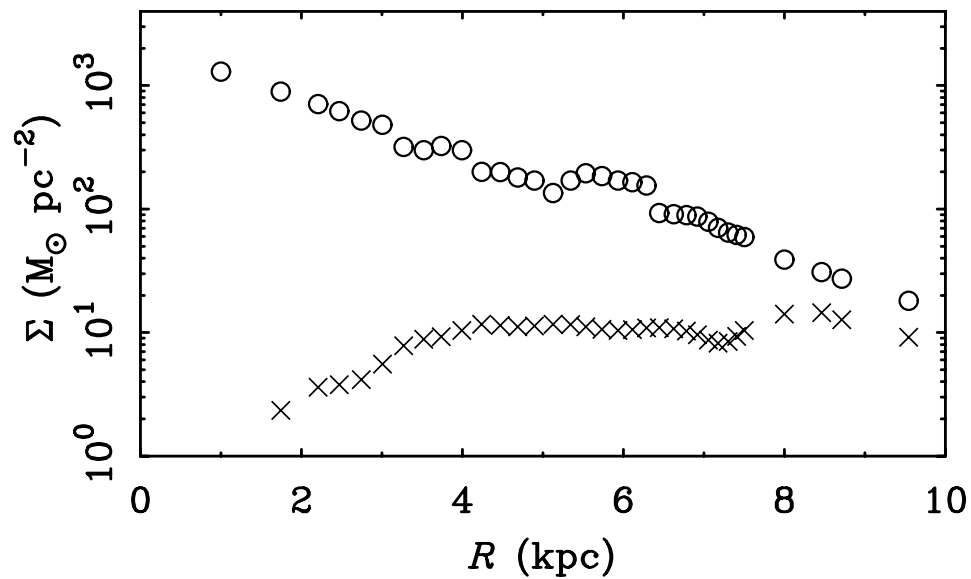




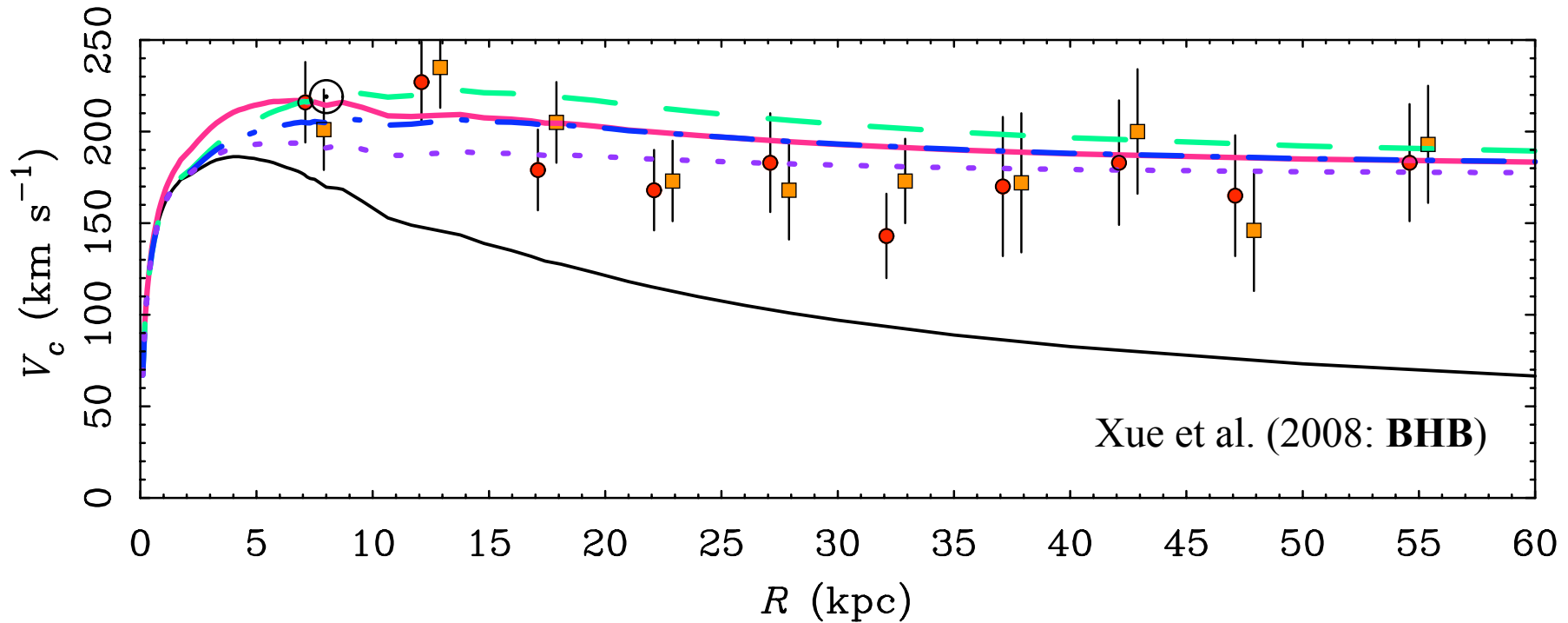




MONDian Milky Way



Get outer rotation curve with no fitting



Different lines represent different assumed interpolation functions



LSB galaxies
got spiral arms!

Need very massive
disks to drive spiral
density waves in
LSBs, as anticipated
by McGaugh & de Blok
(1998), *ApJ*, 499, 66

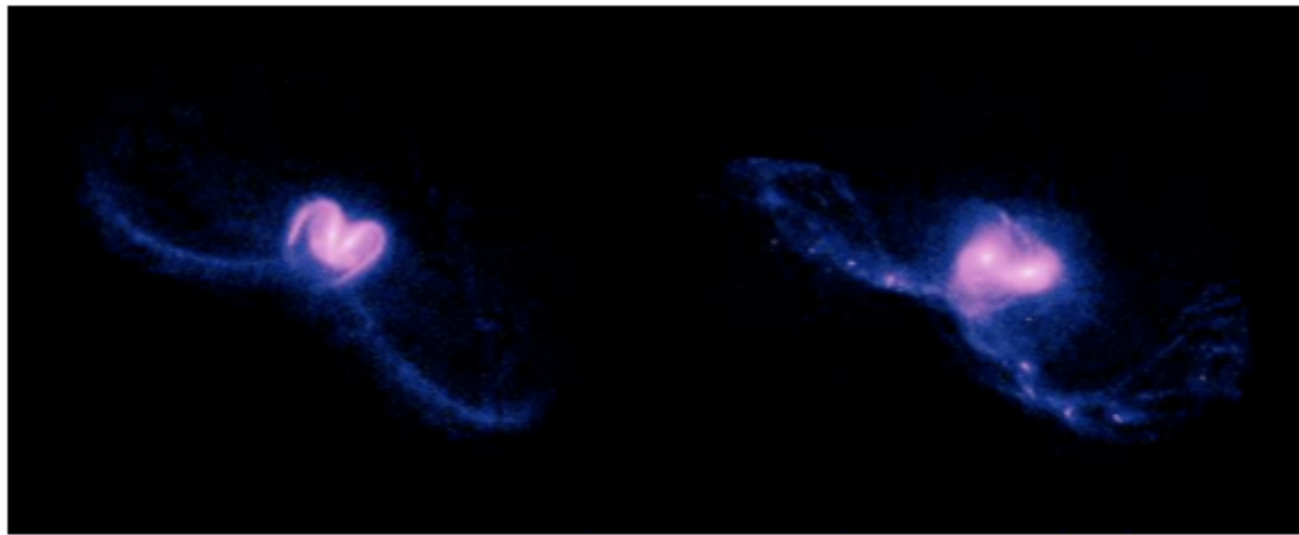
DISK STABILITY
MAKES MORE SENSE
IN MOND

Disk Masses from Density Waves

Galaxy	$(M/L)_*$
F568-1	14
F568-3	7
F568-6	11
F568-VI	16
UGC 128	4
UGC 1230	6
UGC 6614	8
ESO 14-40	4

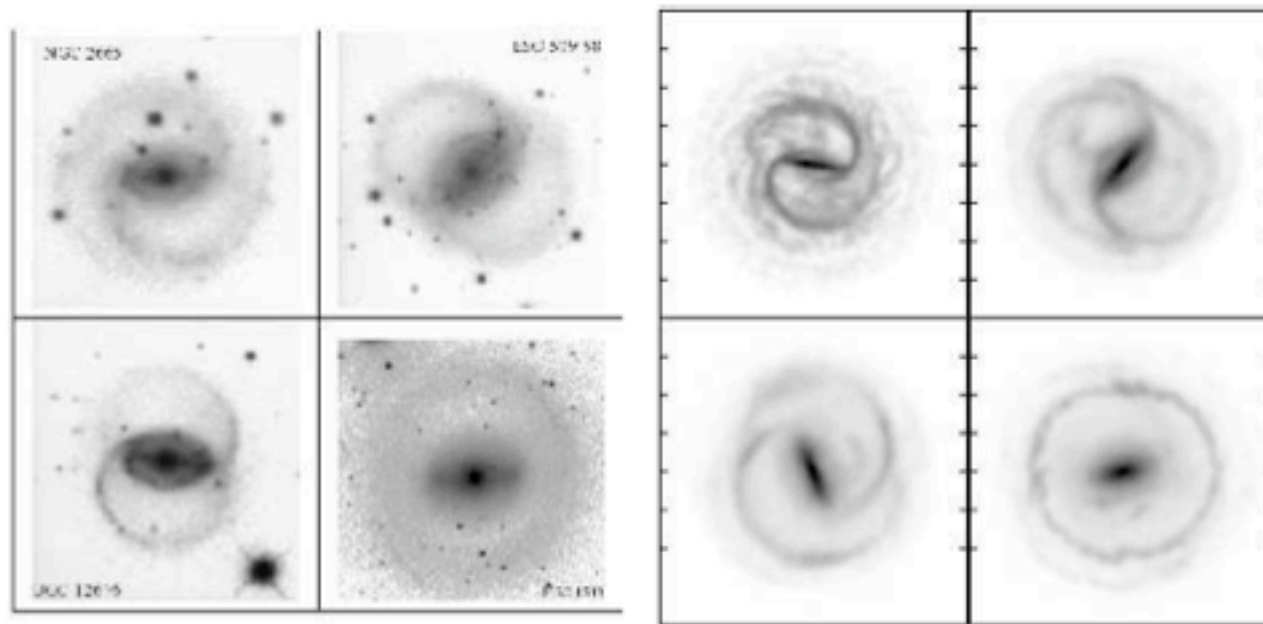
Big $(M/L)_*$'s!

