



Empirical Tests of MOND in Galaxies

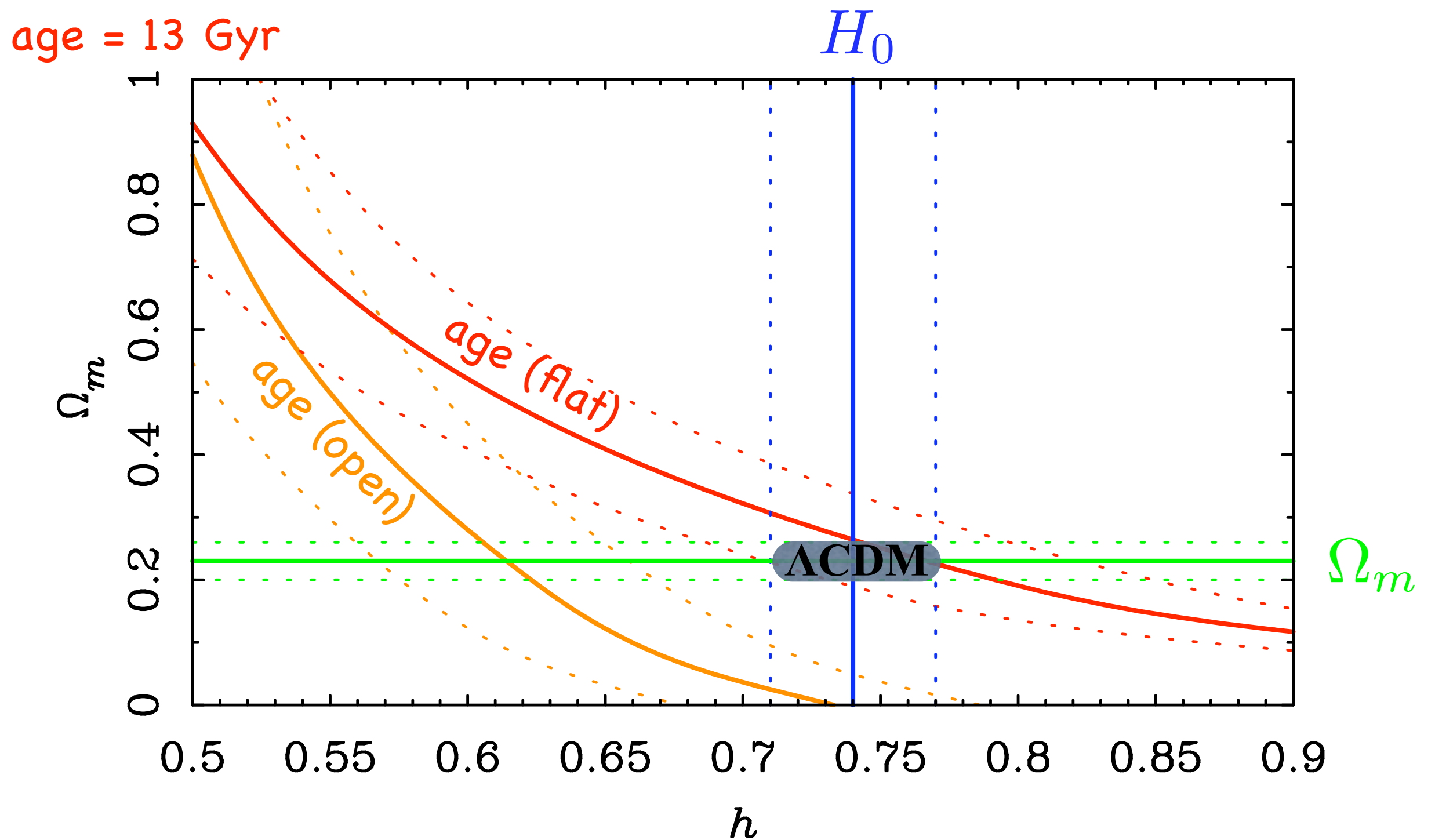
Stacy McGaugh



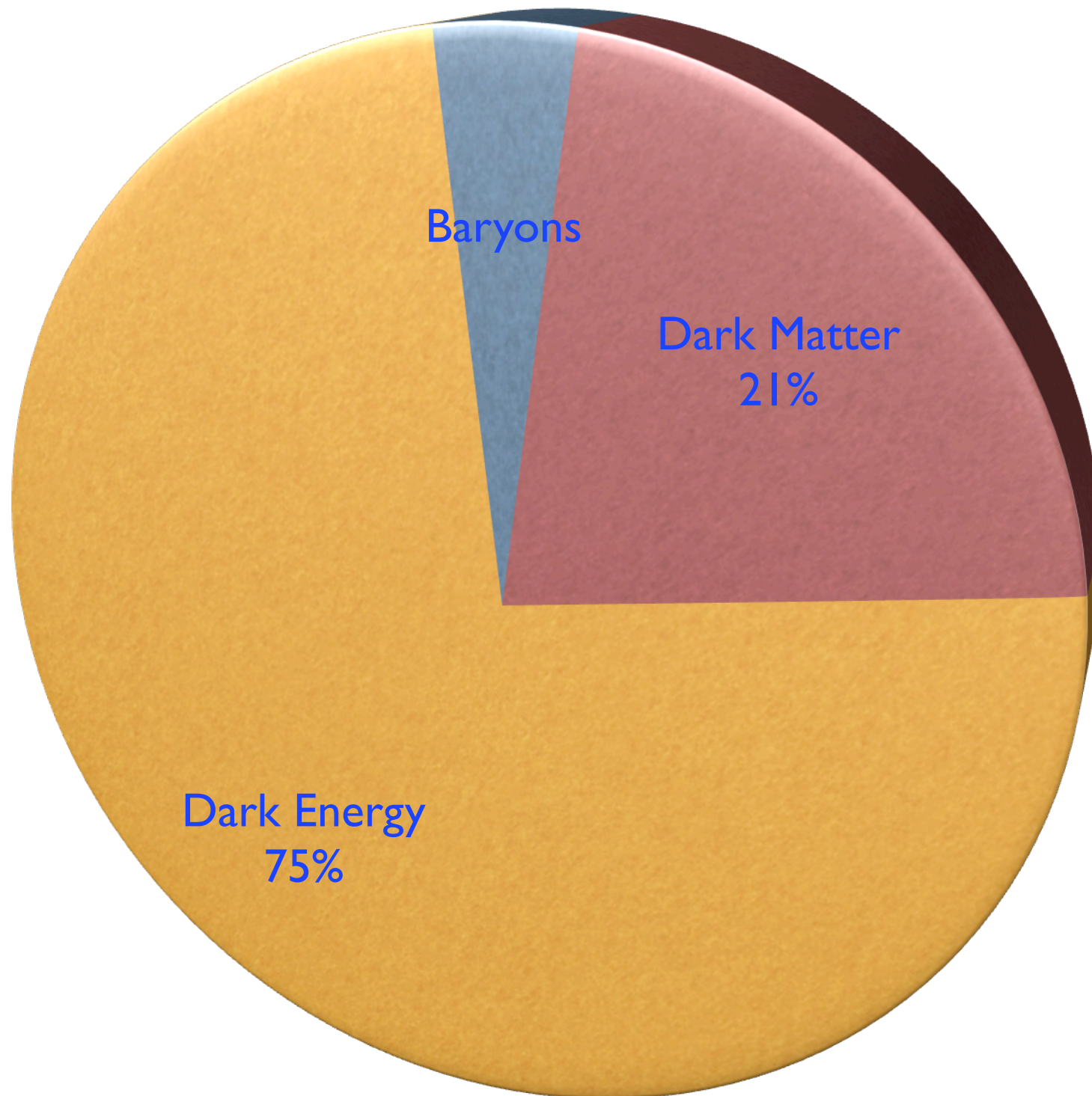
University of Maryland



Cosmological Constraints



SN Ia and CMB want similar numbers, so it must be true!

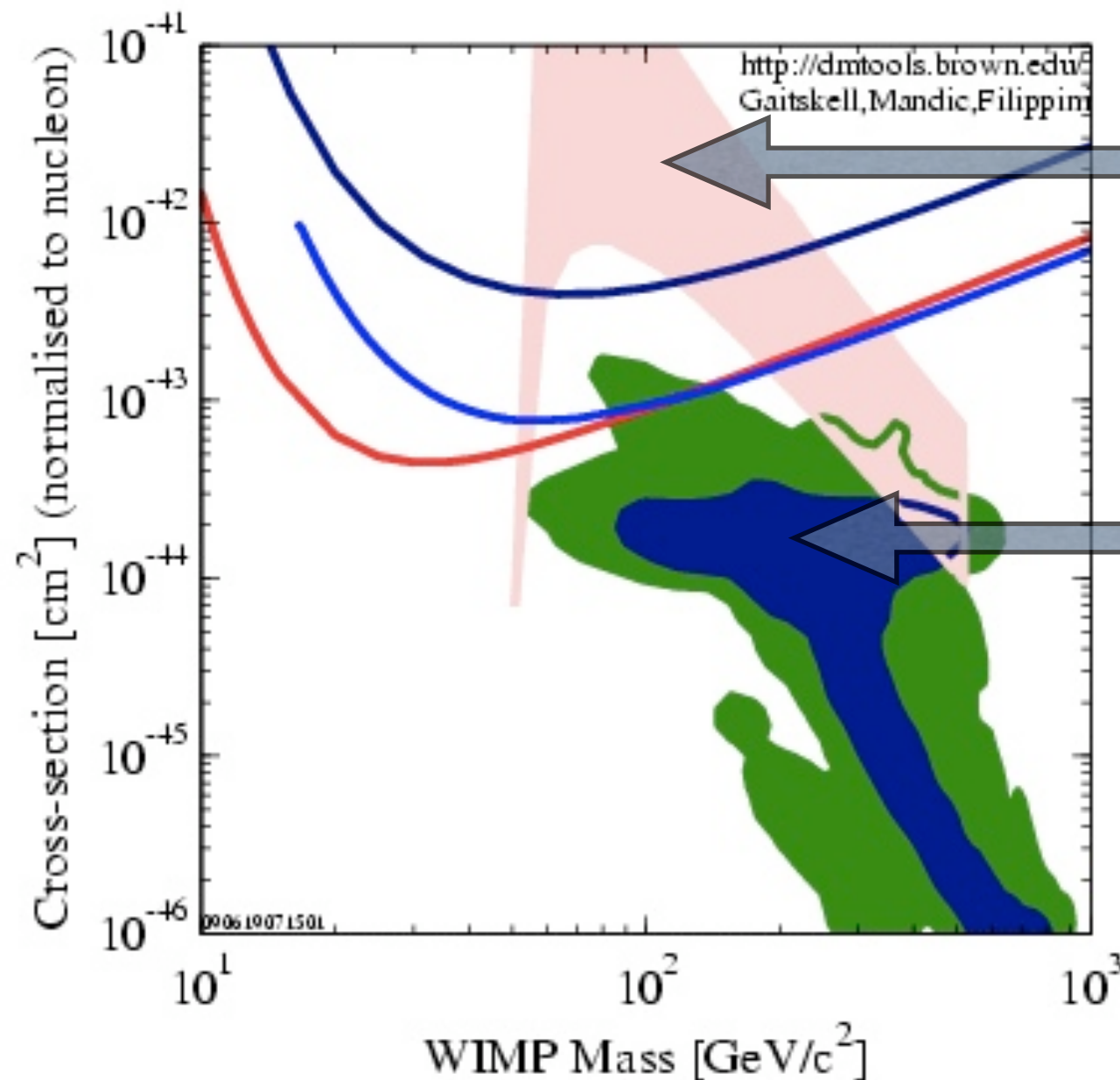


FLRW cosmology
only works with

- dark matter
- dark energy

We don't know what
dark matter is and
we don't understand
what dark energy means

Does dark matter exist?



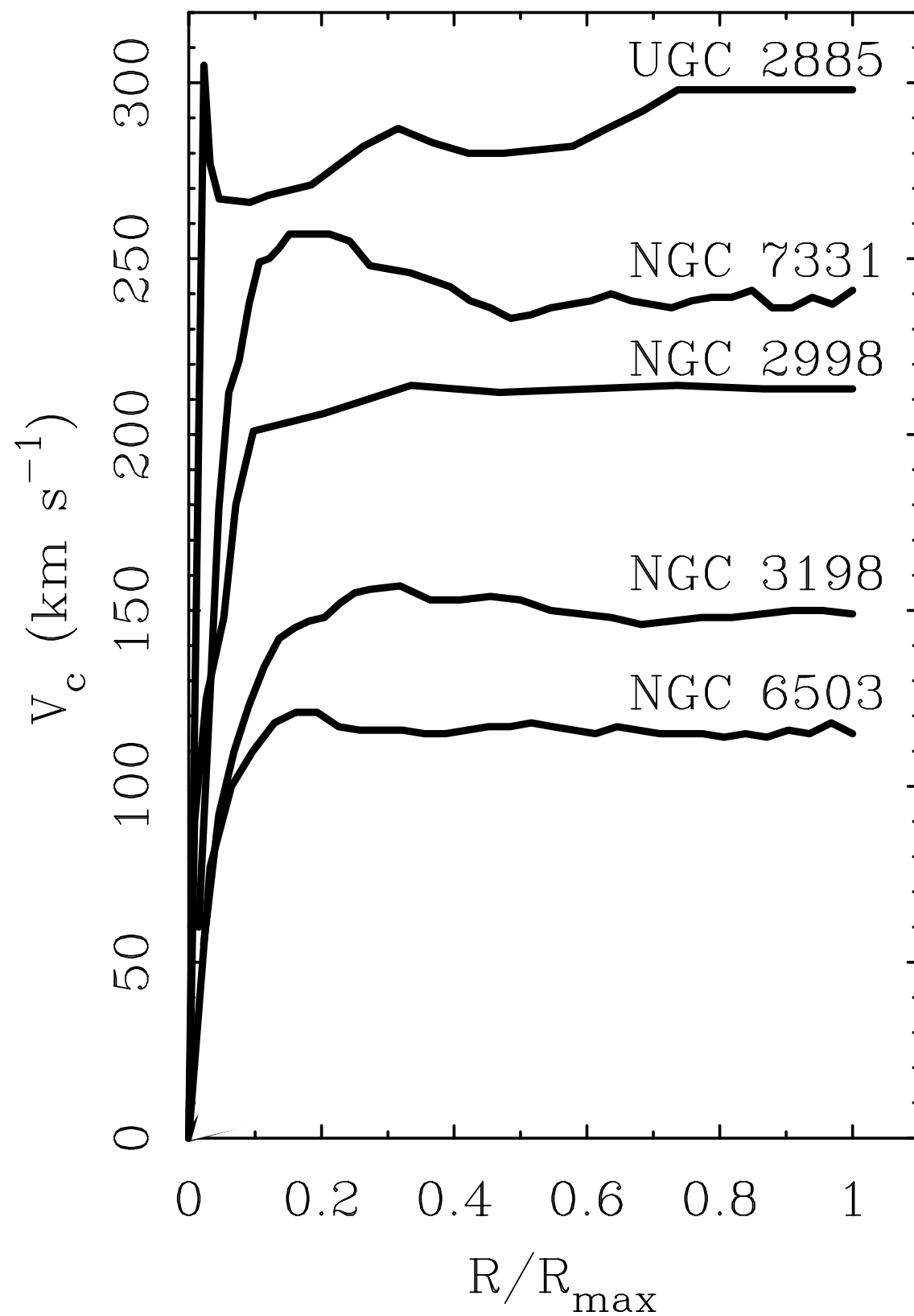
original prediction

current expectation

But we behave like
we're pretty darn sure
that dark matter is
made of WIMPs.

DATA listed top to bottom on plot
CDMS (Soudan) 2004 Blind 53 raw kg-days Ge
ZEPLIN III (Dec 2008) result
XENON10 2007 (Net 136 kg-d)
Ellis et al., Spin dep. sigma in CMSSM
Trotta et al 2008, CMSSM Bayesian: 68% contour
Trotta et al 2008, CMSSM Bayesian: 95% contour
0906.1907 [501]

Rotation curves of spirals



Interpretation in terms
of dark matter leads to
fine-tuning problems.

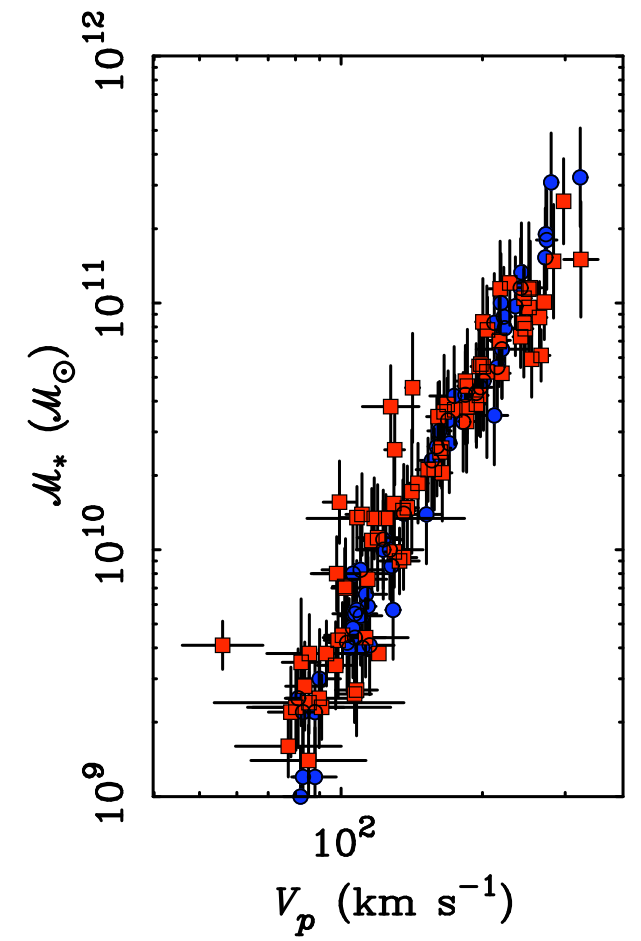
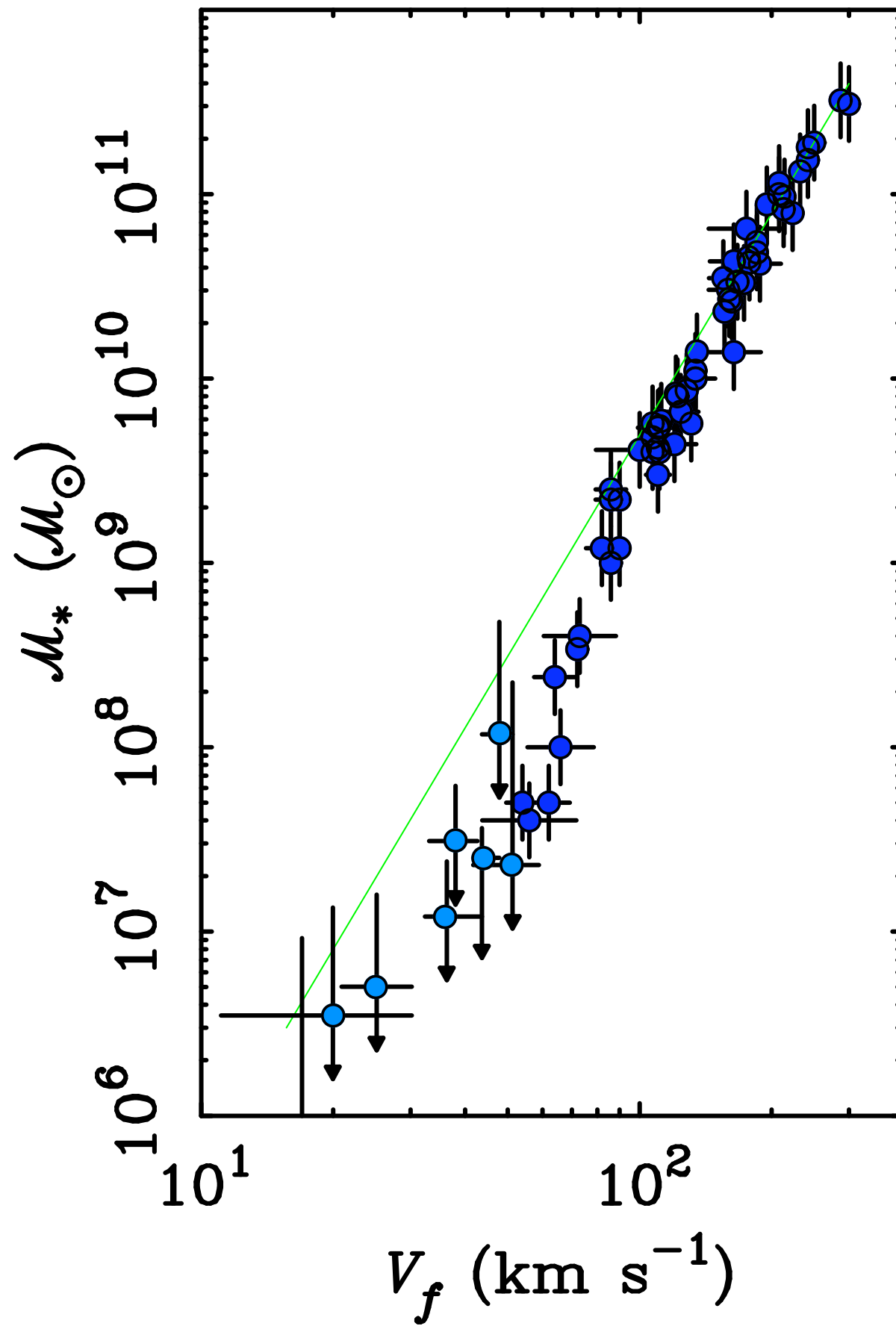


Global Relation:

Tully-Fisher Relation

Stars only

$$M_{\star} = \Upsilon_{\star} L$$



● McGaugh (2005)

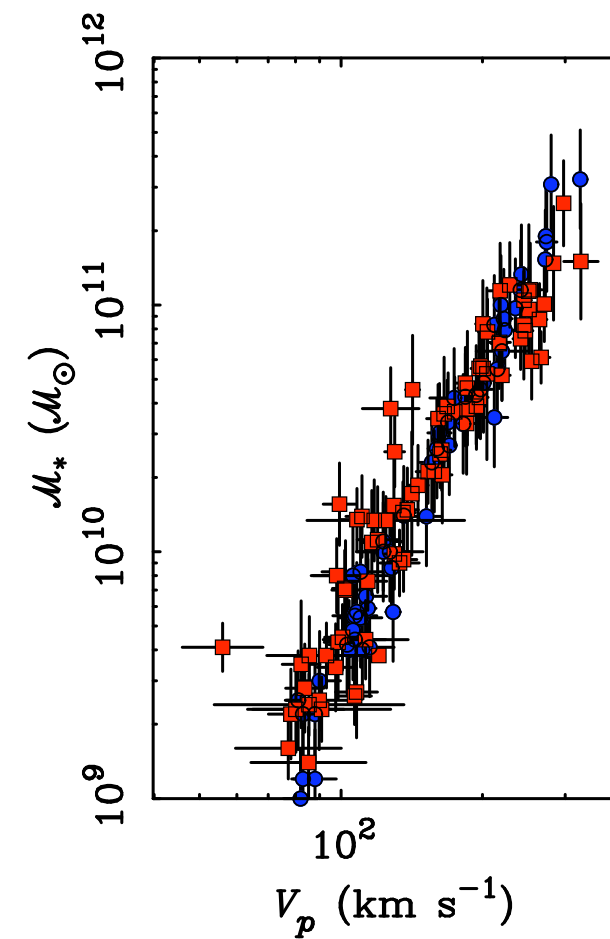
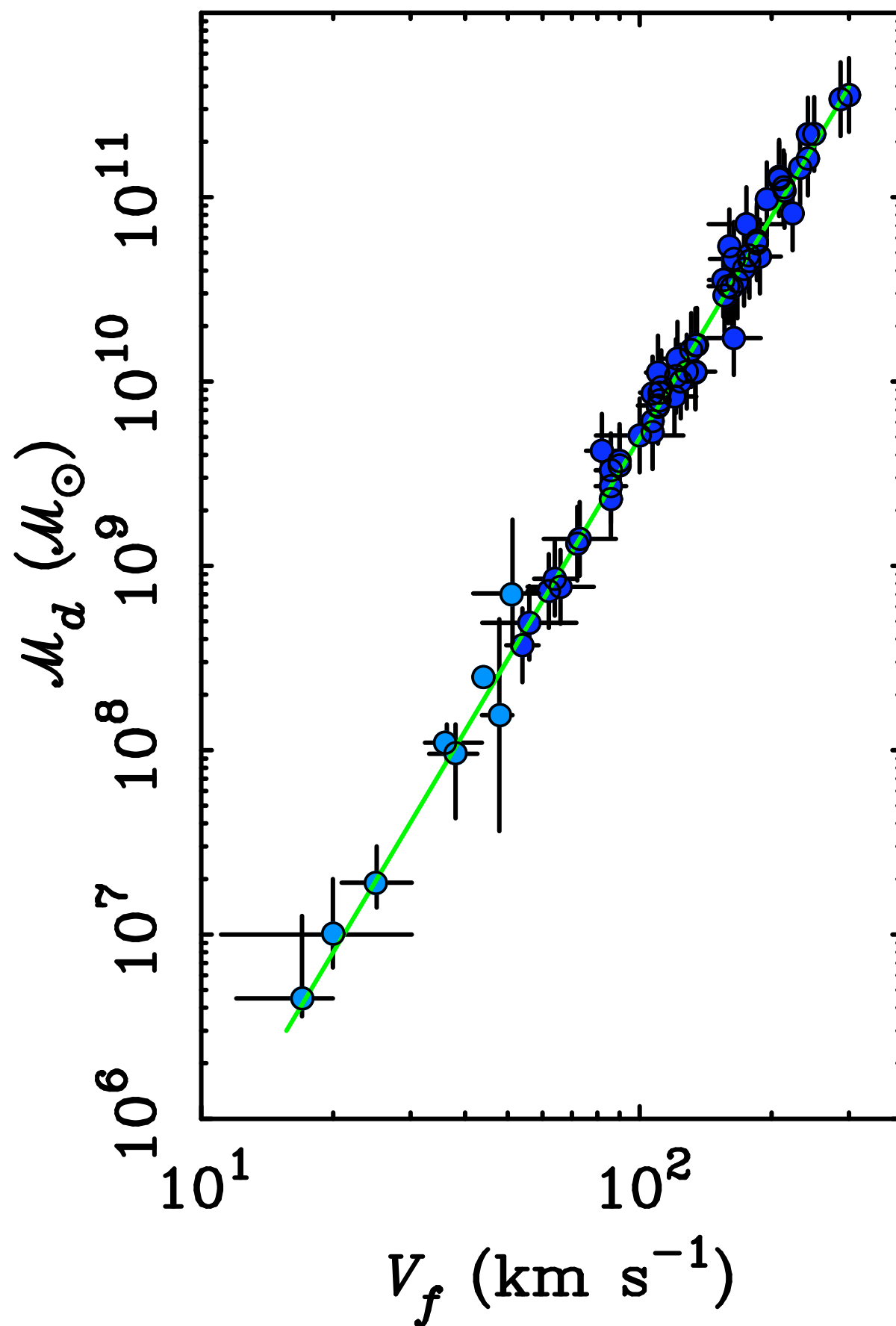
■ Pizagno *et al.*

Global Relation:

Baryonic Tully-Fisher

Stars plus gas

$$M_b = M_\star + M_{gas}$$



line:

$$\log M_b = 4 \log V_f + 1.7$$

(McGaugh 2005)

Implies no other substantial
reservoirs of baryonic mass.

NGC 2403



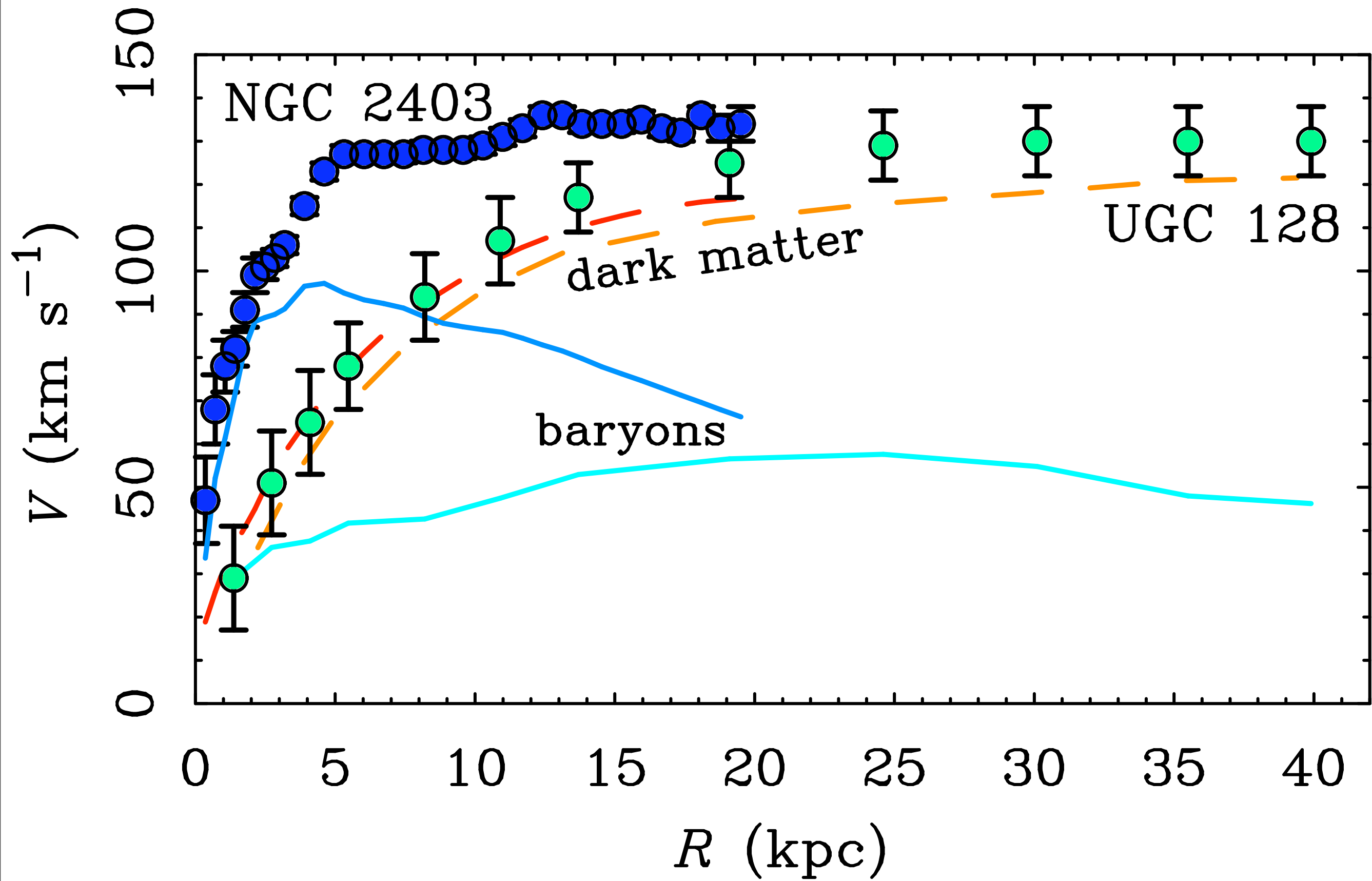
HSB

UGC 128
LSB

Same global M_b, V

Very different
mass distributions

de Blok & McGaugh (1996)
Tully & Verheijen (1997)



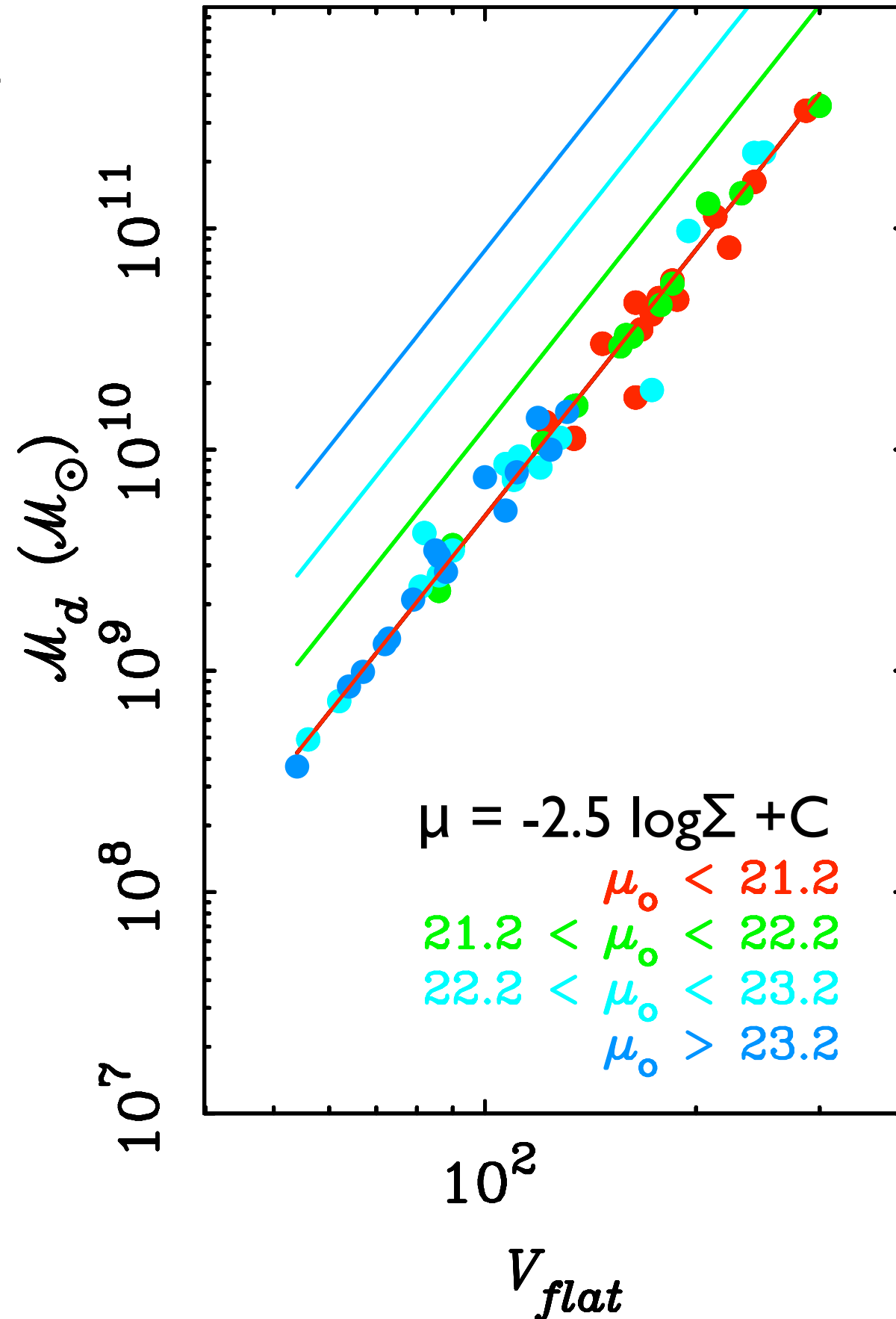
Newton says

$$V^2 = GM/R.$$

Equivalently,

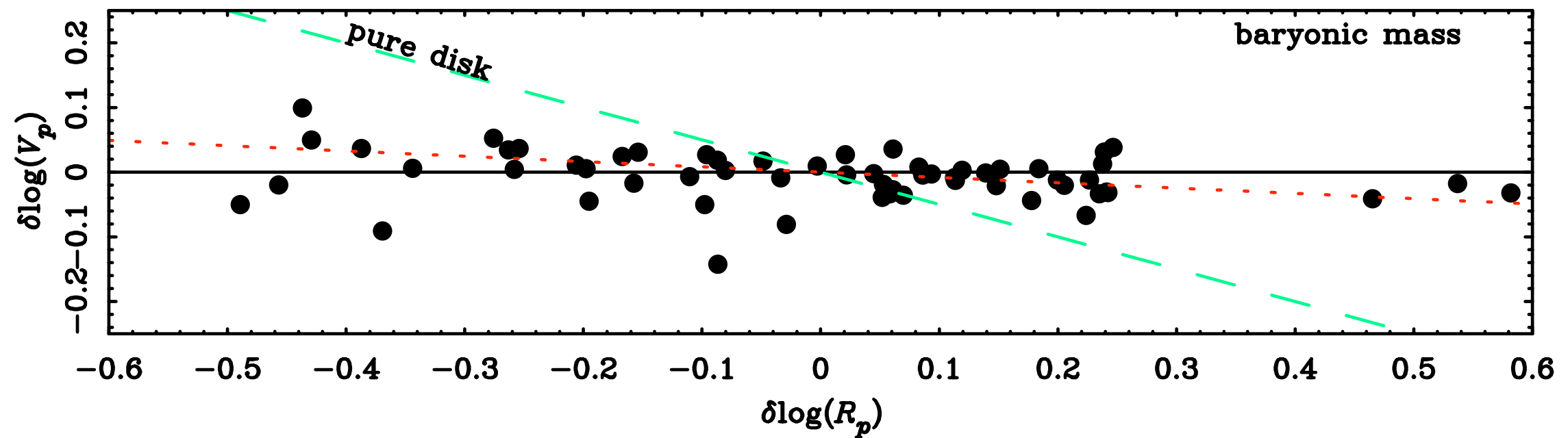
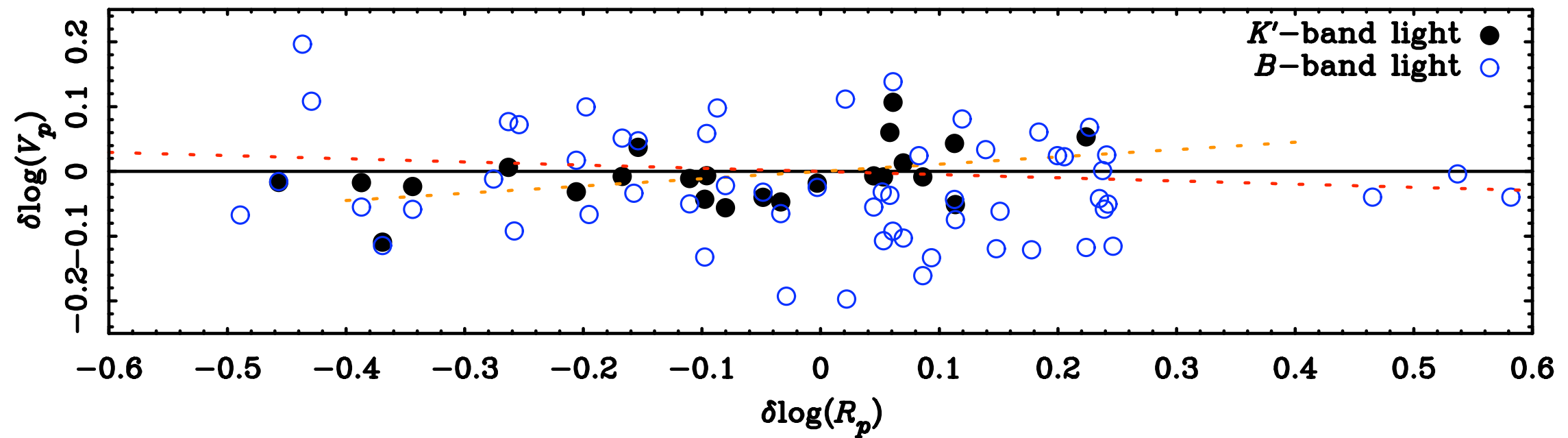
$$\Sigma = M/R^2$$

$$V^4 = G^2 M \Sigma$$



Therefore
Different Σ
should mean
different TF
normalization.