from B. Fuchs, *astro-ph/0209157*



Tiret & Combes

Fig. 5. Simulations of the Antennae galaxies in the DM model (left) and MOND model(right).



Several examples showing the morphological structures of NGC 2665, ESO 509-98, UGC 12646 and NGC 1543 (top panel) compared to simulated galaxies in MOND (bottom panel). Rings and pseudo-rings structures are well reproduced with modified gravity.

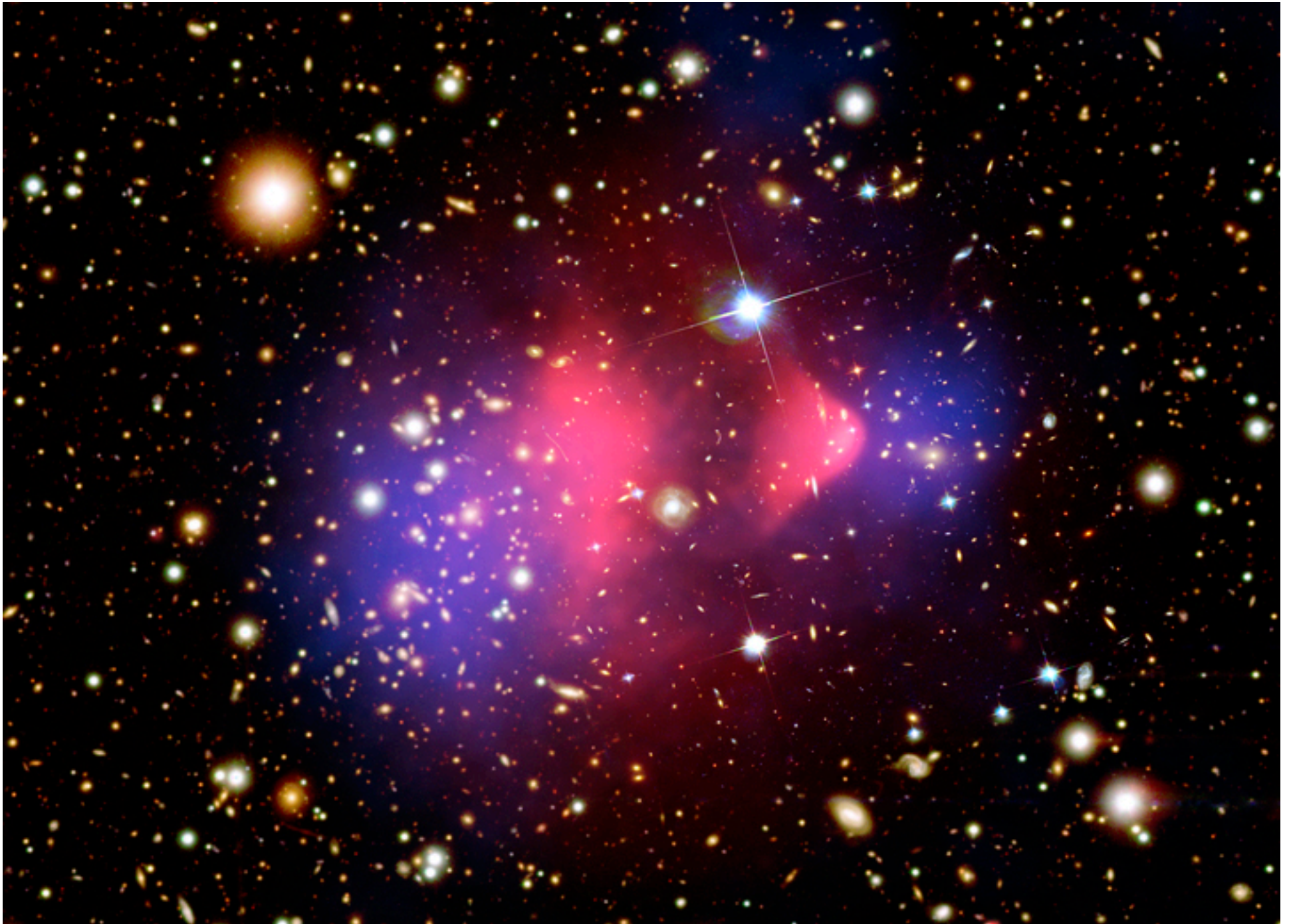
What are the downsides?

- You don't know the Power of the Dark Side
- Can MOND explain large scale structure?
- Can it provide a satisfactory cosmology?
- Can it be reconciled with General Relativity? **TeVes**
Bekenstein (2004)
- Does it survive other tests?

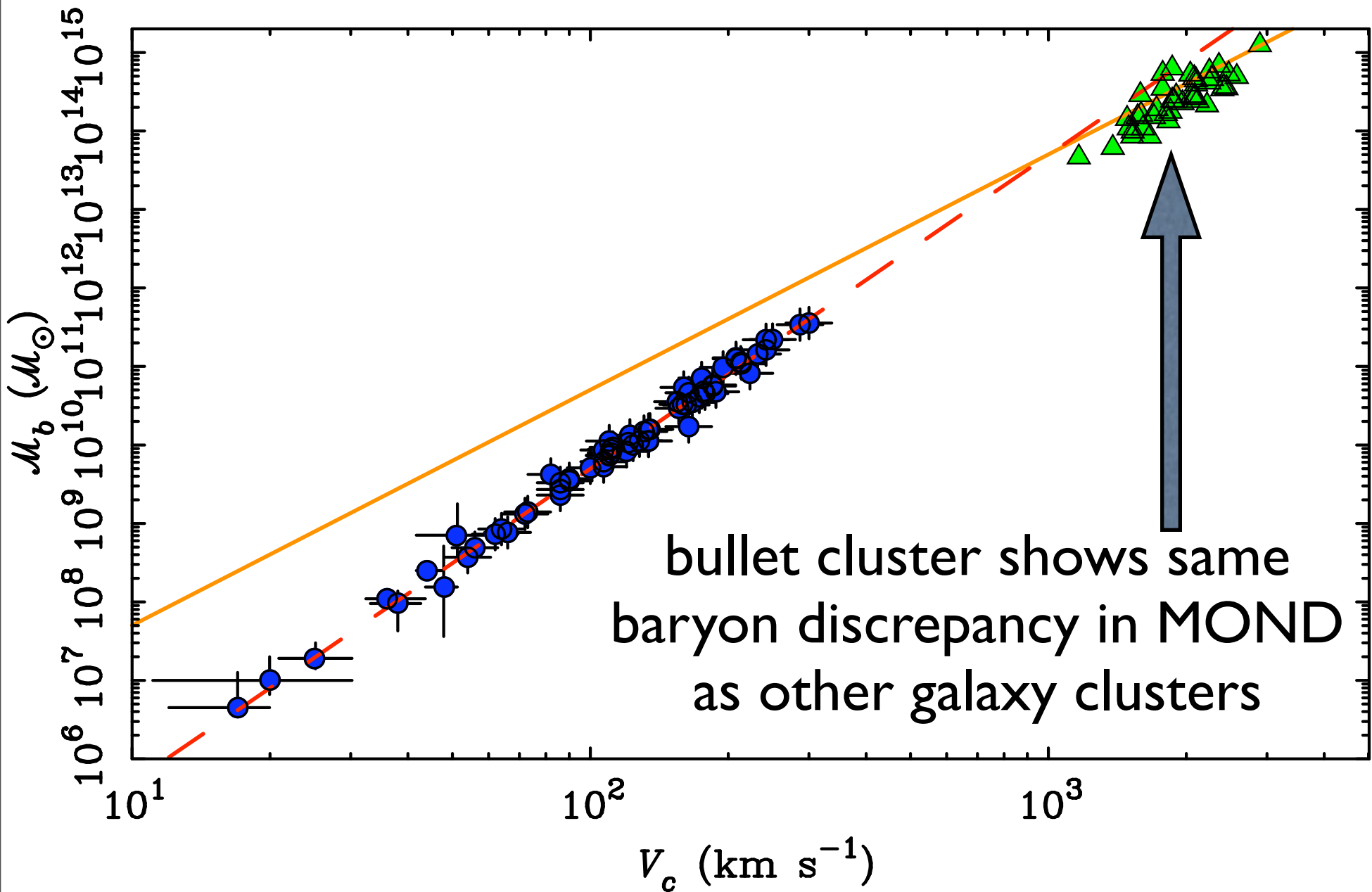


Clusters problematic

1E 0657-56 - “bullet” cluster (Clowe et al. 2006)

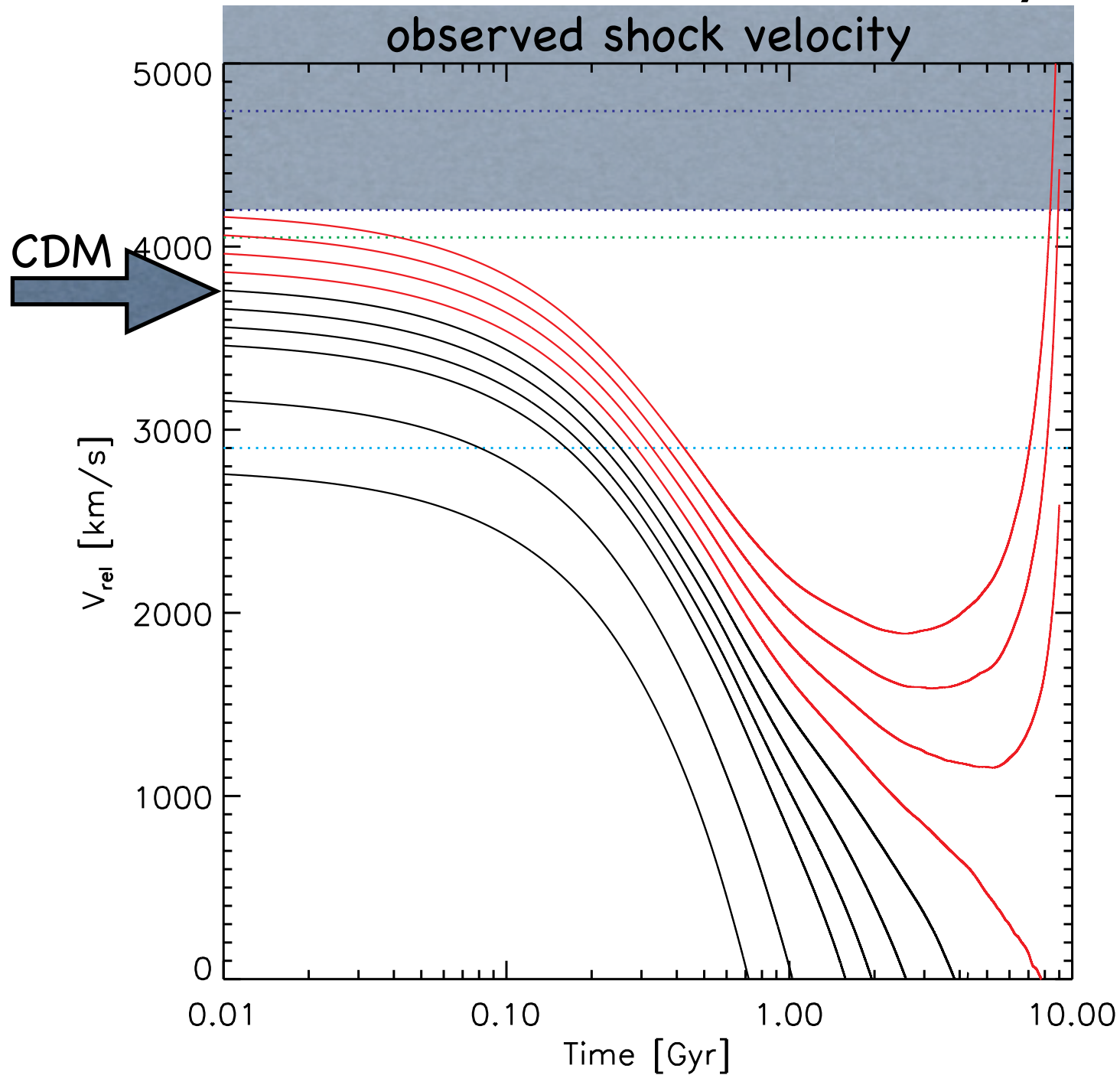


direct proof of dark matter?

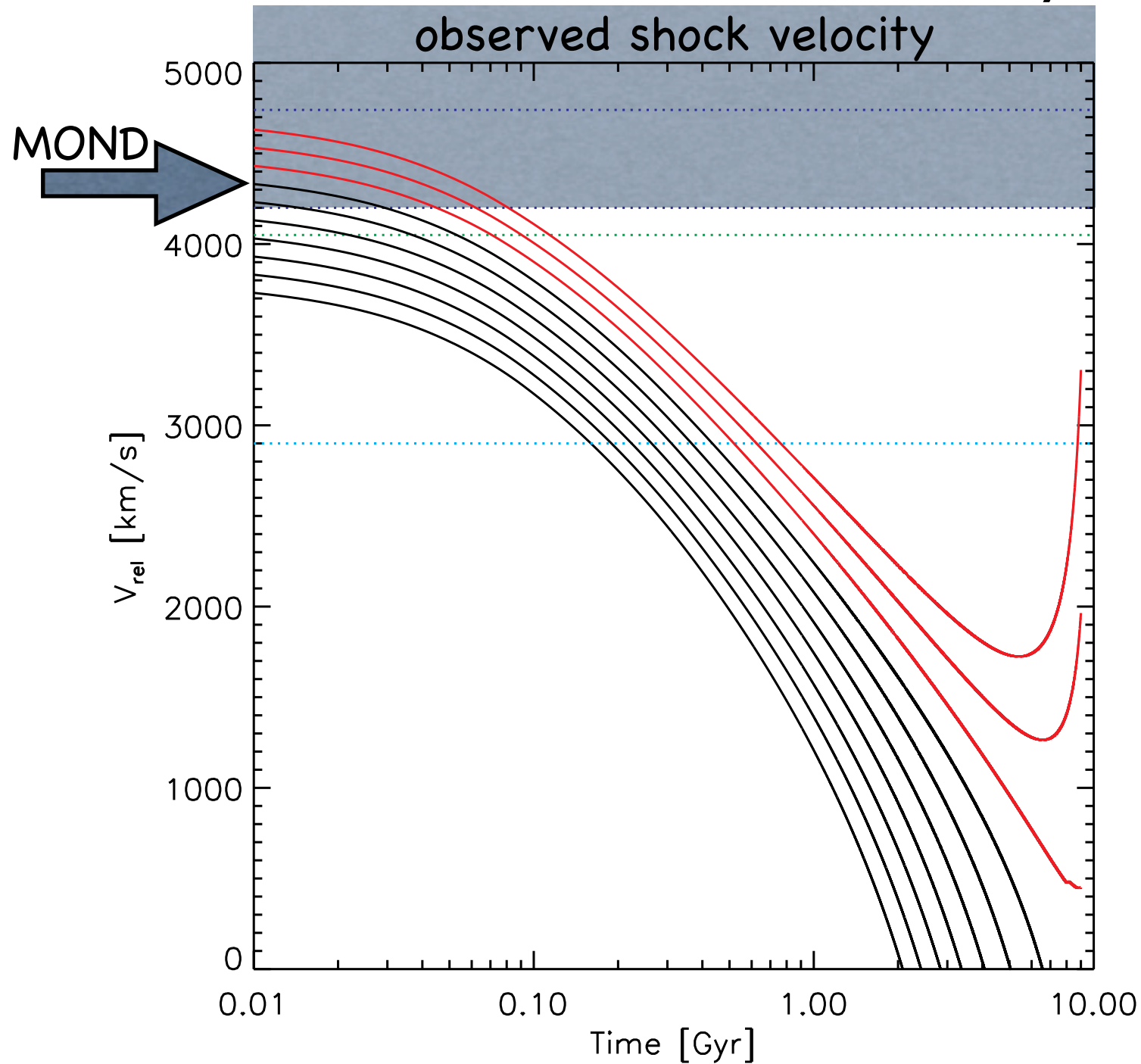


MOND suffers a missing mass problem!
unseen baryons? heavy neutrinos?