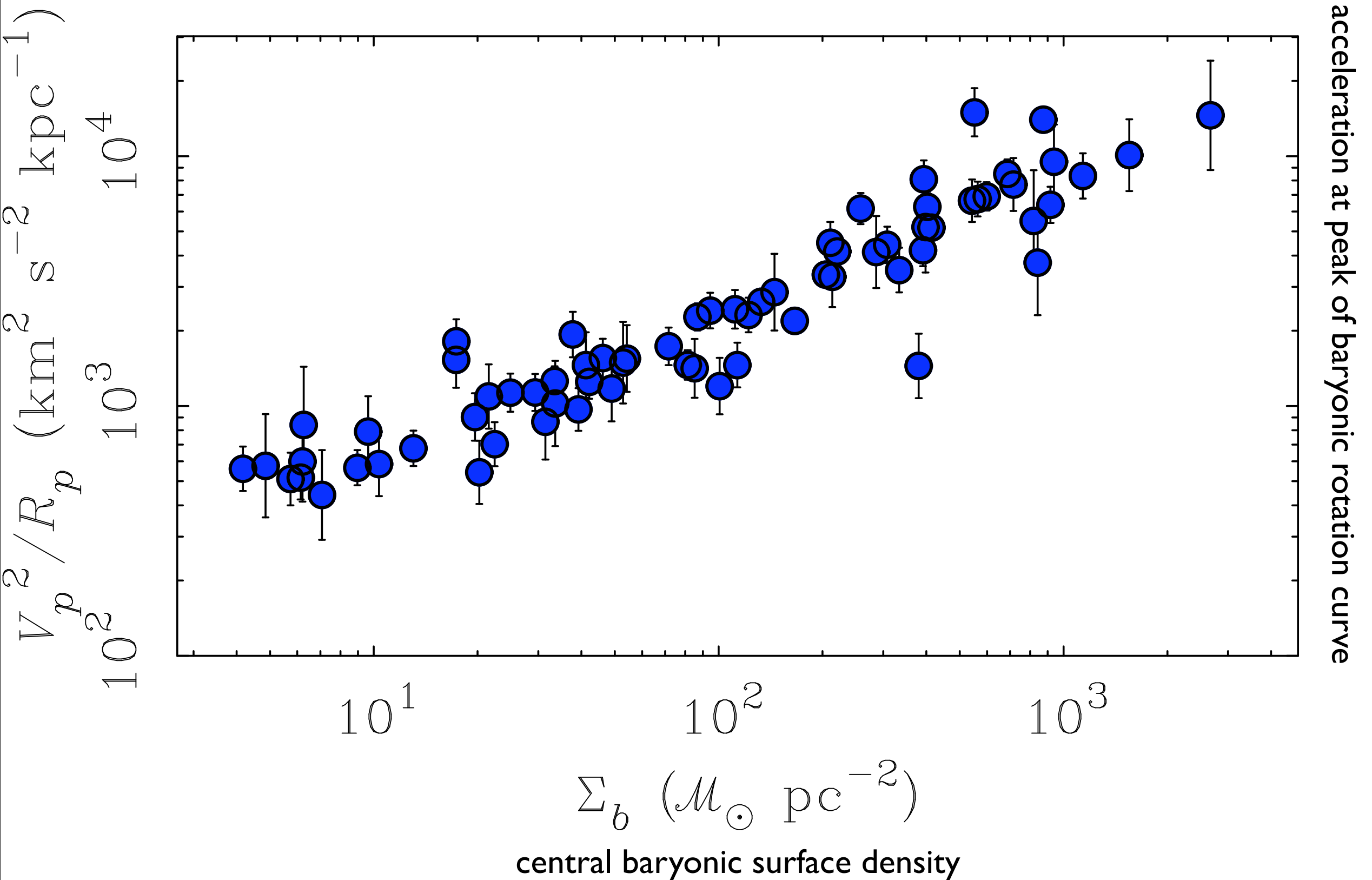
No Residuals from TF rel'n



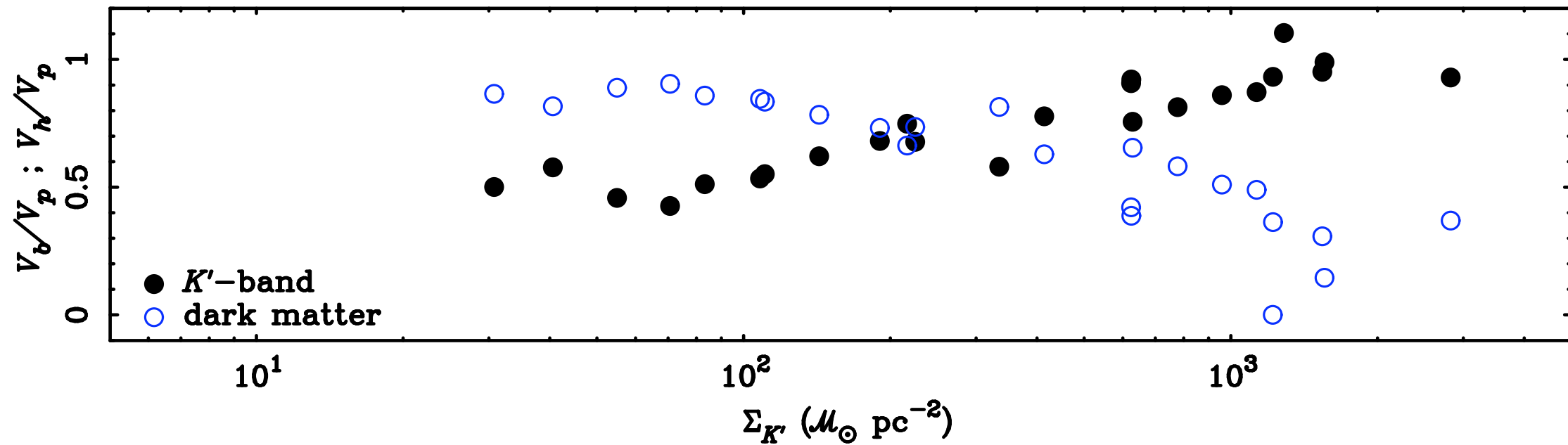
Sometimes interpreted to mean that dark matter dominates over disk mass

Acceleration related to baryonic surface density

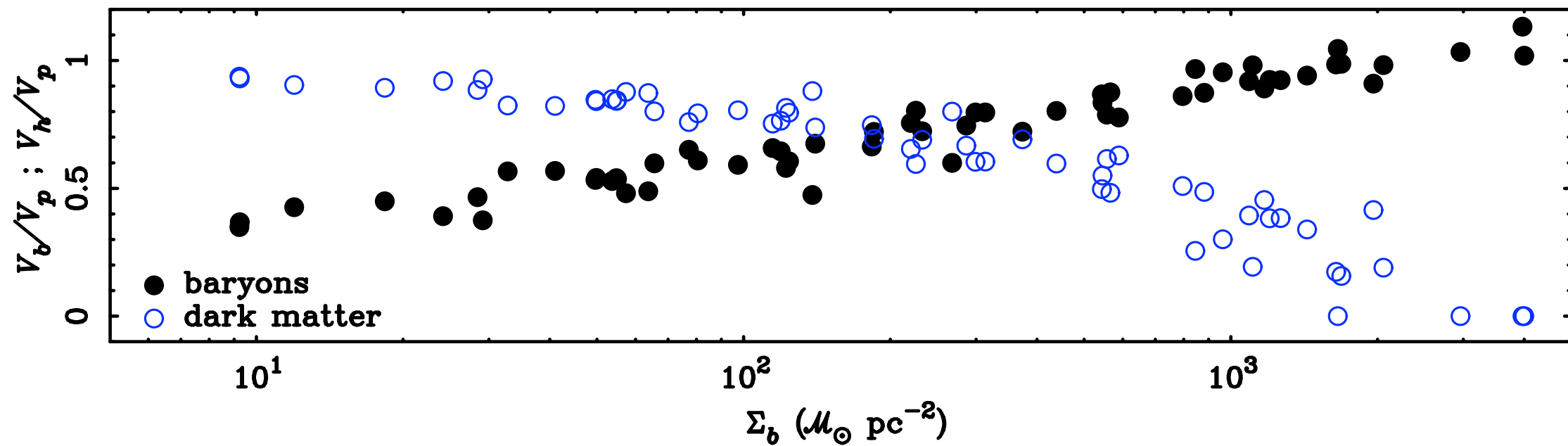
Baryons important to dynamics - dark matter does not dominate.
A contradiction to purely Newtonian dynamics?



Fine-tuning unavoidable



Phys. Rev. Lett. 95, 171302 (2005)



“...working on the thing can drive you mad.”

MOND

$$a \gg a_0$$

$$a \rightarrow g_N$$

$$a \ll a_0$$

$$a \rightarrow \sqrt{g_N a_0}$$

$$a_0 \approx 1 \text{ Å s}^{-2}$$

$$\mu \left(\frac{a}{a_0} \right) = \frac{g_N}{a}$$

$$\mu \rightarrow 1$$

$$a \gg a_0$$

$$\mu \rightarrow \frac{a}{a_0}$$

$$a \ll a_0$$

Milgrom 1983

No. 2, 1983

MODIFICATION OF NEWTONIAN DYNAMICS

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A major step in understanding ellipticals can be made if we can identify them, at least approximately, with idealized structures such as the FRCL spheres discussed above. I have also studied isotropic and nonisotropic isothermal spheres, in the modified dynamics, as such possible structures. I found that they have properties which resemble those of ellipticals and galactic bulges. I describe them in Milgrom (1983).

VIII. PREDICTIONS

The main predictions concerning mass are as follows.

1. Velocity curves calculated with the modified dynamics on the basis of the observed mass in galaxies should agree with the observed curves. Elliptical and S0 galaxies may be the best for this purpose since (a) practically no uncertainty due to obscuration is involved and (b) there is not much uncertainty due to the possible presence of molecular hydrogen.

2. The relation between the asymptotic velocity (V_∞) and the mass of the galaxy (M) ($V_\infty^4 = MG a_0$) is an absolute one.

3. Analysis of the π -dynamics in disk galaxies using the modified dynamics should yield surface densities which agree with the observed ones. Accordingly, the same analysis using the conventional dynamics should yield a discrepancy which increases with radius in a predictable manner.

4. Effects of the modified dynamics are predicted to be particularly strong in dwarf elliptical galaxies (for review of properties see, e.g., Hodge 1971 and Zinn 1980). For example, those dwarfs believed to be bound to our Galaxy would have internal accelerations typically of order $a_{in} \sim a_0/30$. Their (modified) acceleration, g , in the field of the Galaxy is larger than the internal ones but still much smaller than a_0 , $g \approx (8 \text{ kpc}/d) a_0$, based on a value of $V_\infty = 220 \text{ km s}^{-1}$ for the Galaxy, and where d is the distance from the dwarf galaxy to the center of the Milky Way ($d \sim 70\text{--}220 \text{ kpc}$). Whichever way the external acceleration turns out to affect the internal dynamics (see the discussion at the end of § II, the section on small groups in Paper III, and Paper I), we predict that when velocity dispersion data is available for the dwarfs, a large mass discrepancy will result when the conventional dynamics is used to determine the masses. The dynamically determined mass is predicted to be larger by a factor of order 10 or more than that which can be accounted for by stars. In case the internal dynamics is determined by the external acceleration, we predict this factor to increase with d and be of order $(d/8 \text{ kpc})$ (as long as $a_{in} \ll g$, $h_{50} = 1$).

Prediction 1 is a very general one. It is worthwhile listing some of its consequences as separate predictions, numbered 5–7 below (note that, in fact, even prediction 2 is already contained in prediction 1).

5. Measuring local M/L values in disk galaxies (assuming conventional dynamics) should give the following results: In regions of the galaxy where $V^2/r \gg a_0$ the local M/L values should show no indication of hidden mass. At a certain transition radius, local M/L should start to increase rapidly. The transition radius should occur where $V^2/r = a_0$. This is just as the following advantages: (a) does not require an absolute calibration of M/L as we are concerned only with variations of this quantity; (b) Effects of the modified dynamics manifest themselves more clearly in local M/L determinations than in the integrated masses and (c) in many cases this requires information on local behavior in the disk only while the spheroid can be neglected. This makes the determination of mass from velocity more certain.

6. Disk galaxies with low surface brightness provide particularly strong tests (a study of a sample of such galaxies is described by Strom 1982 and by Romanishin *et al.* 1982). As low surface brightness means small accelerations, the effects of the modification should be more noticeable in such galaxies. We predict, for example, that the proportionality factor in the $M \propto V_\infty^4$ relation for these galaxies is the same as for the high surface density galaxies. In contrast, if one wants to obtain a correlation $M \propto V_\infty^4$ in the conventional dynamics (with additional assumptions), one is led to the relation $M \propto \Sigma^{-1} V_\infty^4$ (see, for example, Aaronson, Huchra, and Mould 1979), where Σ is the average surface brightness. This implies that low surface density galaxies, of a given velocity, have a mass higher than predicted by the M - V relation derived for normal surface density galaxies.

We also predict that the lower the average surface density of a galaxy is, the smaller is the transition radius, defined in prediction 5, in units of the galaxy's scale length. In fact, if the average surface density is very small we may have a galaxy in which $V^2/r < a_0$ everywhere, and analysis with conventional dynamics should yield local M/L values starting to increase from very small radii.

7. As the study of model rotation curves shows, we predict a correlation between the value of the average surface density (or brightness) of a galaxy and the steepness with which the rotational velocity rises to its asymptotic value (as measured, for example, by the radius at which $V = V_\infty/2$ in units of the scale length of the disk). Small surface densities imply slow rise of V .

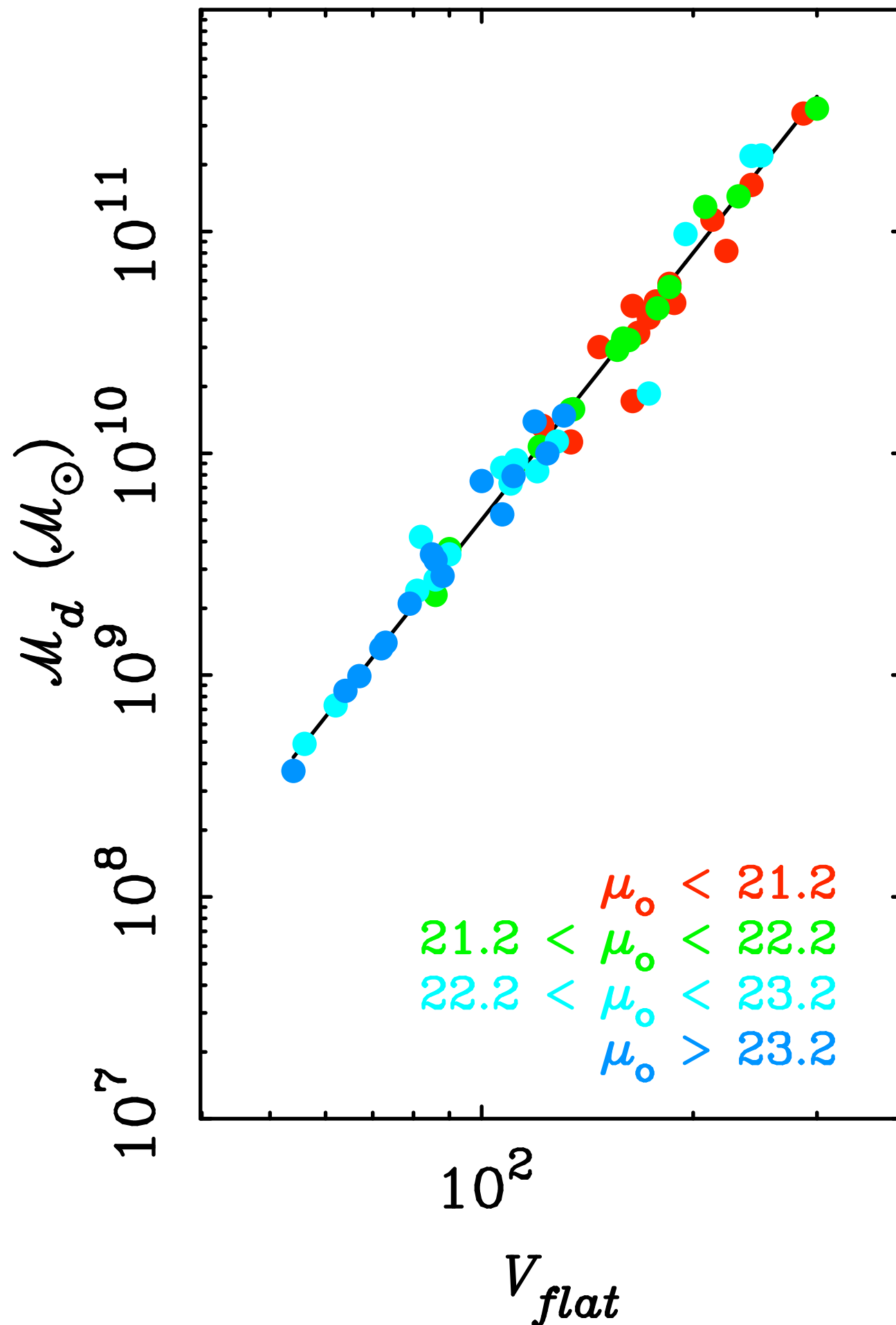
IX. DISCUSSION

The main results of this paper can be summarized by the statement that the modified dynamics eliminates the need to assume hidden mass in galaxies. The effects in galaxies which I have considered, and which are commonly attributed to such hidden mass, are readily explained by the modification. More specifically:

MOND predictions

- The Tully-Fisher Relation
- slope = 4
- Normalization = $1/(a_0 G)$
- Fundamentally a relation between Disk Mass and V_{flat}
- No Dependence on Surface Brightness
- Dependence of conventional M/L on radius and surface brightness
- Rotation Curve Shapes
- Surface Density \sim Surface Brightness
- Detailed Rotation Curve Fits
- Stellar Population Mass-to-Light Ratios

“Disk Galaxies with low surface brightness provide particularly strong tests”



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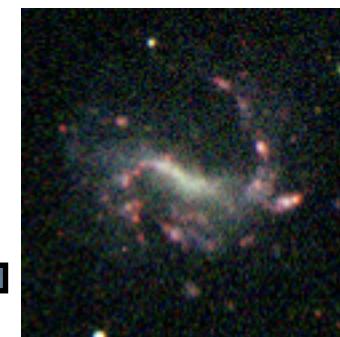
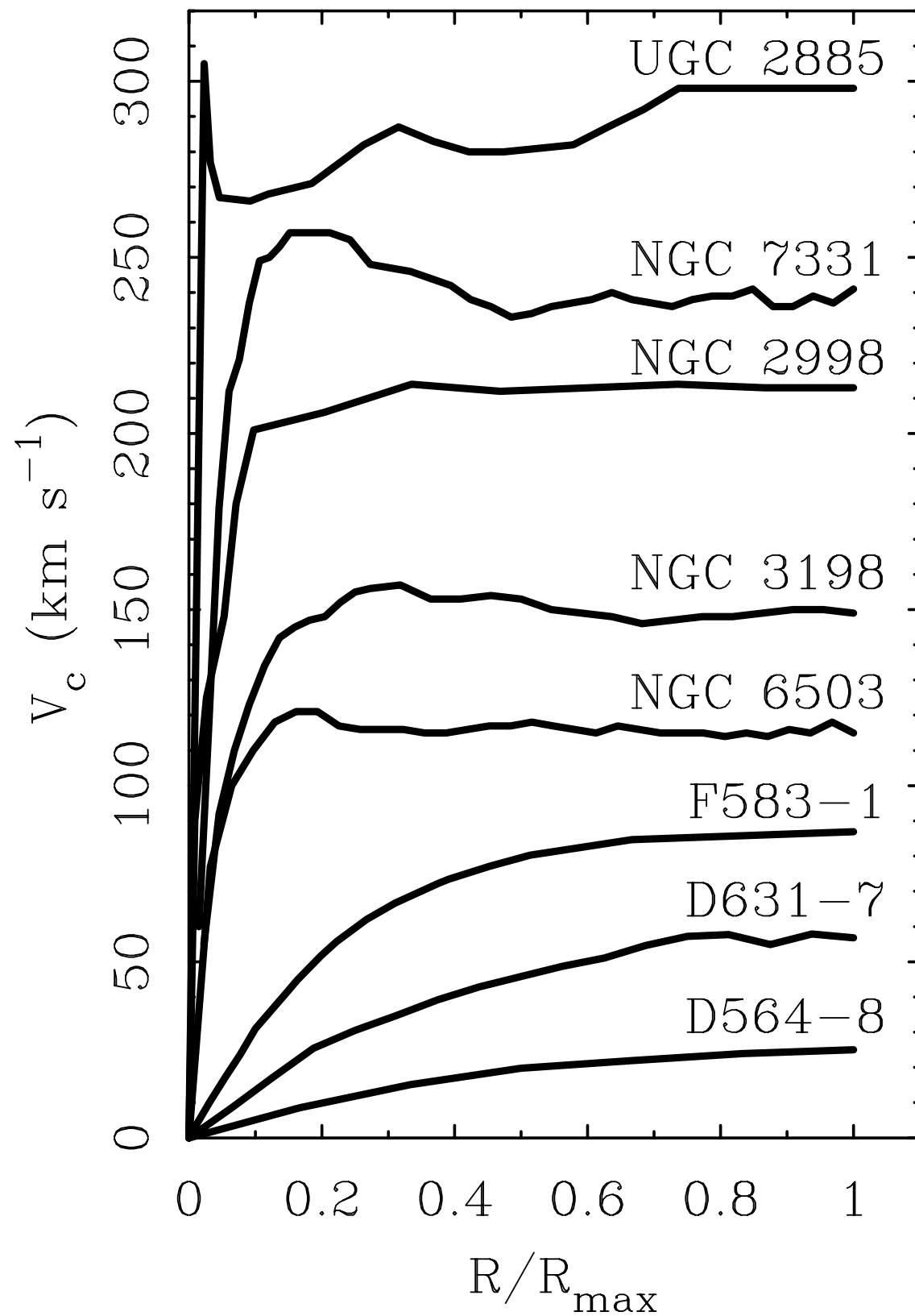
- Rotation Curve Shapes

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Rotation curves of spirals
and low mass dIrrs with $M_* < M_g$.

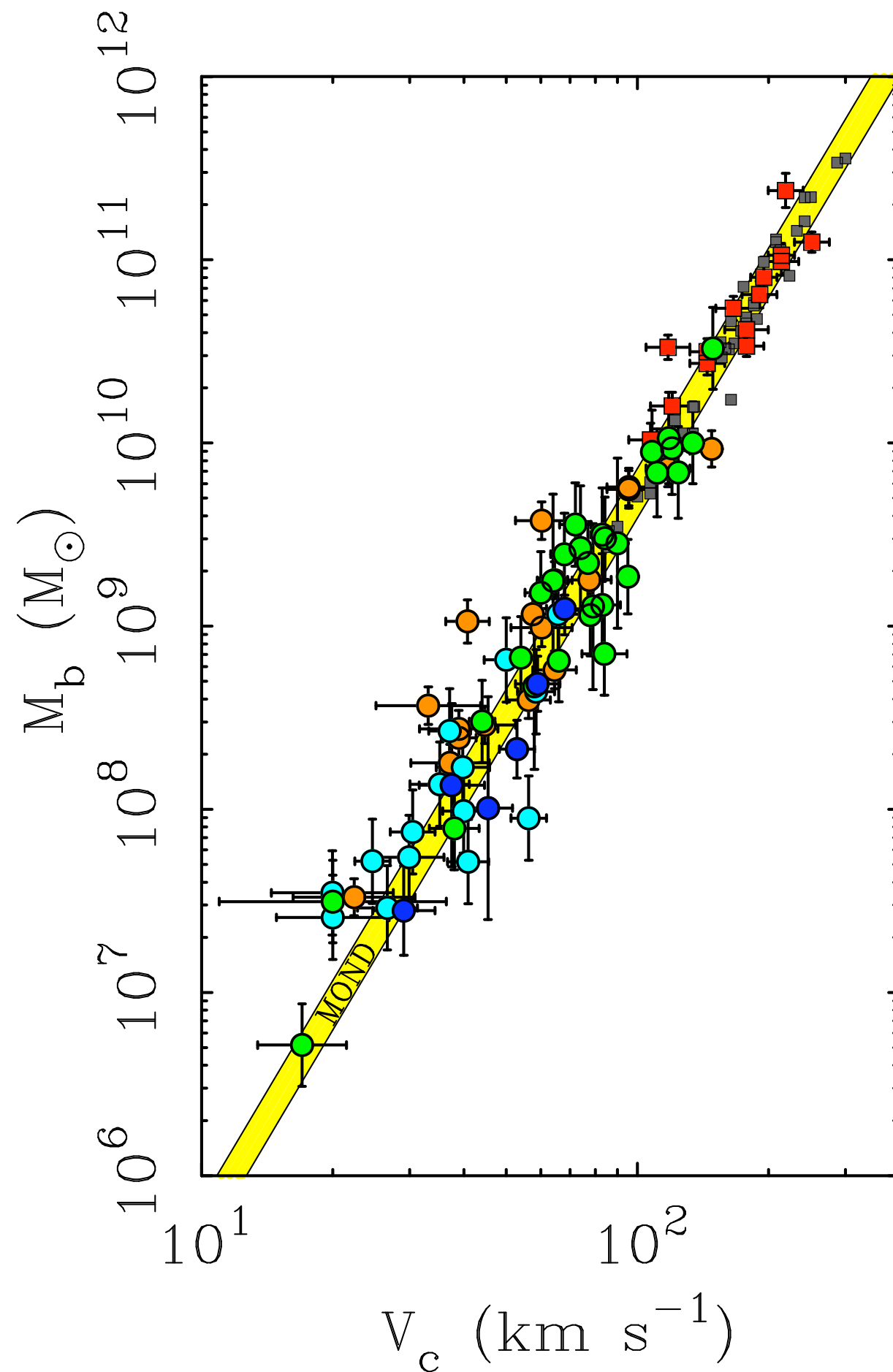


Rotation curves of
late type disks
(Sd, Sm, Irr)



Kuzio de Naray et al.
(2006, 2008, 2009);
Trachternach et al.
(2009)

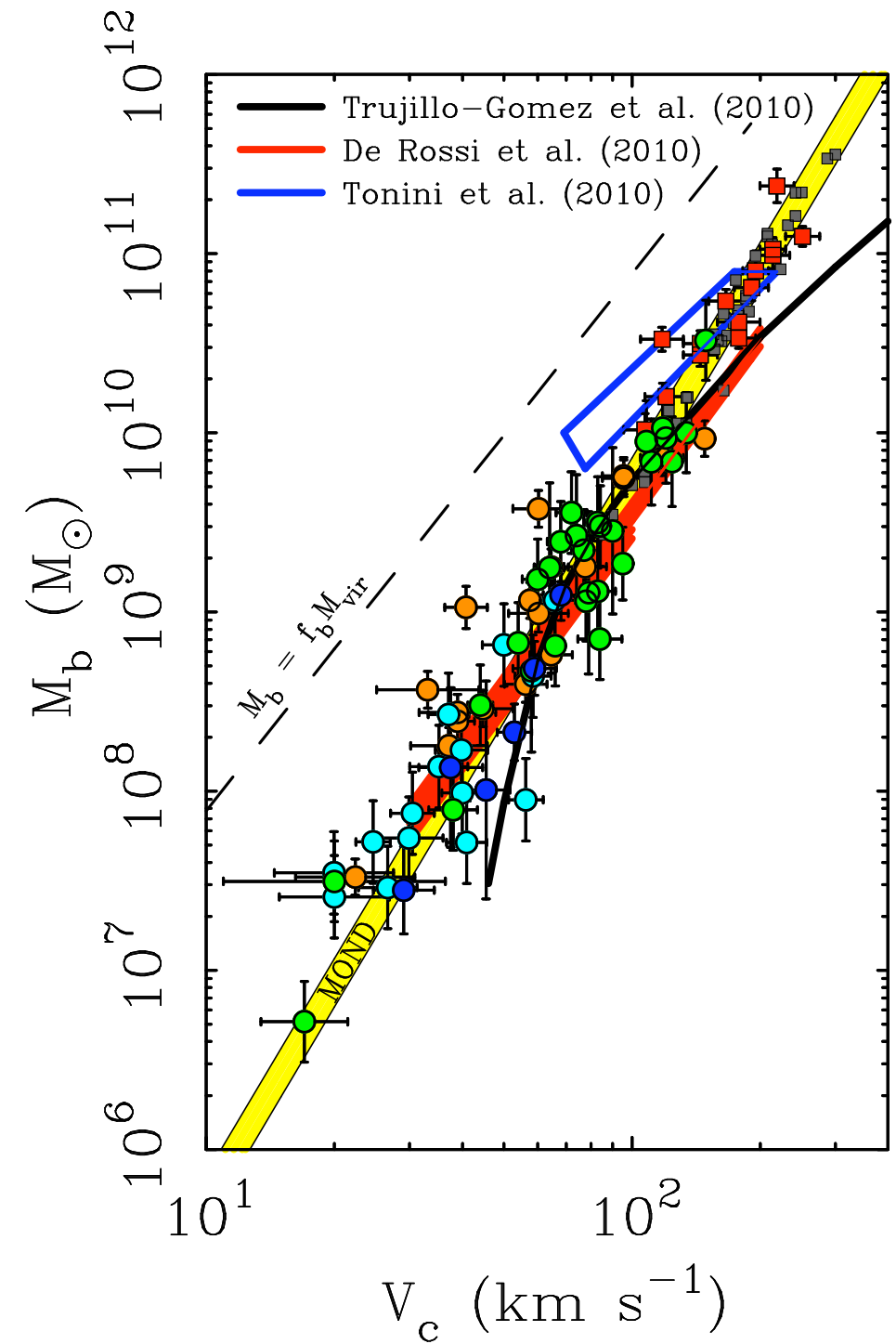


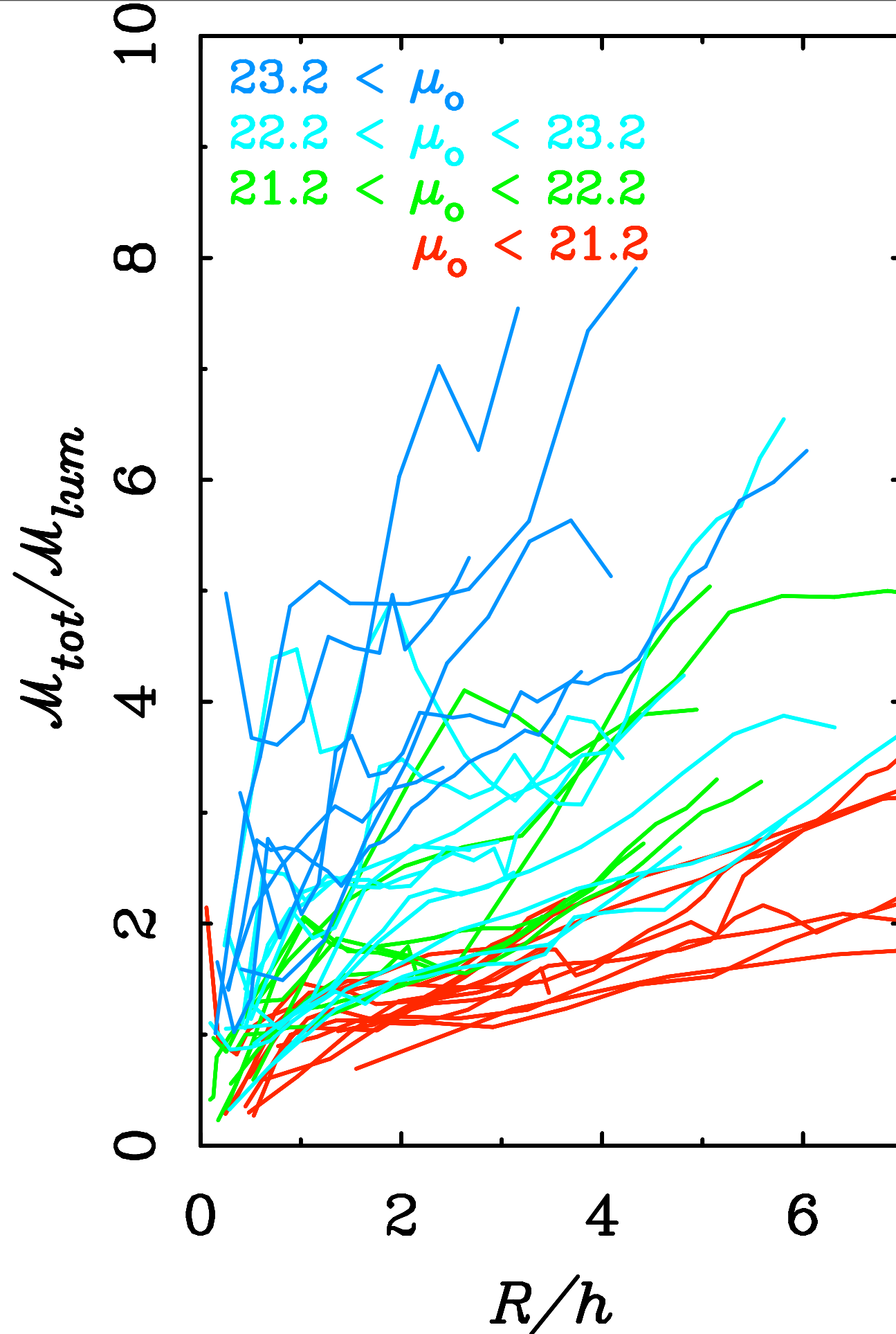


- $M_* > M_g$ (MOND fits)
McGaugh (2005)
- $M_* > M_g$ (H-band ppsynth)
Sakai (2000); Gurovich et al. (2010)
- $M_* < M_g$ ($V_c = W_{20}/2$)
Gurovich et al. (2010)
- $M_* < M_g$ $\sin(i_{opt}) < 1.12 \sin(i_{HI})$
Begum et al. (2008)
- $M_* < M_g$
Stark et al. (2009)
- $M_* < M_g$
Trachternach et al. (2008)

Position on BTFR independent
of stellar M_*/L for $M_* < M_g$

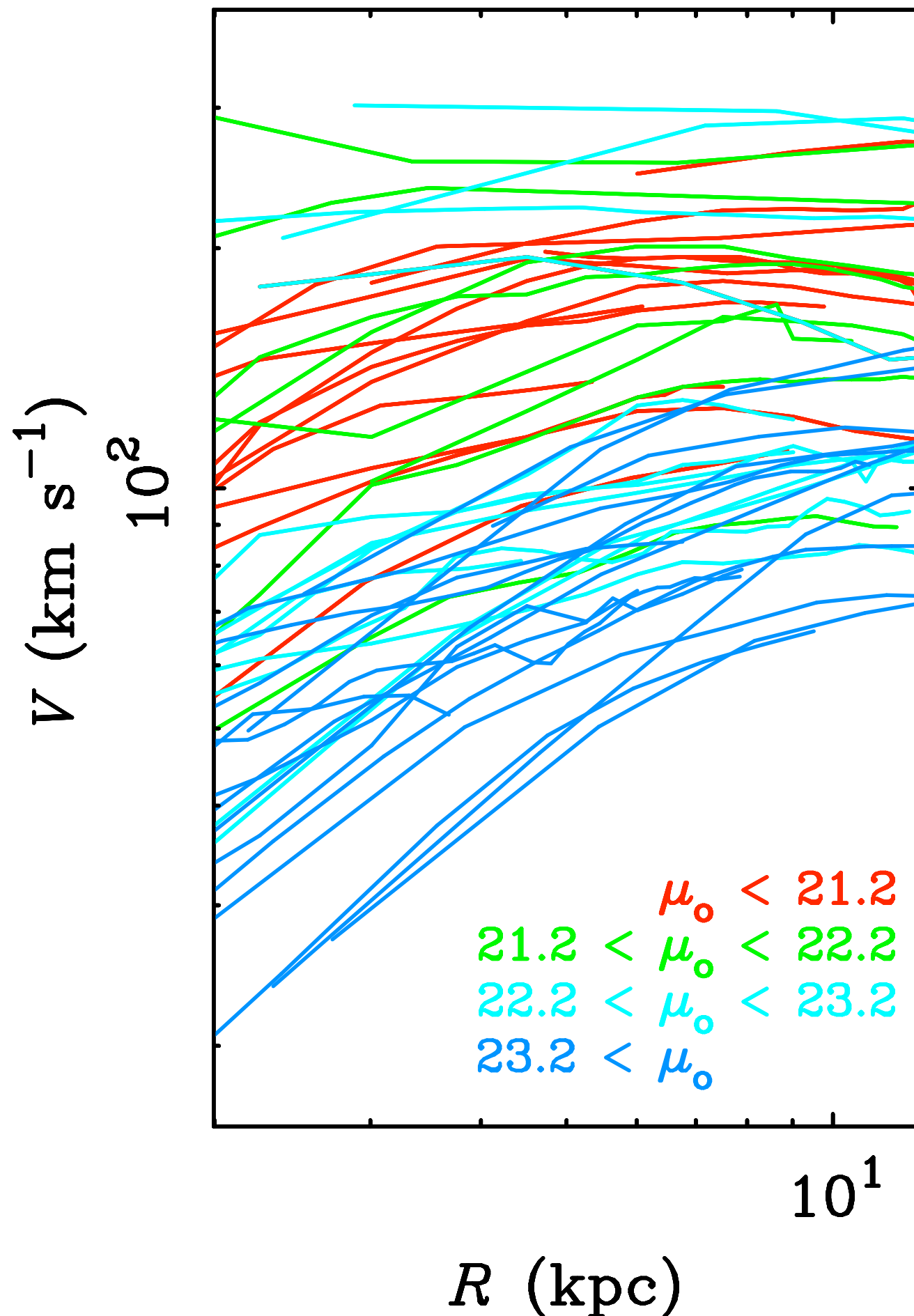
- MOND accurately predicts the BTG location of gas dominated galaxies with zero free parameters.
- CDM does not do this.





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- Rotation Curve Shapes

- Surface Density \sim Surface Brightness

- Detailed Rotation Curve Fits

- Stellar Population Mass-to-Light Ratios



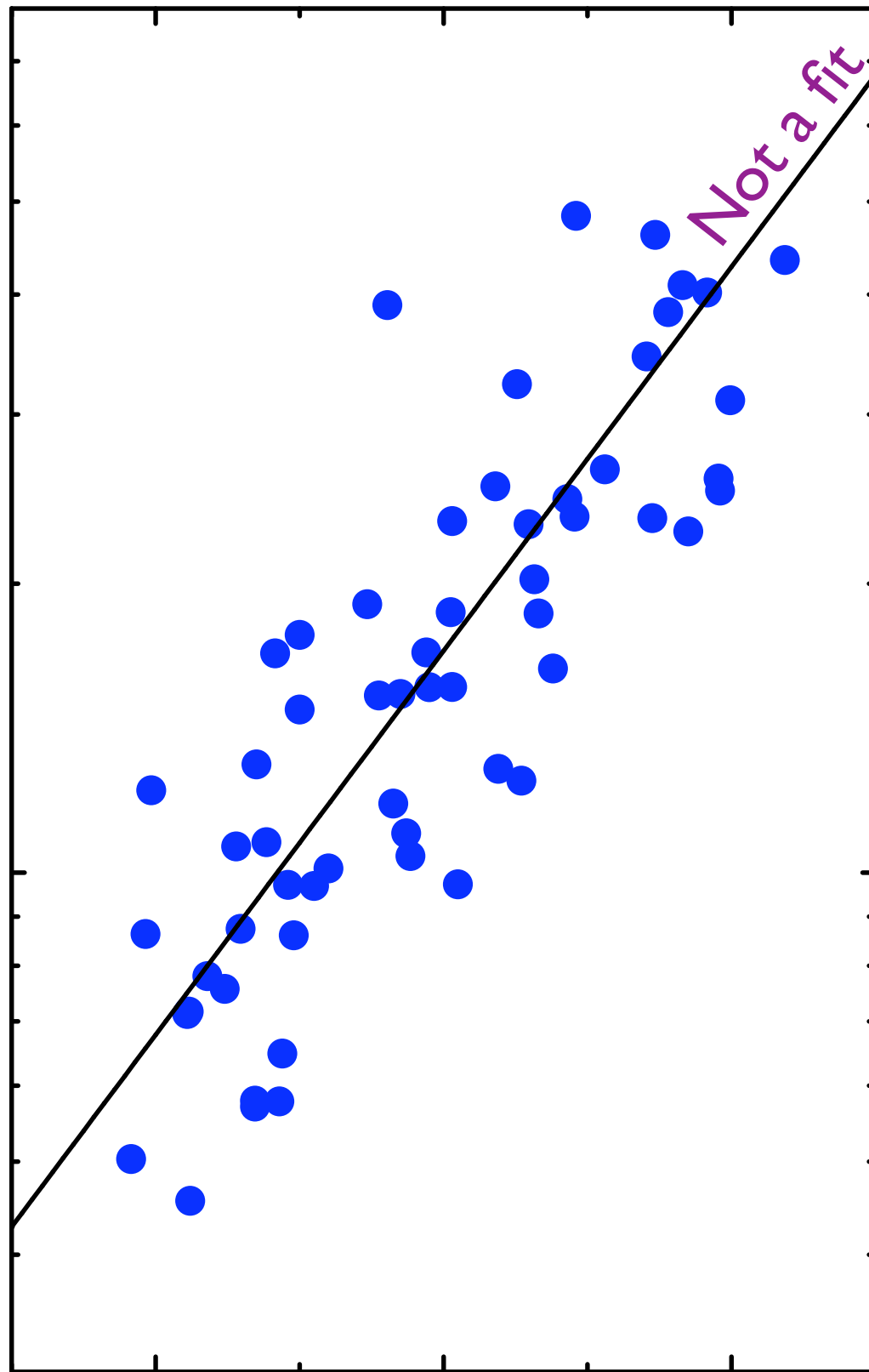
mass surface density

$$\xi = V^2/(Gh)$$

5

1

0.5



24

22

20

μ_o

surface brightness



MOND predictions

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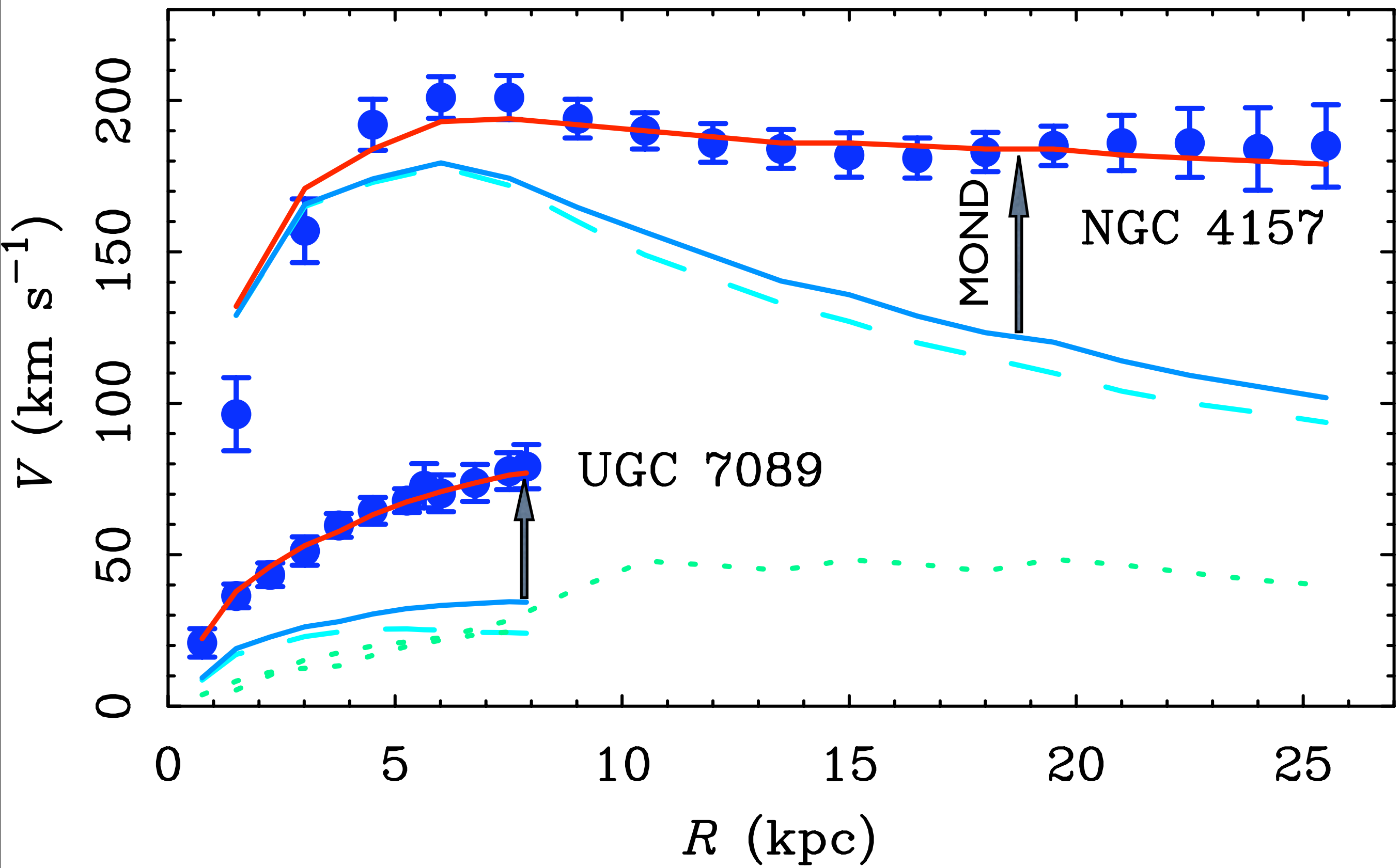
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• Surface Density \sim Surface Brightness

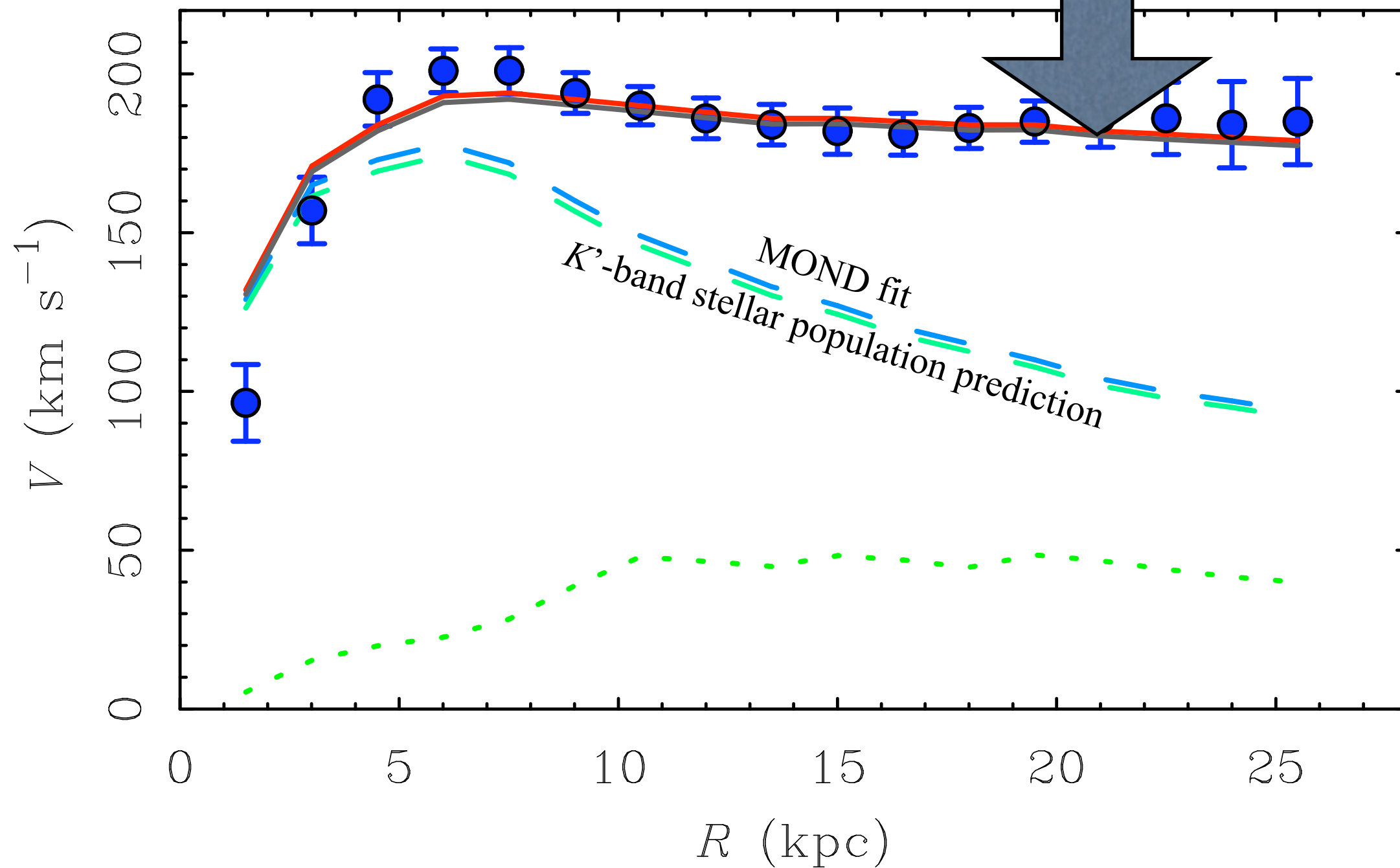
- Detailed Rotation Curve Fits

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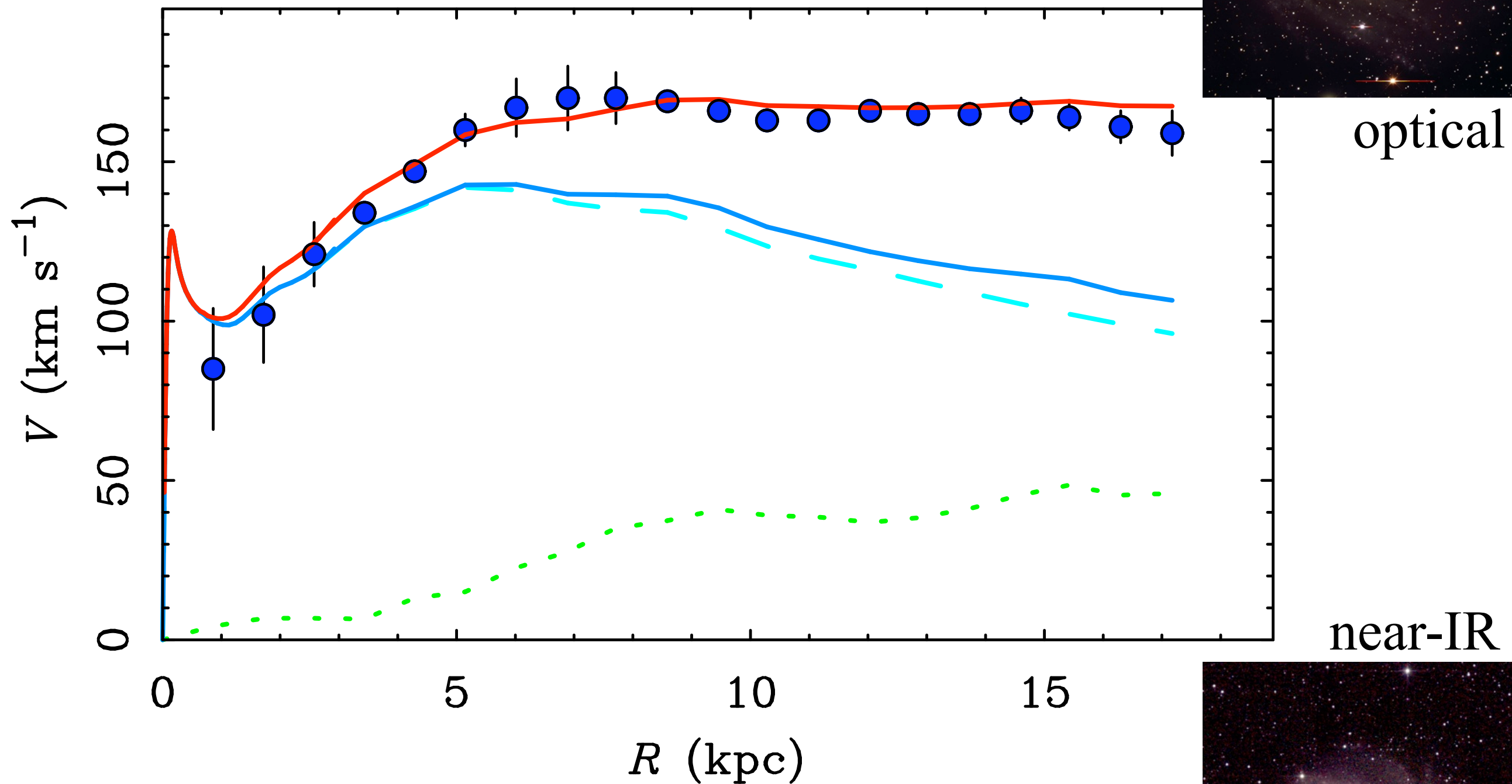


predictive power: zero free parameters

NGC 4157



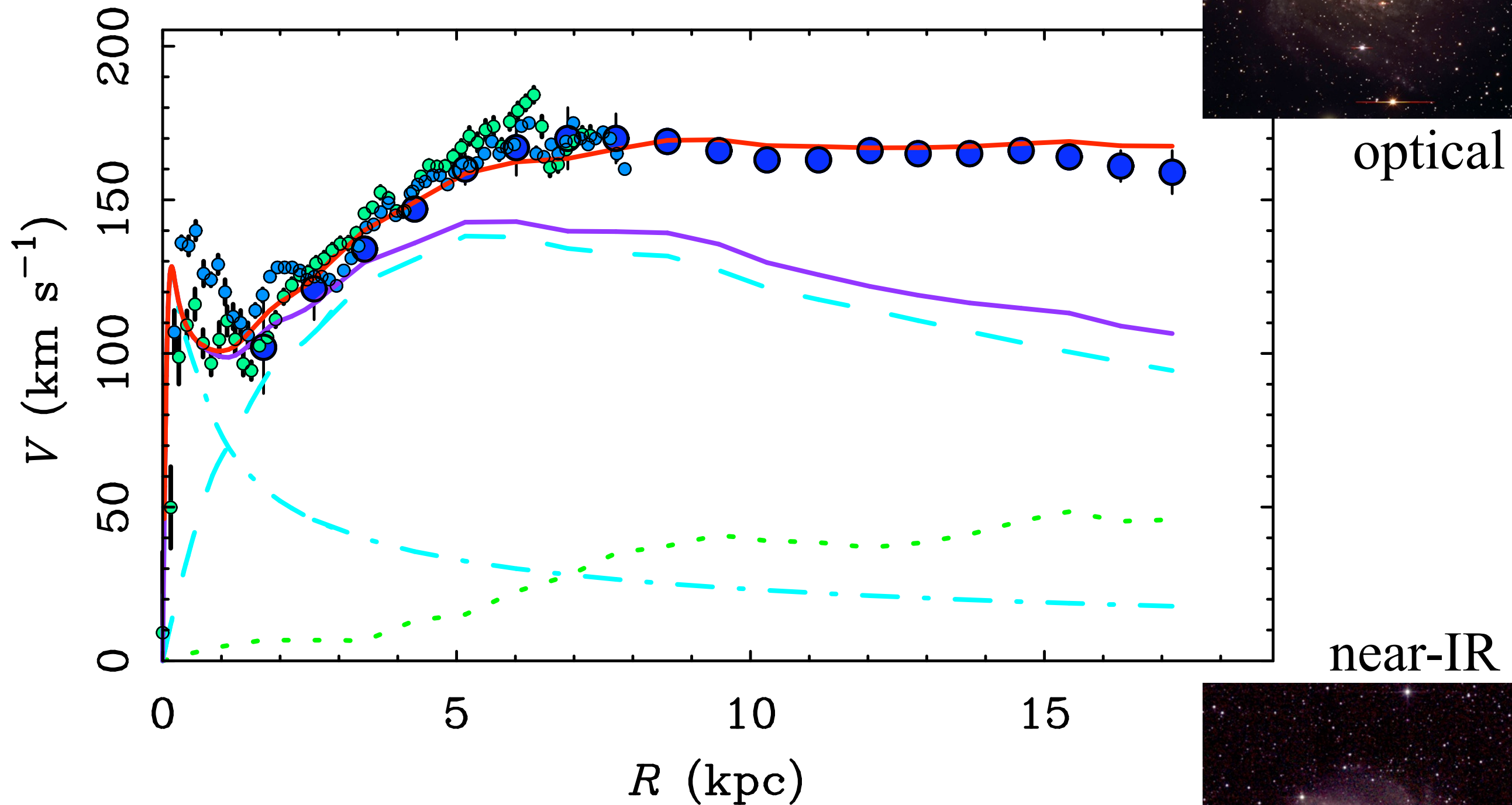
NGC 6946 - small bulge predicted



Renzo's Rule:

"When you see a feature in the light, you see a corresponding feature in the rotation curve."