bullet cluster collision velocity

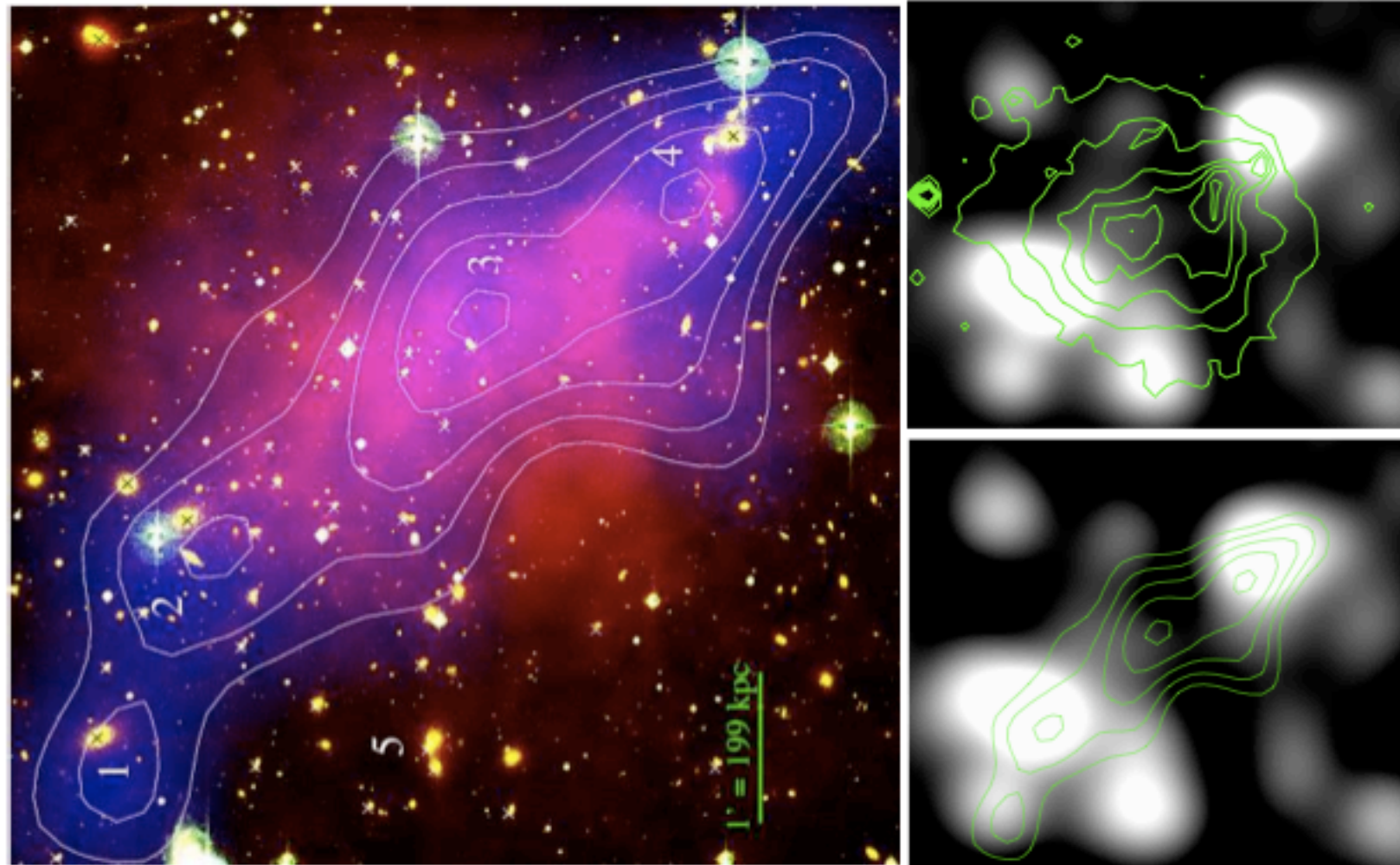


bullet cluster collision velocity



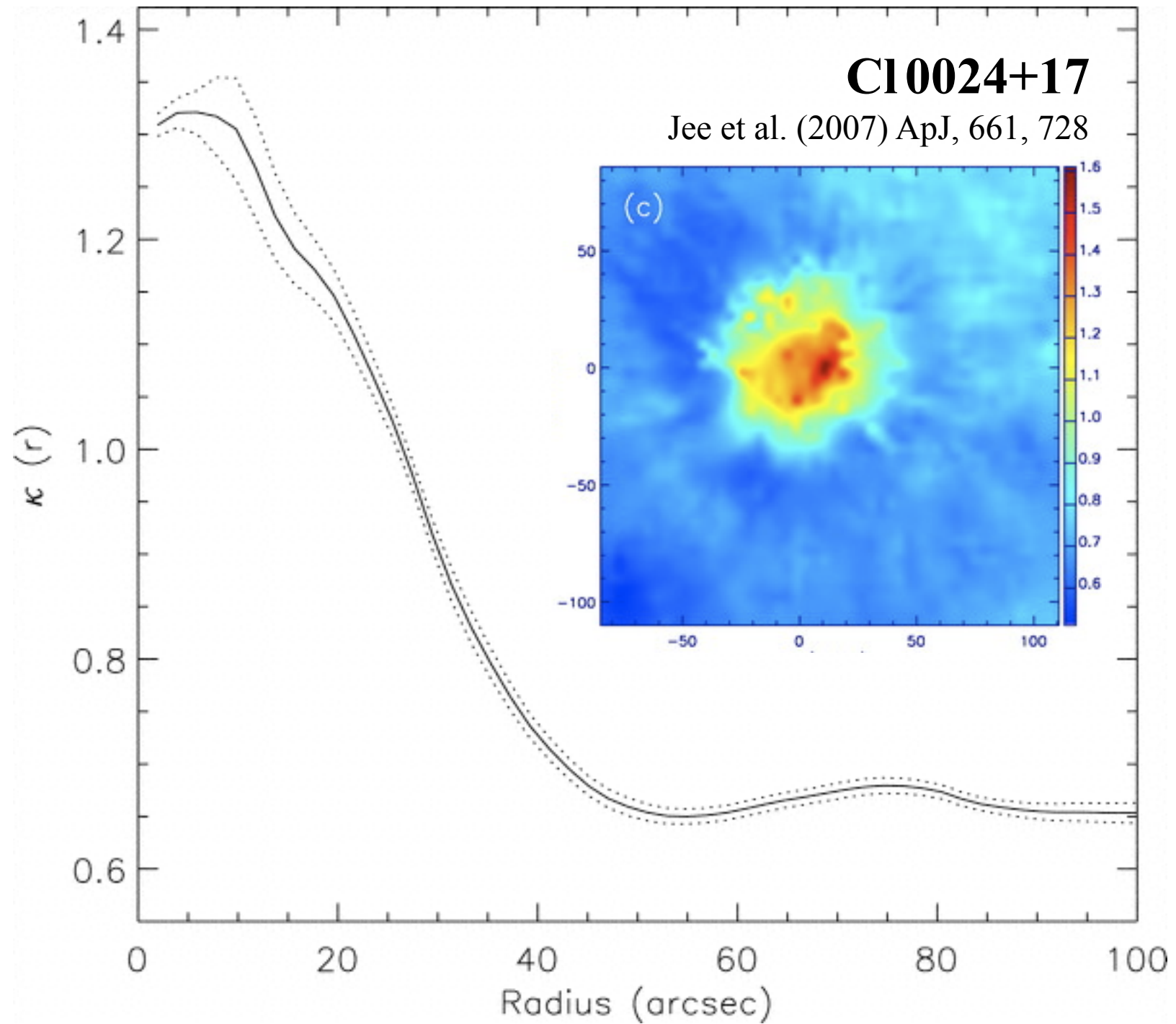
Mahdavi et al. (2007) arXiv:0706.3048

Abell 520 - Counter-example to bullet cluster
with a mass peak devoid of galaxies