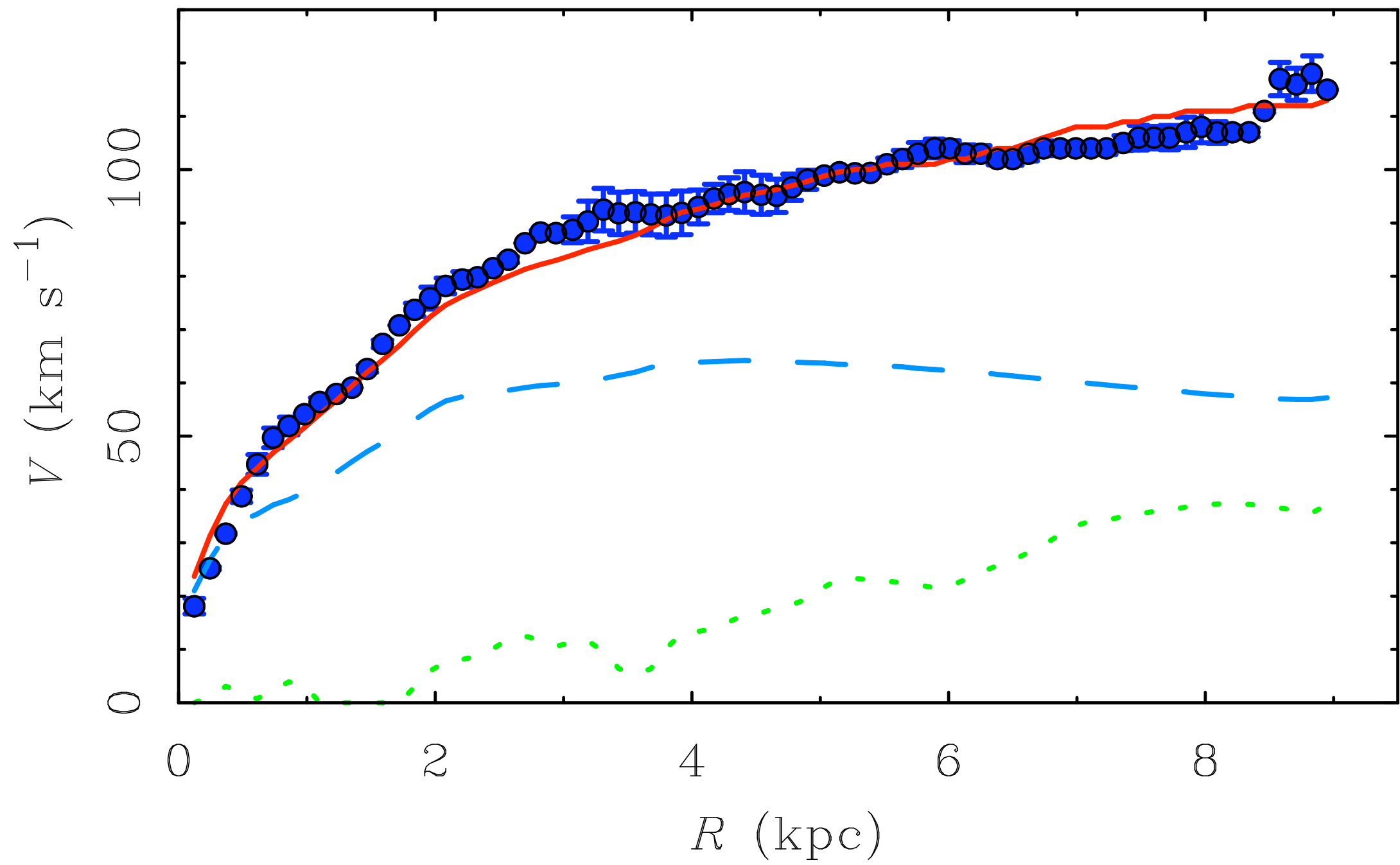
NGC 6946 - small bulge observed



Renzo's Rule:

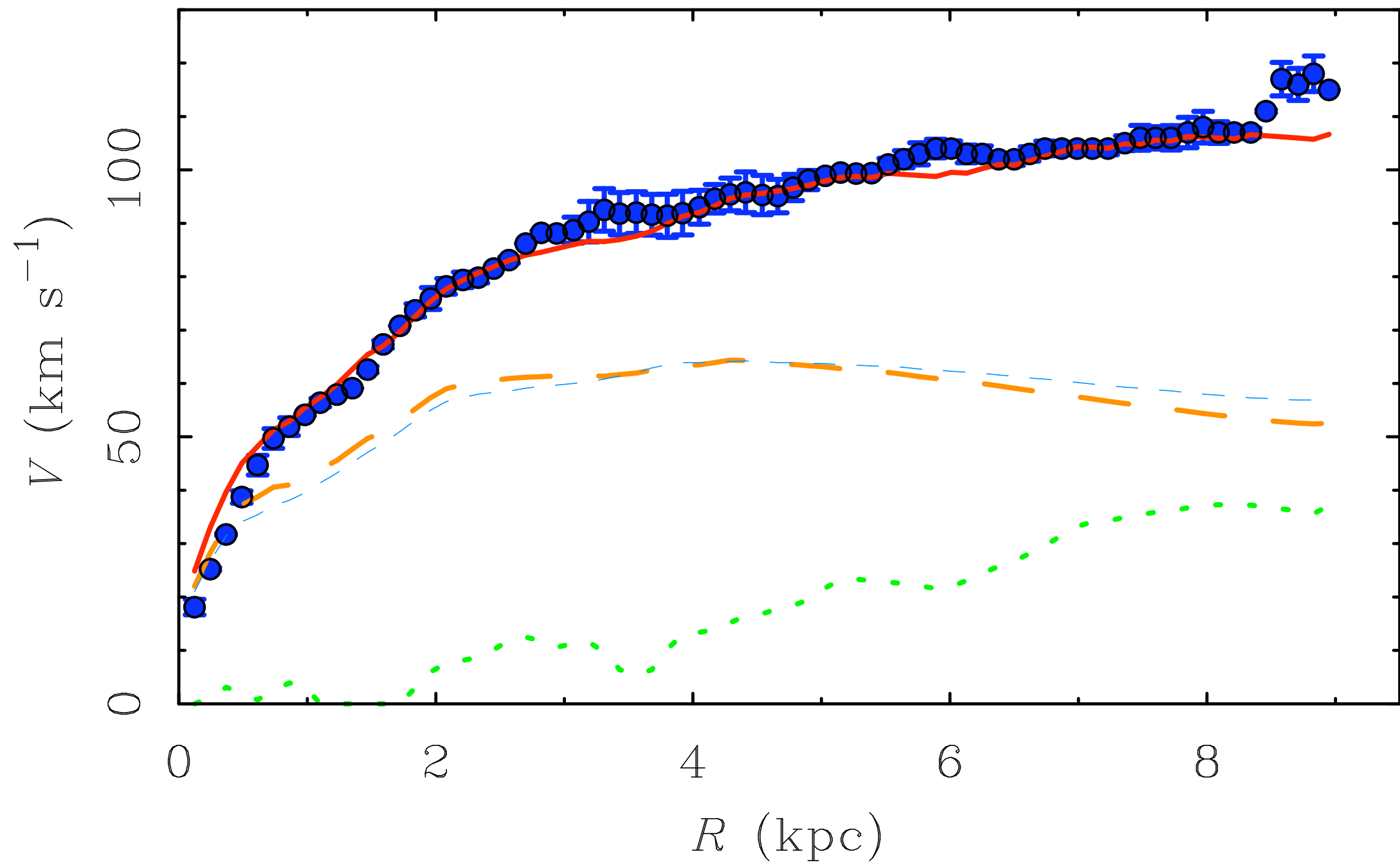
"When you see a feature in the light, you see a corresponding feature in the rotation curve."

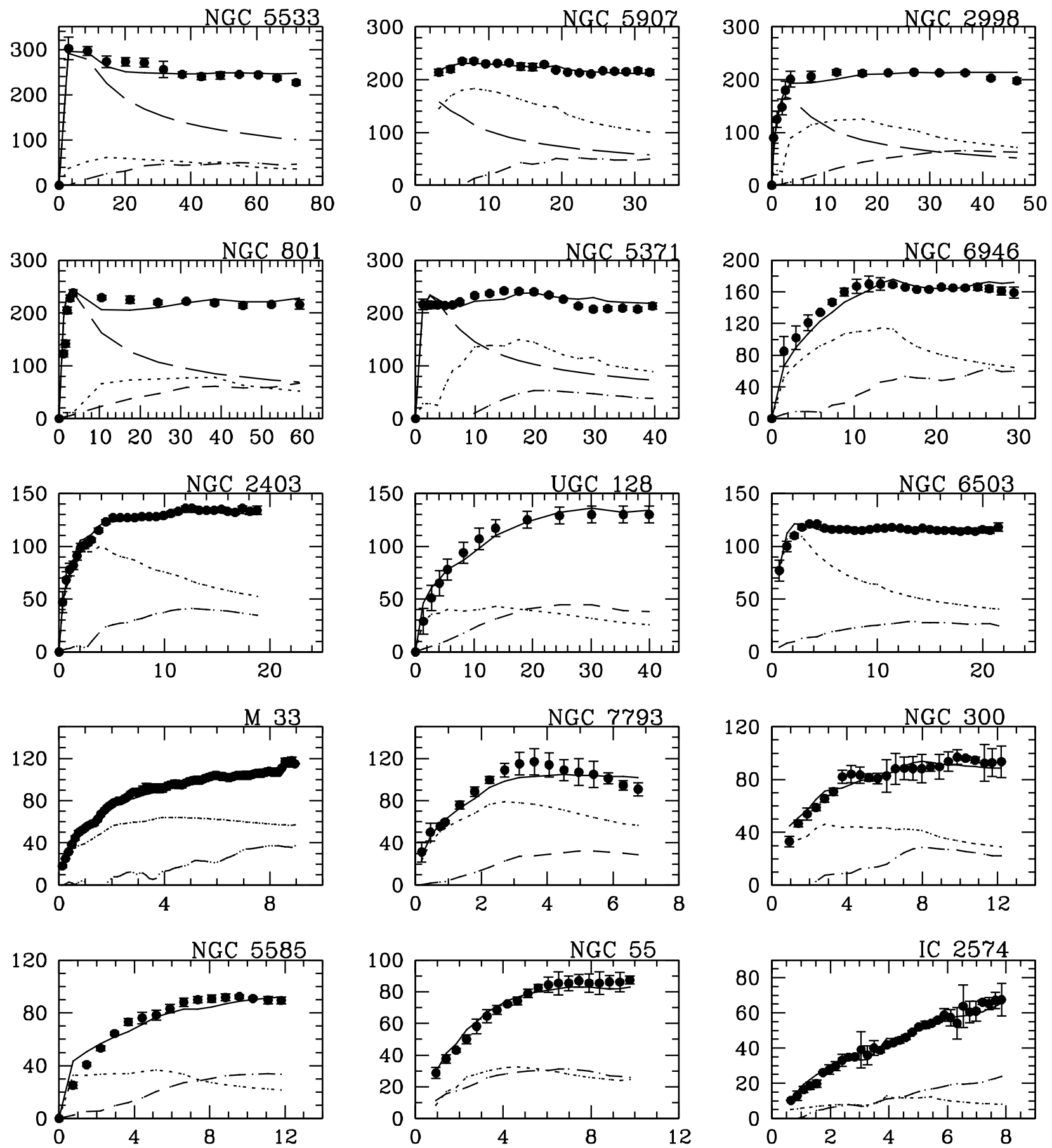
M33

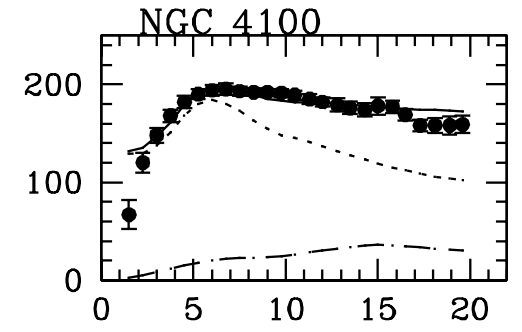
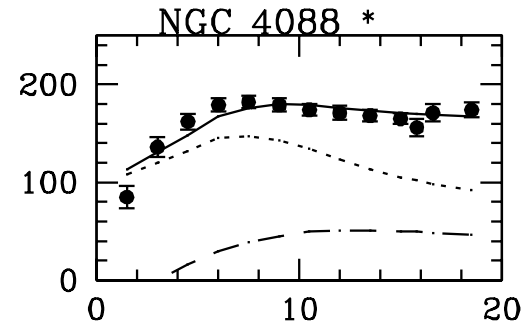
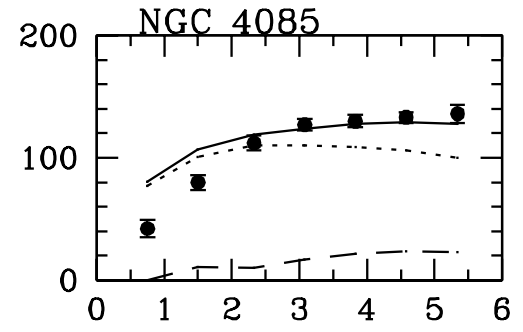
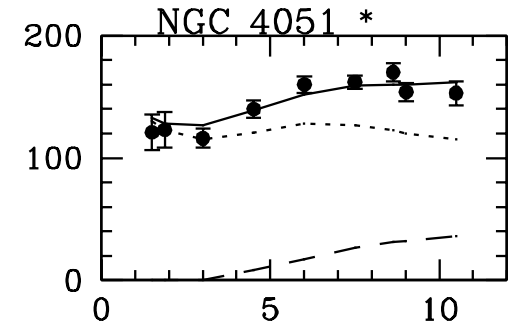
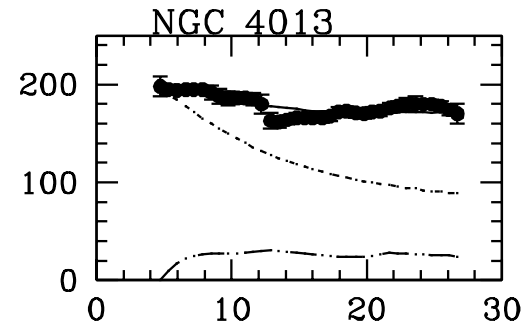
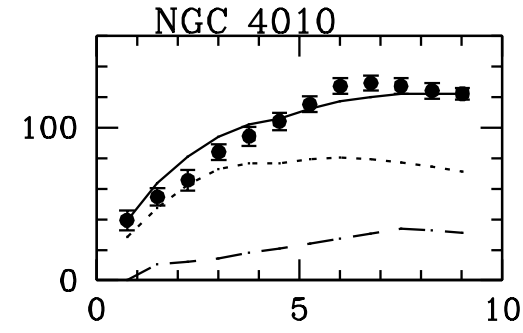
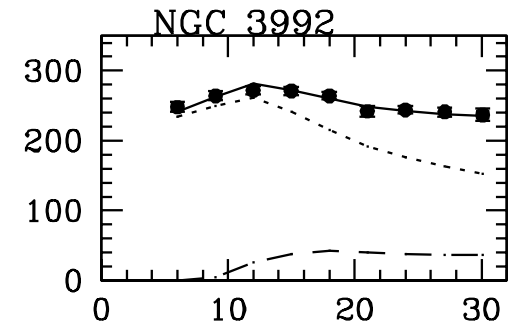
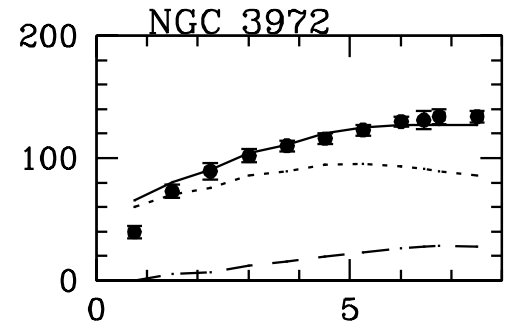
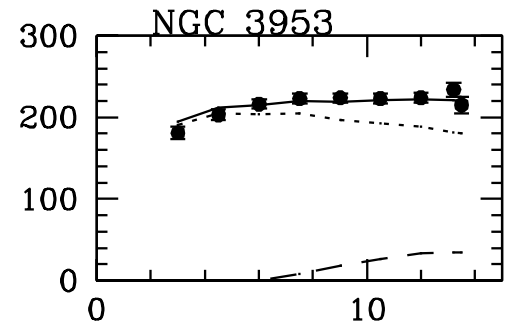
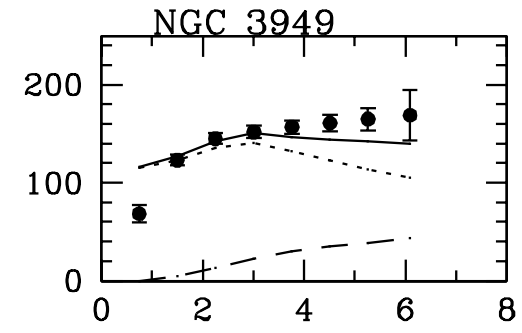
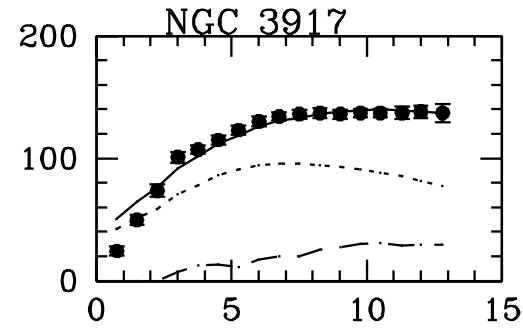
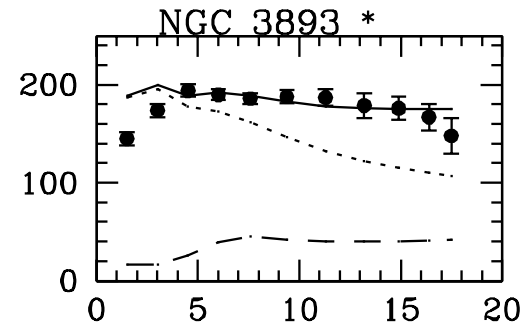
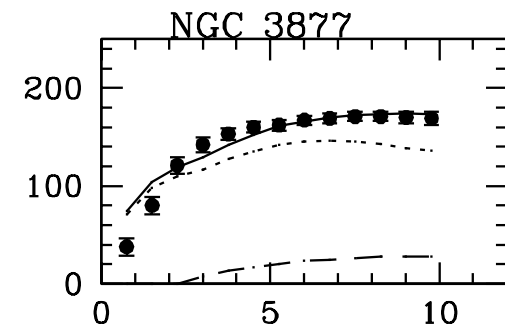
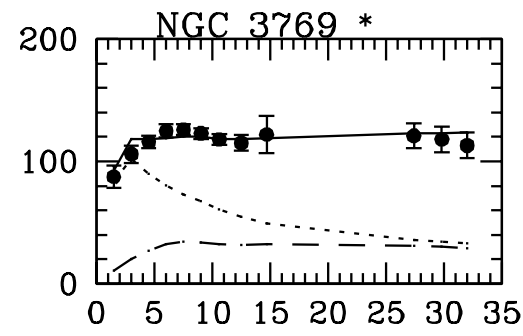
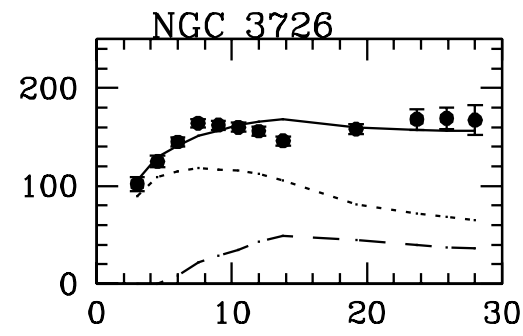