CI0024+17

Jee et al. (2007) ApJ, 661, 728



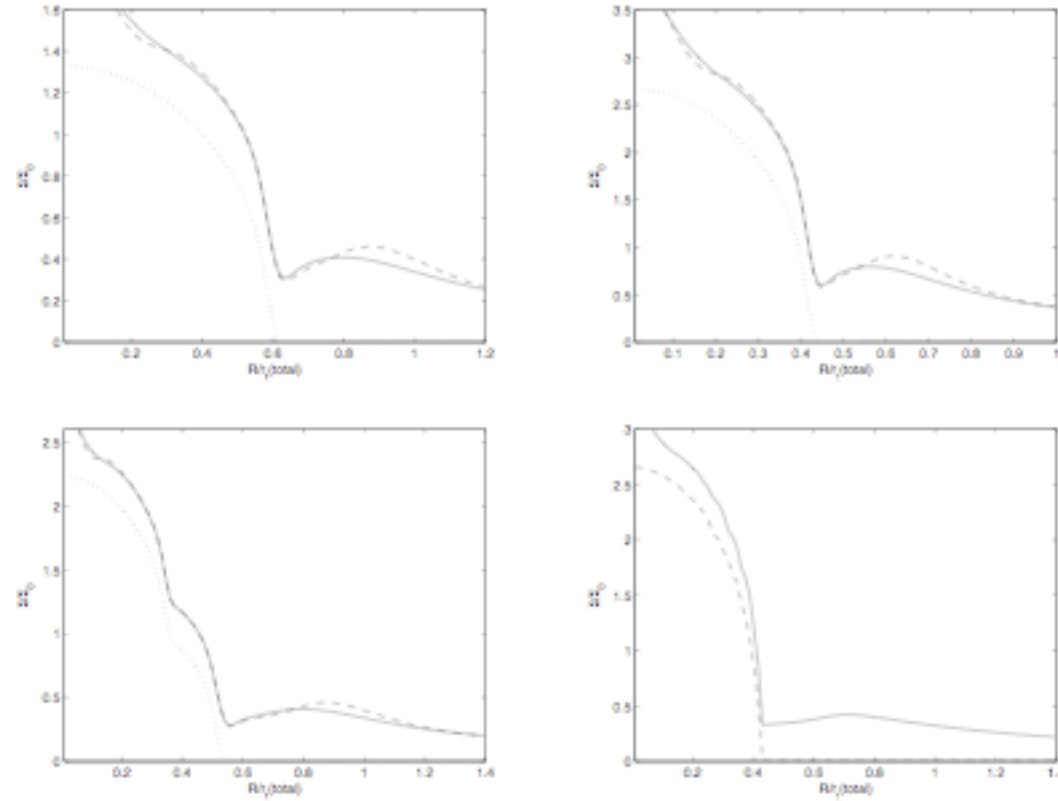


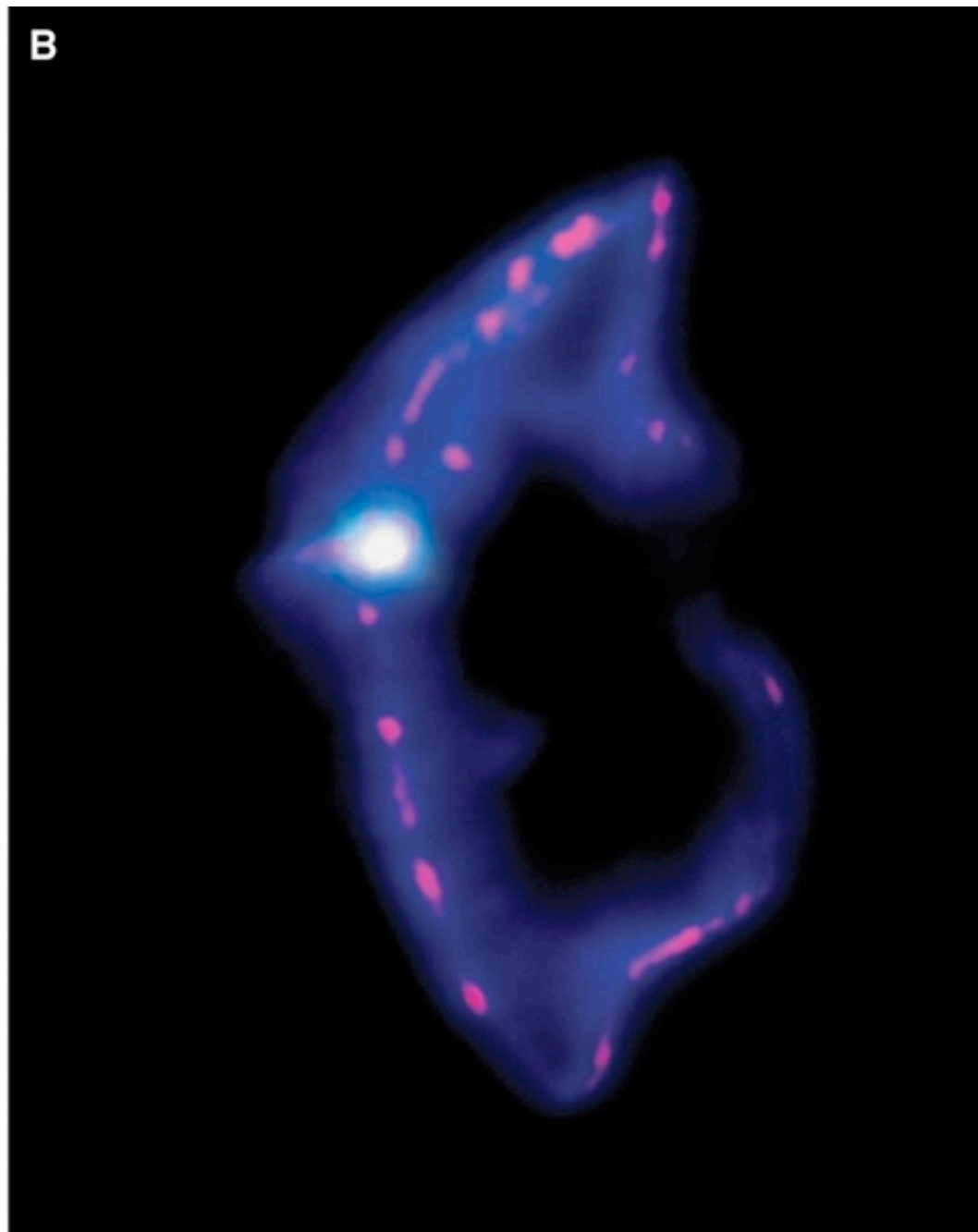
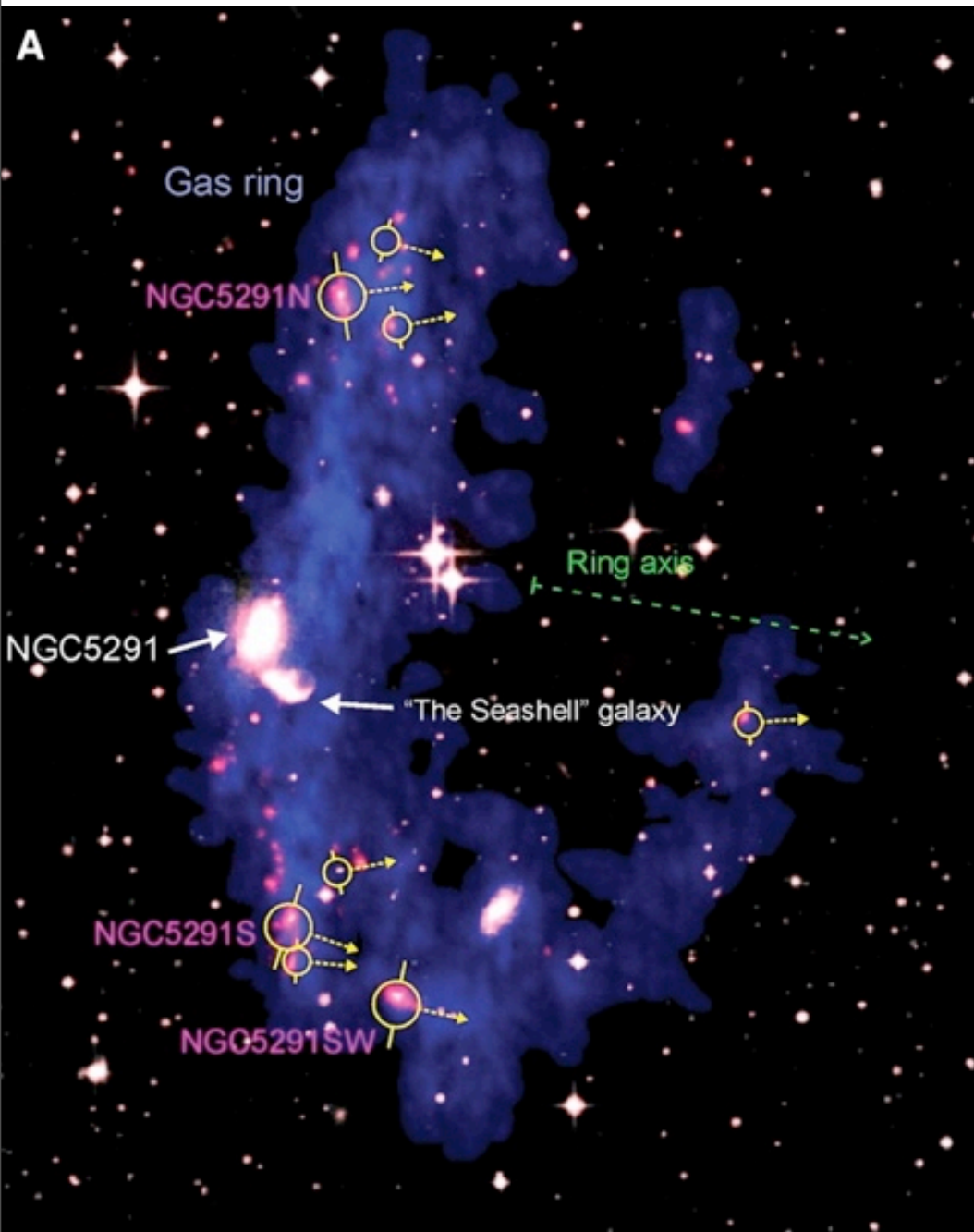
Fig. 6.— The total projected Surface density in units of Σ_0 . Upper left: for a single sphere of constant density with a radius that is 0.6 the transition radius. Upper right: for two spheres of constant density far apart from each other along the line of sight, each has a radius that is 0.6 of its own transition radius. Lower left: for two concentric spheres of constant densities of masses 1 and 0.3 and radii 0.53 and 0.35 of the total transition radius. All these for two interpolating functions: $\bar{\mu}_2$ (solid) and $\bar{\mu}_3$ (dashed). In each case the baryon contribution alone is shown as the dotted line. Lower right: a dumbbell of two equal spherical masses of constant density far apart along the line of sight with μ_{10} (the source, baryon, contribution in dashed line).

$$r_t = \sqrt{\frac{GM}{a_0}}$$

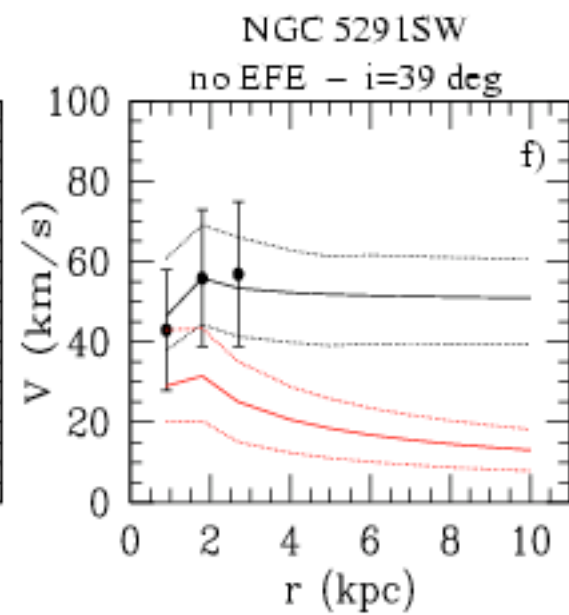
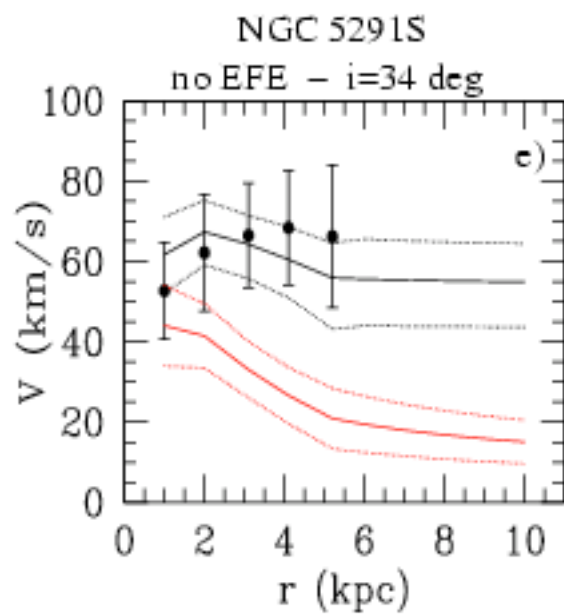
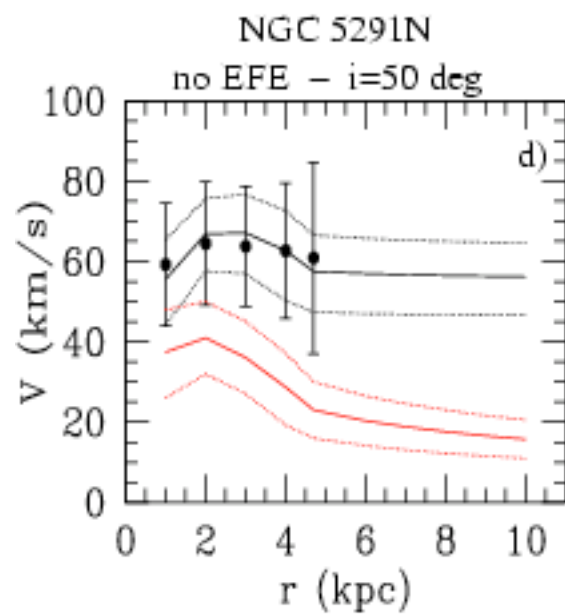
There can be a feature around the transition radius, depending on the interpolation function.

The ring reported by Jee et al. may be such a feature.

Tidal Debris Dwarfs - should be devoid of Dark Matter

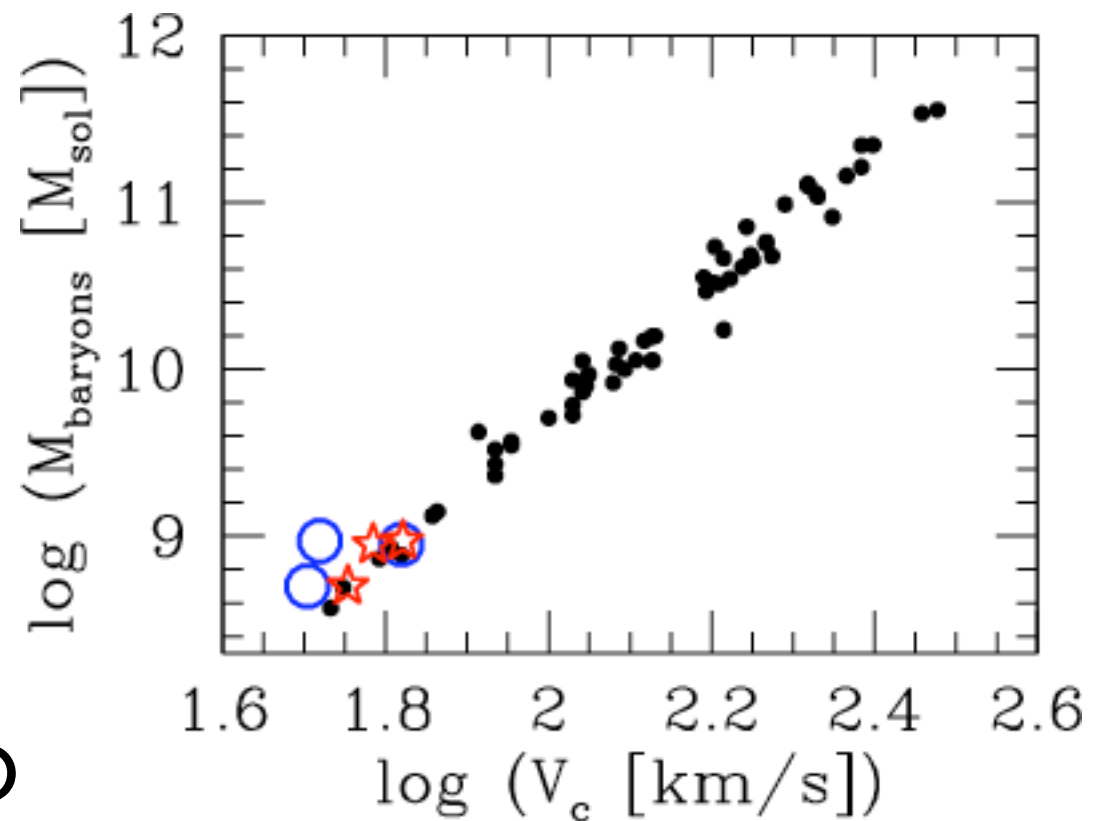


Bournaud et al. (2007) *Science*, 316, 1166



Gentile et al. (2007)
A&A, 472, L25

Tidal dwarfs
do show mass
discrepancies as
expected in MOND



Tiret & Combes

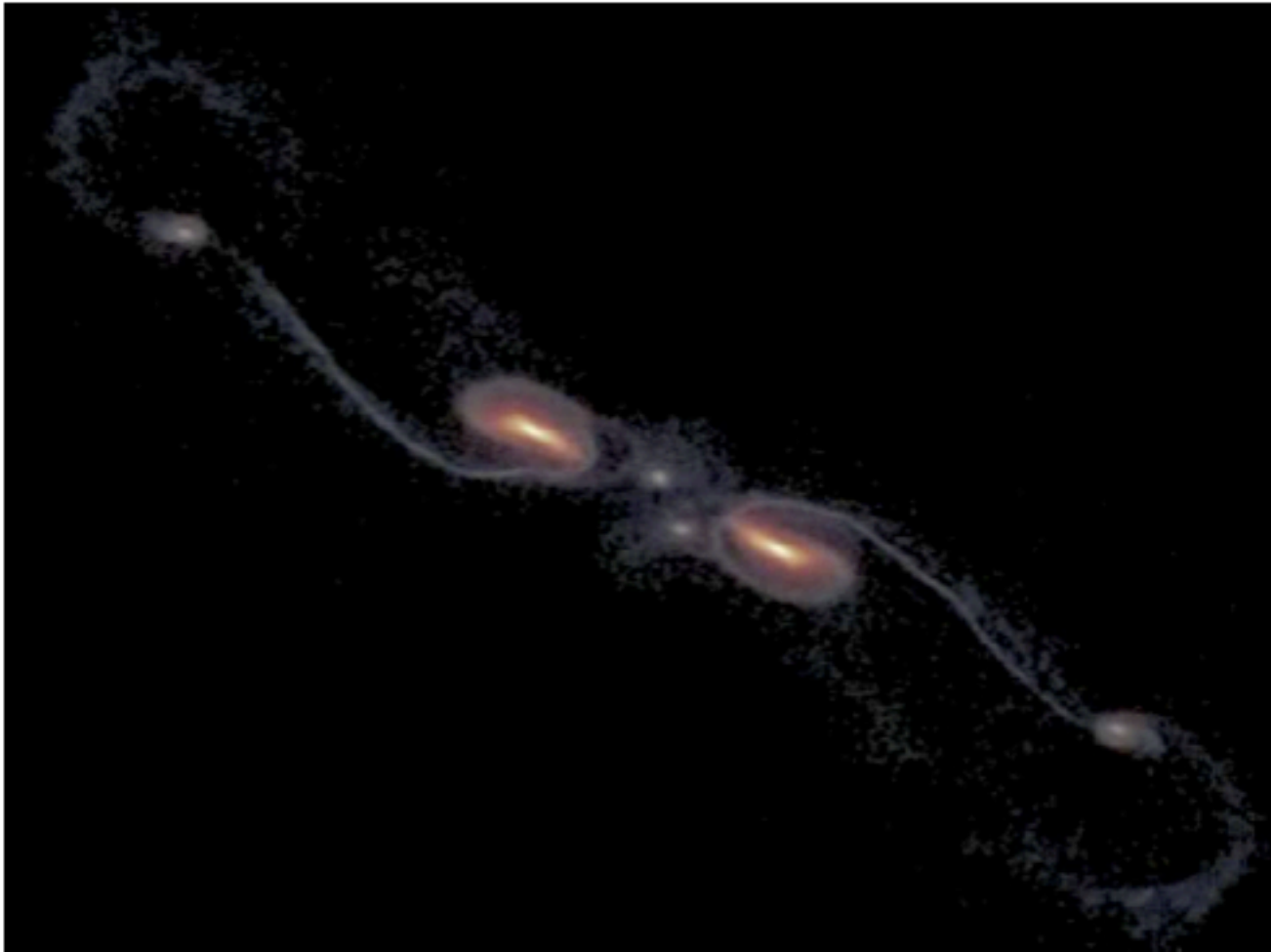


Fig. 6. Tidal dwarf formation at the tip of the tidal tail in MOND.