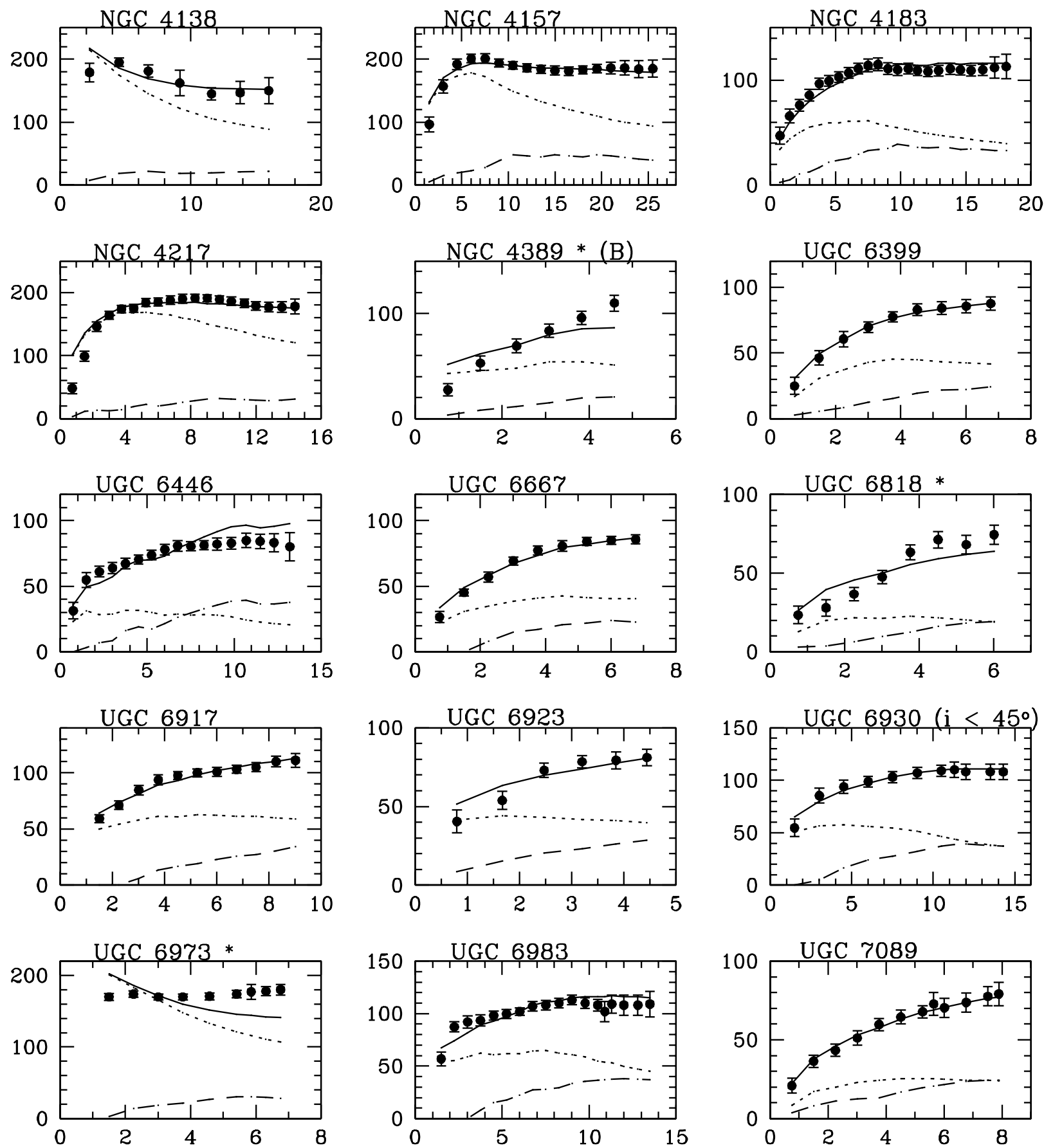
M33

color gradient corrected

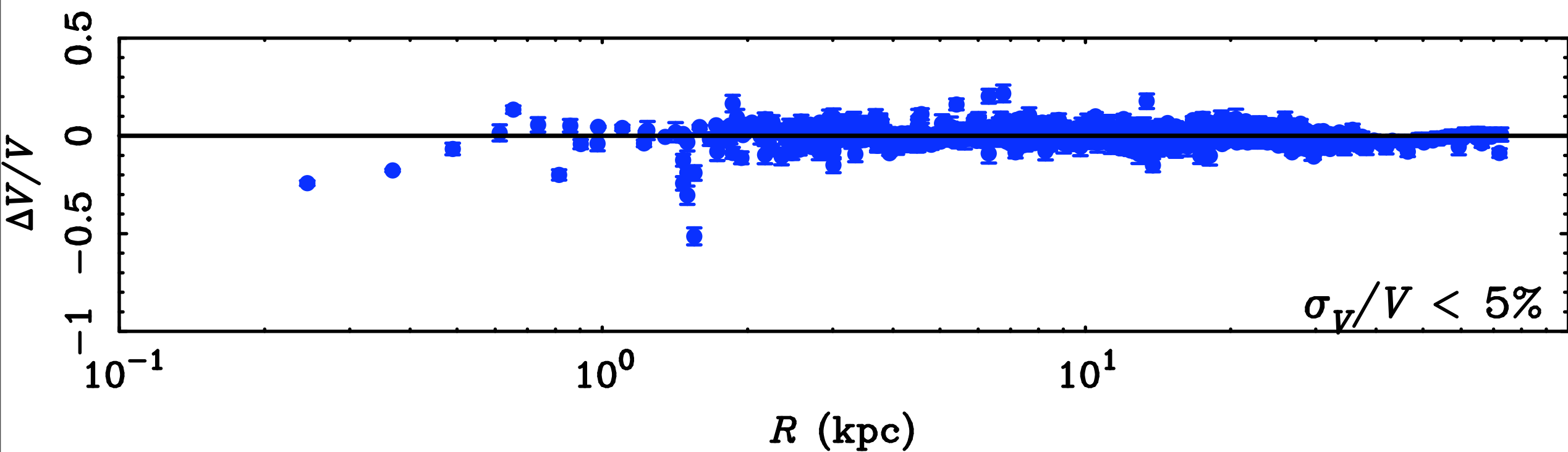
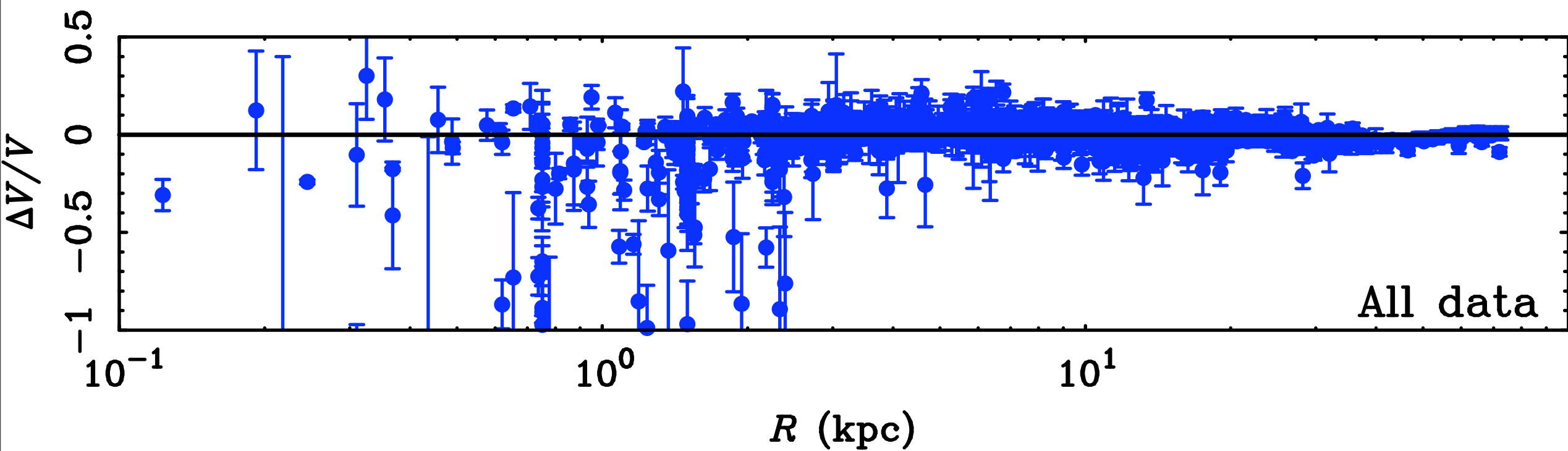




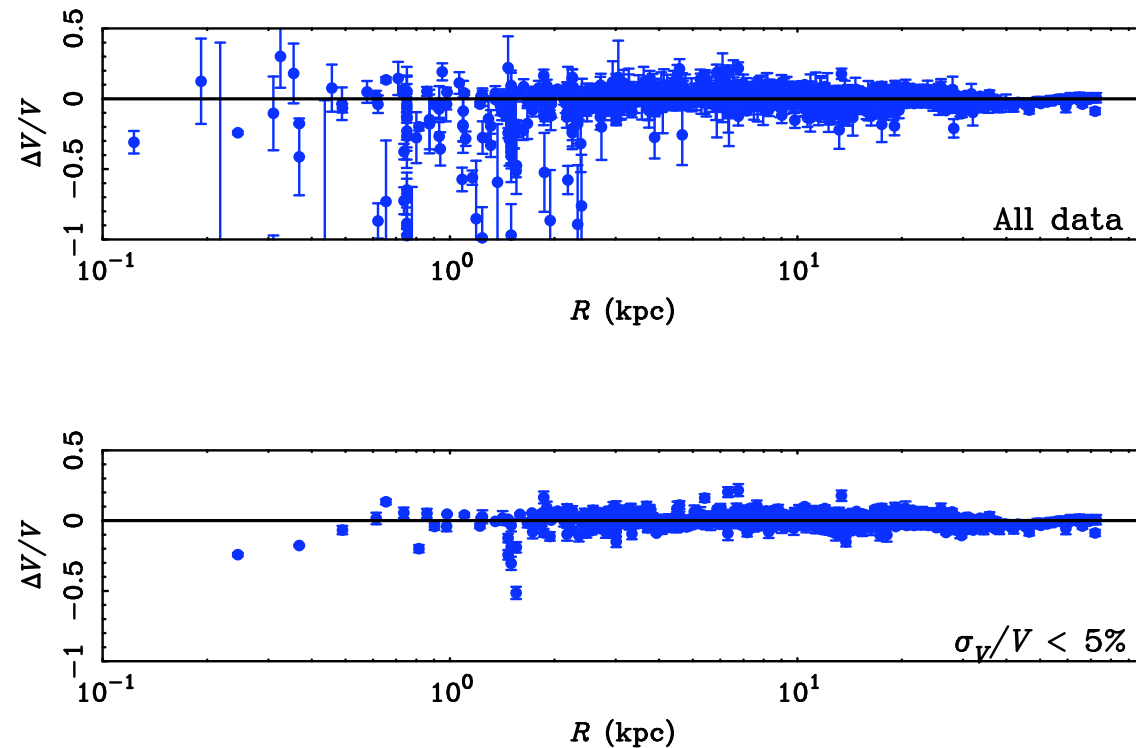




Residuals of MOND fits



MOND predictions



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- ✓ • Slope = 4

- ✓ • Normalization = $1/(a_0 G)$

- ✓ • Fundamentally a relation between Disk Mass and V_{flat}

- ✓ • No Dependence on Surface Brightness

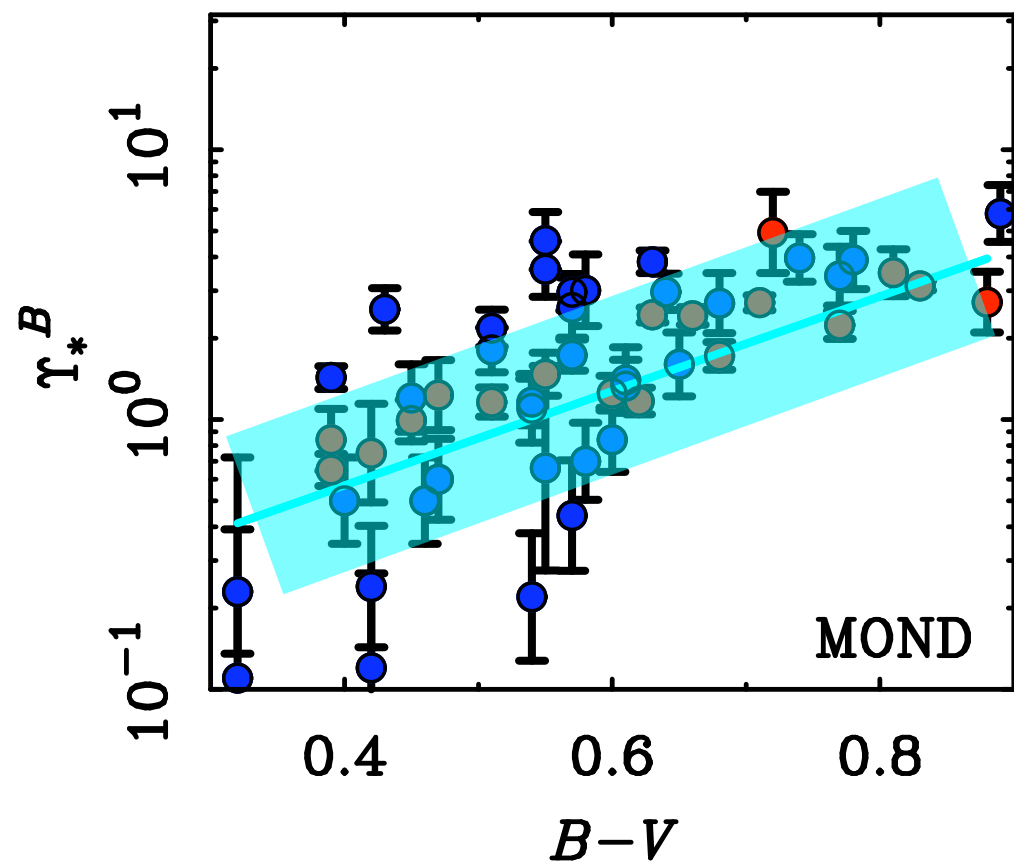
- ✓ • Dependence of conventional M/L on radius and surface brightness

- ✓ • Rotation Curve Shapes

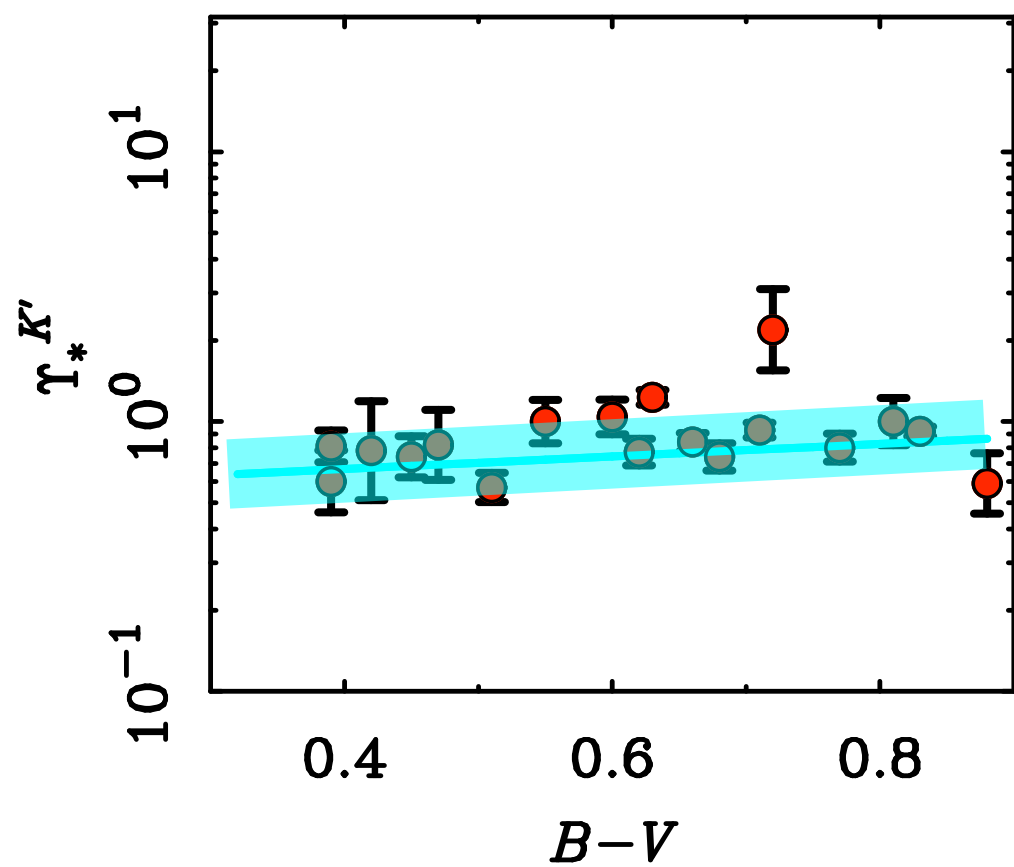
- ✓ • Surface Density \sim Surface Brightness