How else can we tell the difference?

McGaugh (1999) proposed a test for the existence of non-baryonic Cold Dark Matter in the power spectrum of the cosmic microwave background.

1st:2nd peak ratio larger w/o CDM

$$A_{1:2} < 1.9 \text{ (CDM)}$$

$$A_{1:2} = 2.4 \text{ (No CDM)}$$

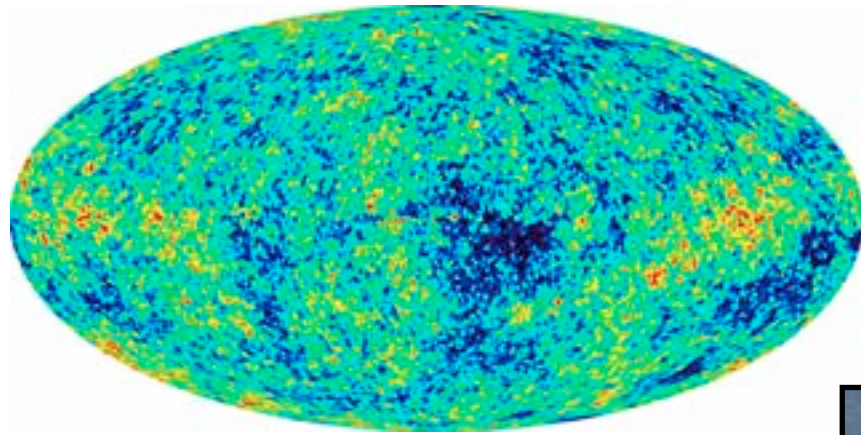
3rd peak lower than 2nd w/o CDM

$$A_{2:3} < 1 \text{ (CDM)}$$

$$A_{2:3} > 1 \text{ (No CDM)}$$

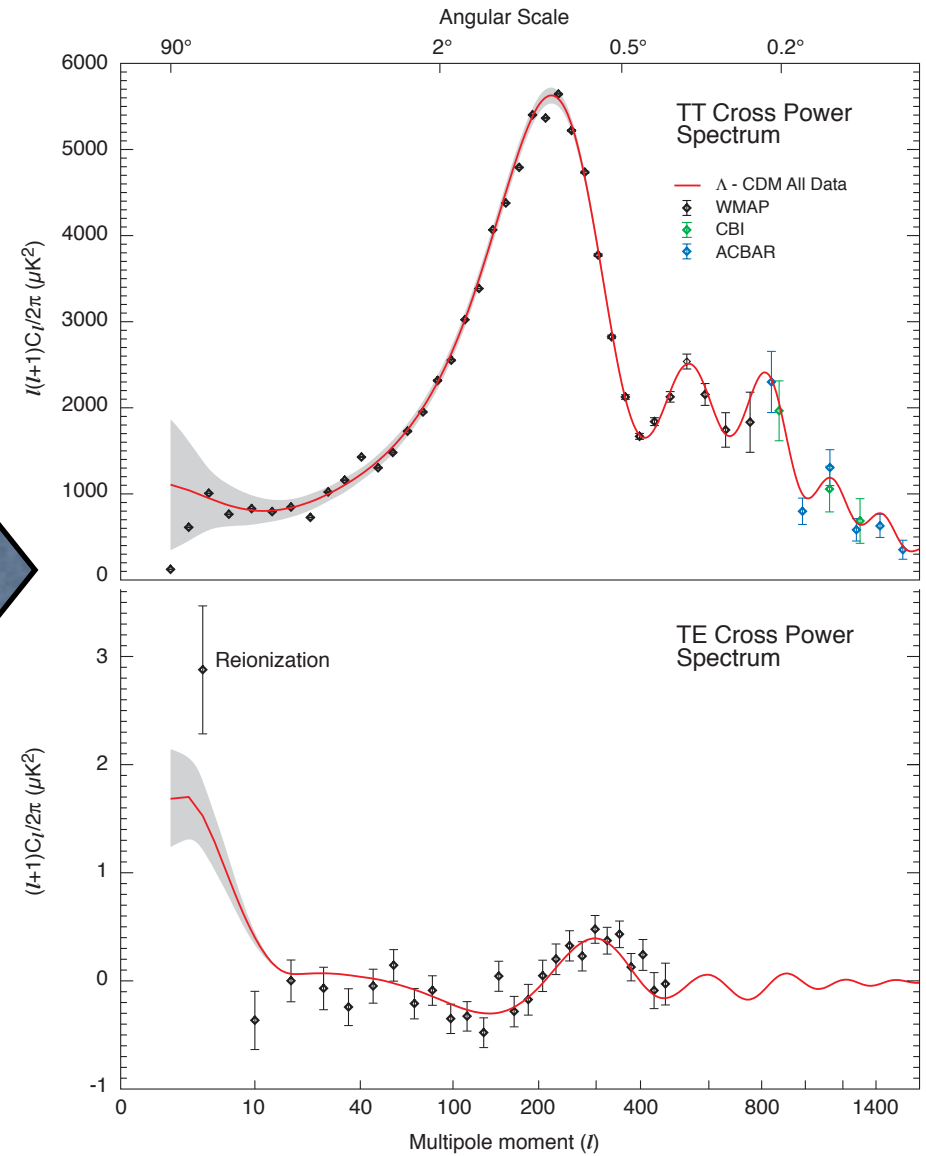
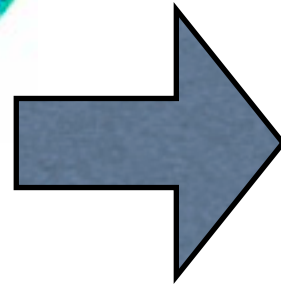
Also expect:

- enhanced ISW
- earlier reionization

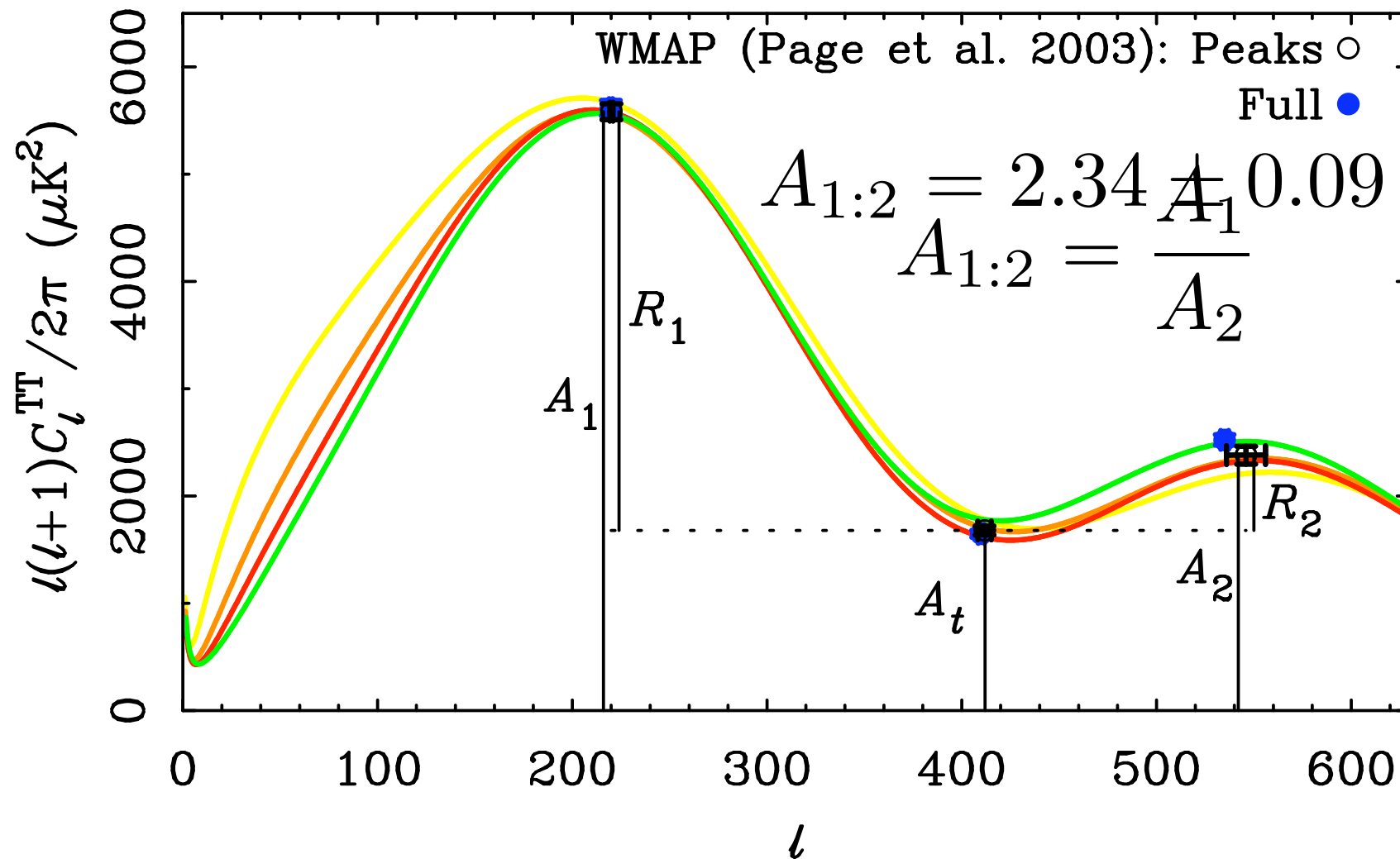


WMAP

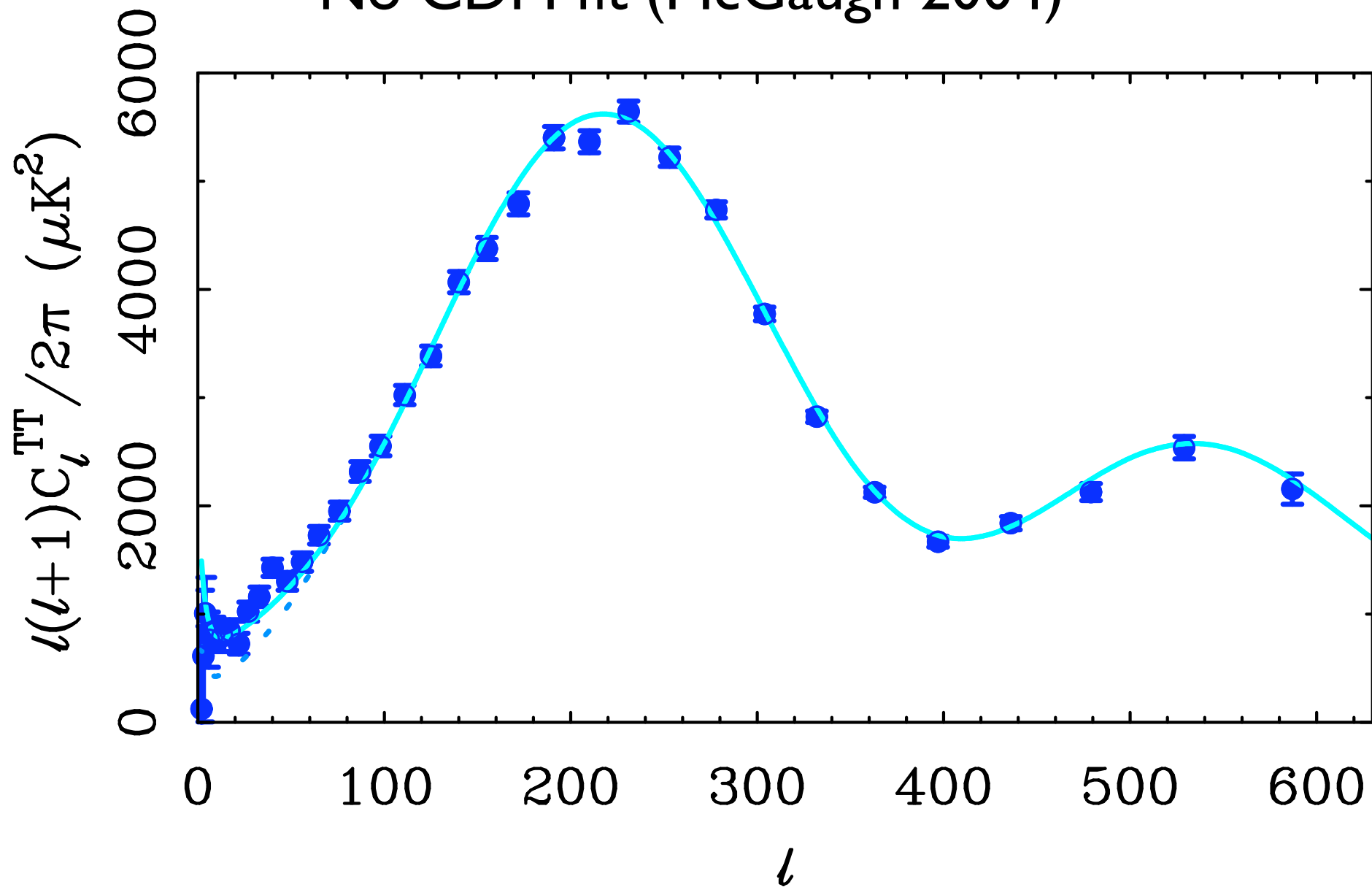
Cosmic Background Radiation



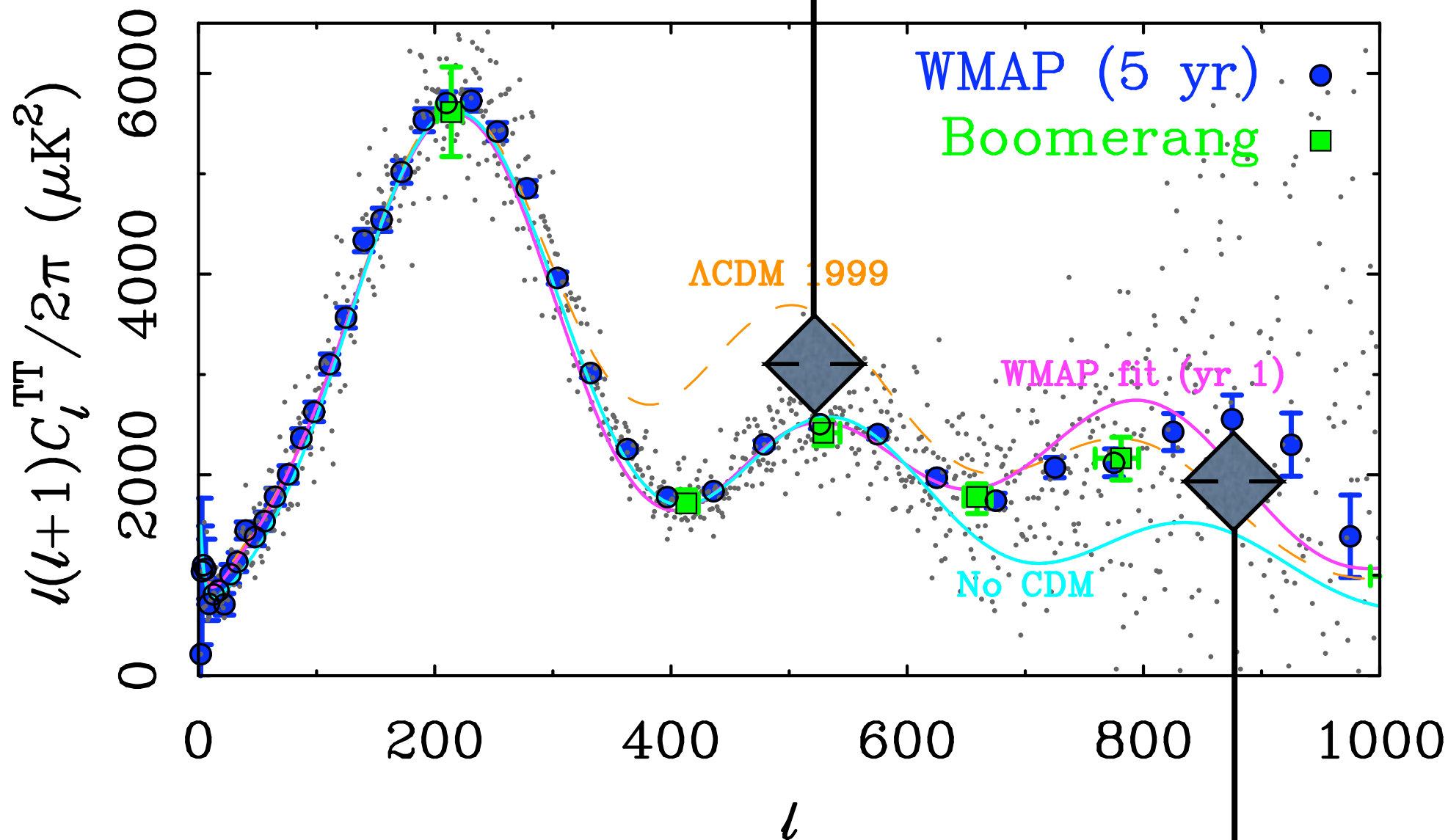
No CDM prediction (McGaugh 1999): $A_{1:2} = 2.4$



No CDM fit (McGaugh 2004)



Depends on baryon density



Driving term from TeVeS? (Ferrara & Skordis 2005)

Big Bang Nucleosynthesis