- ✓ • Detailed Rotation Curve Fits

- Stellar Population Mass-to-Light Ratios



Line: stellar population model
(mean expectation)



MOND predictions

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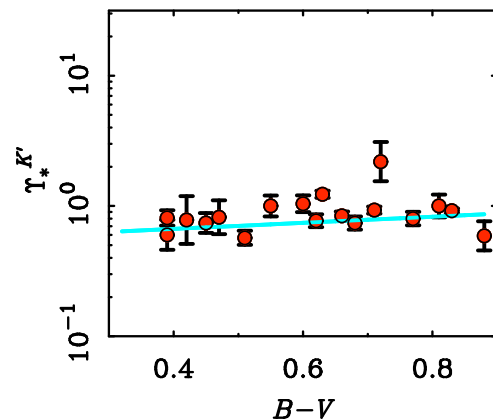
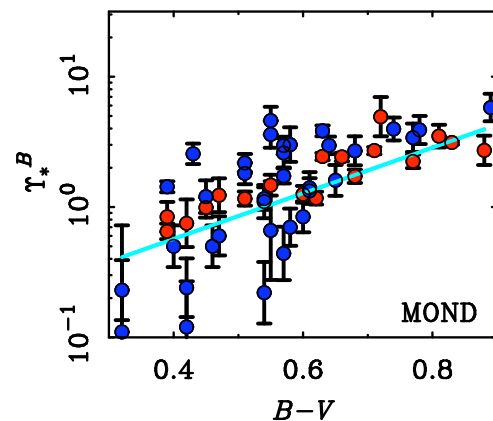
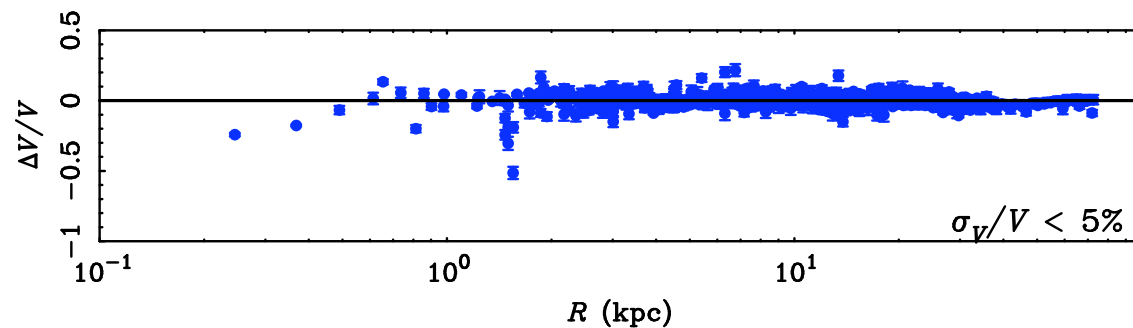
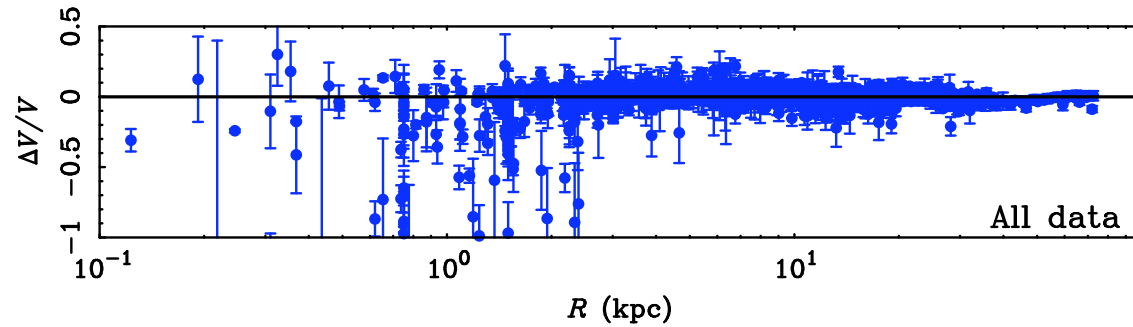
- Surface Density \sim Surface Brightness



- Detailed Rotation Curve Fits

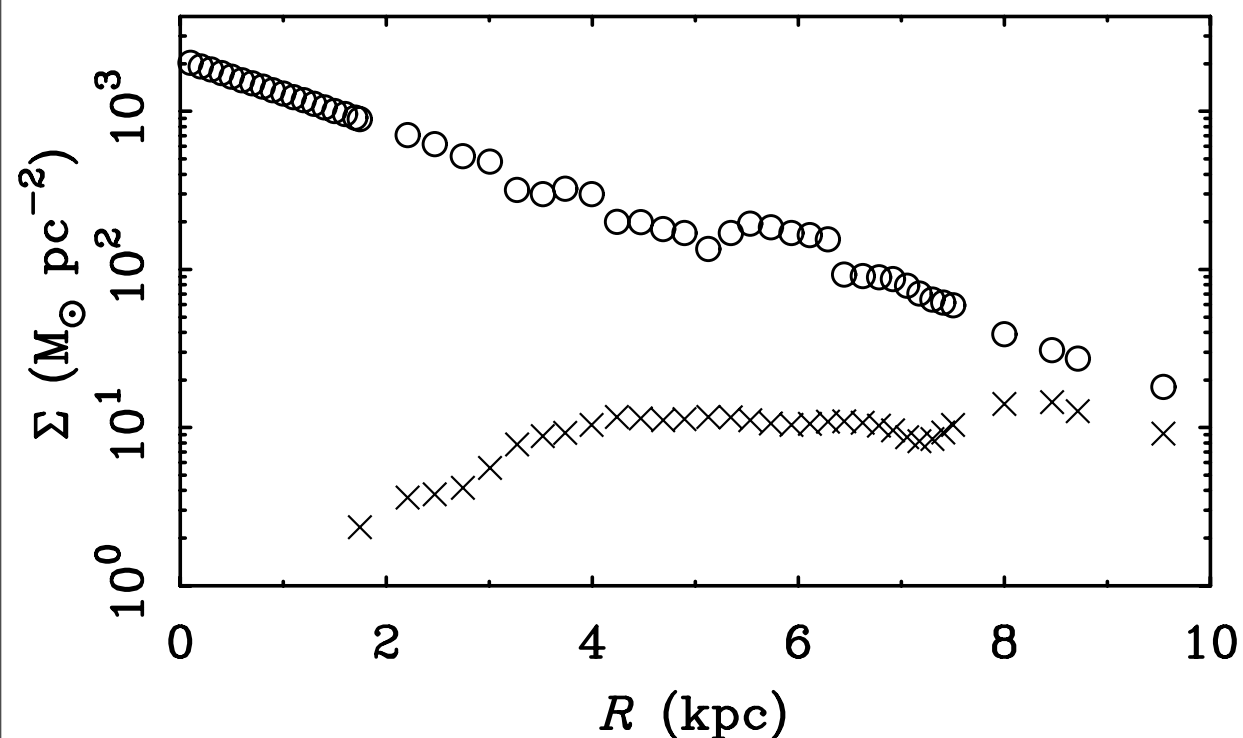


- Stellar Population Mass-to-Light Ratios

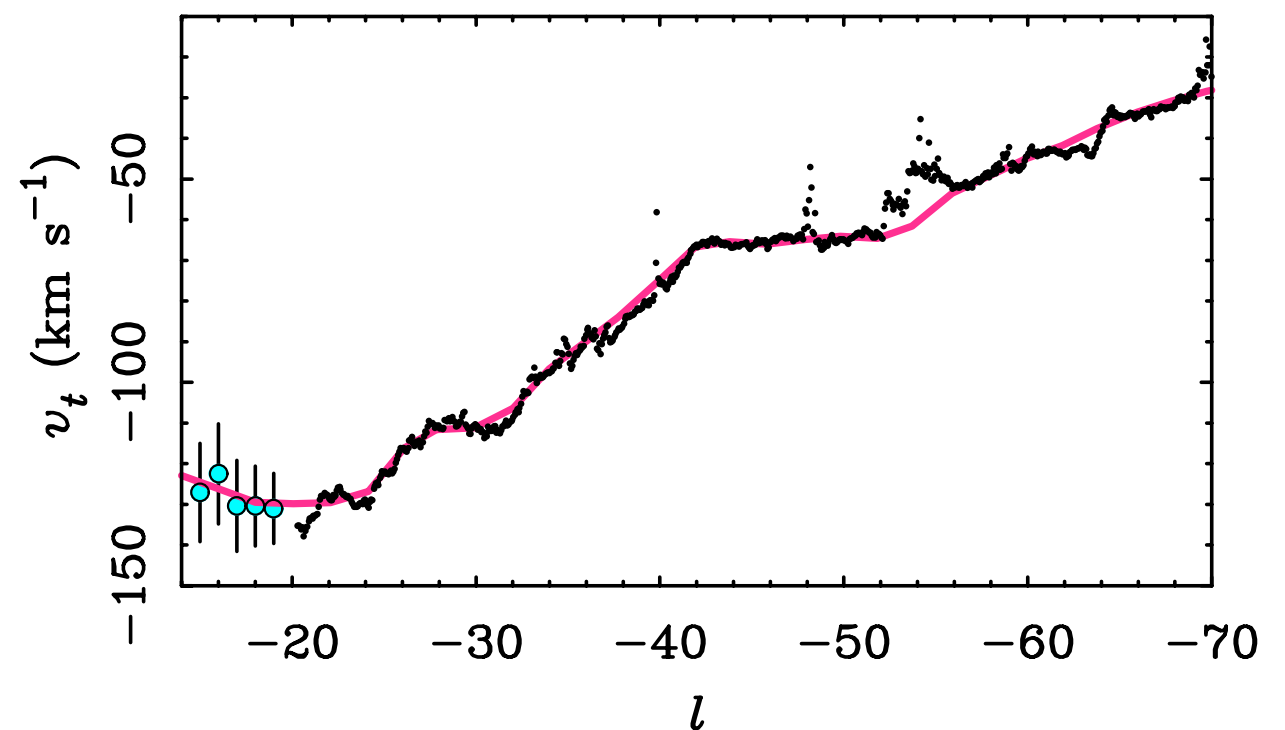


Can we reverse the procedure?

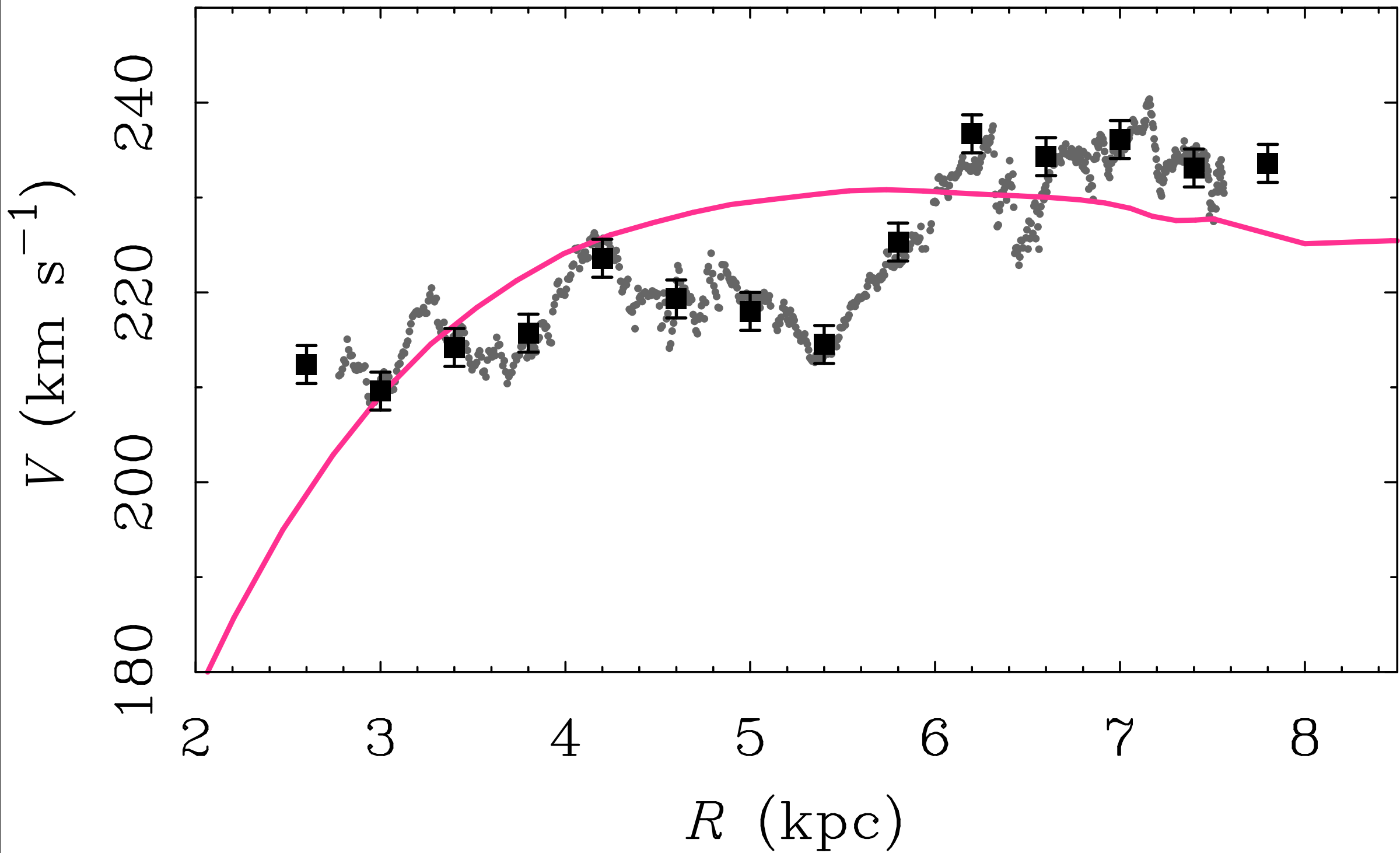
Rather than fitting rotation curves
given the photometry, can we infer
the baryonic mass distribution from
the rotation curve?



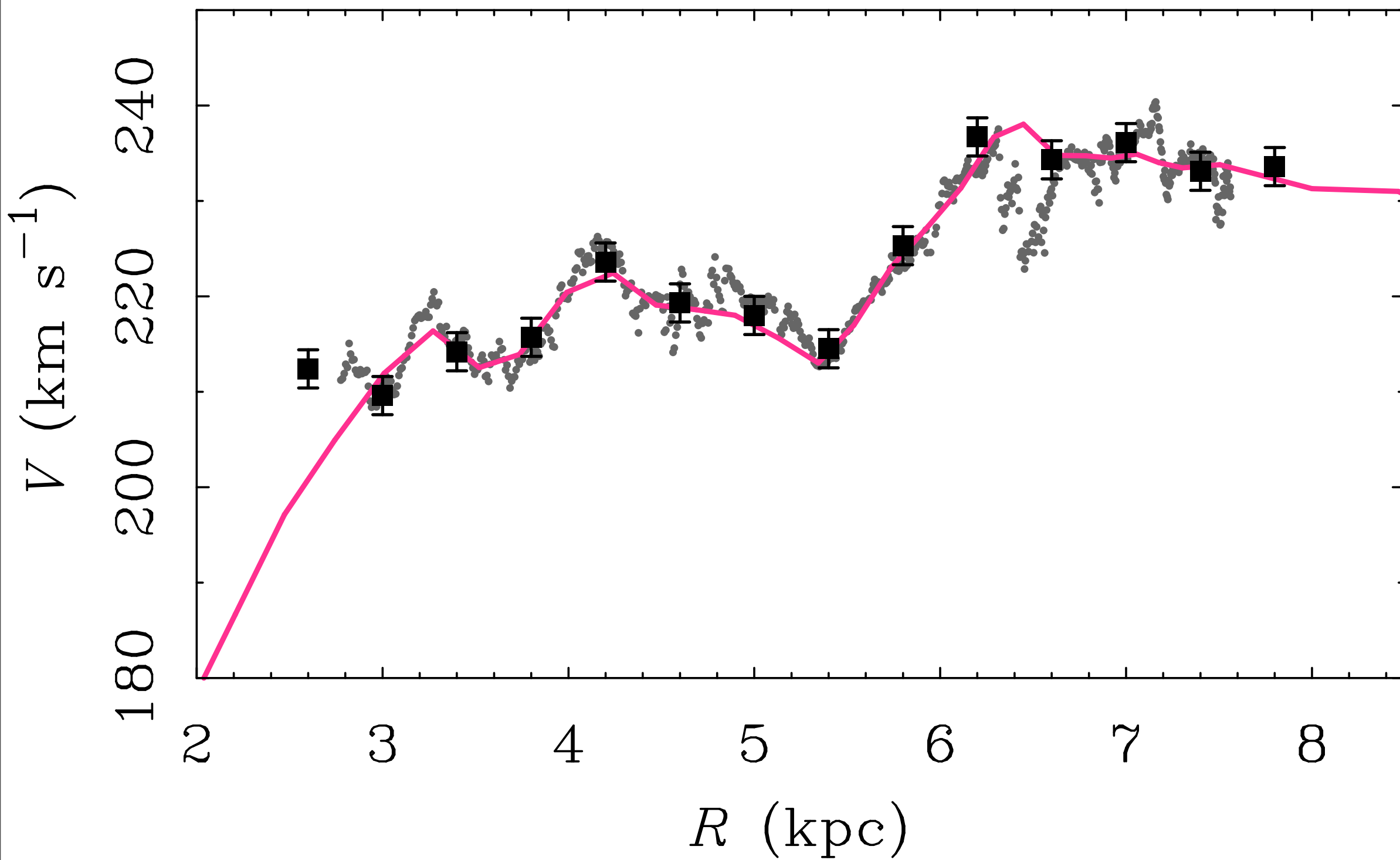
Milky Way terminal velocities



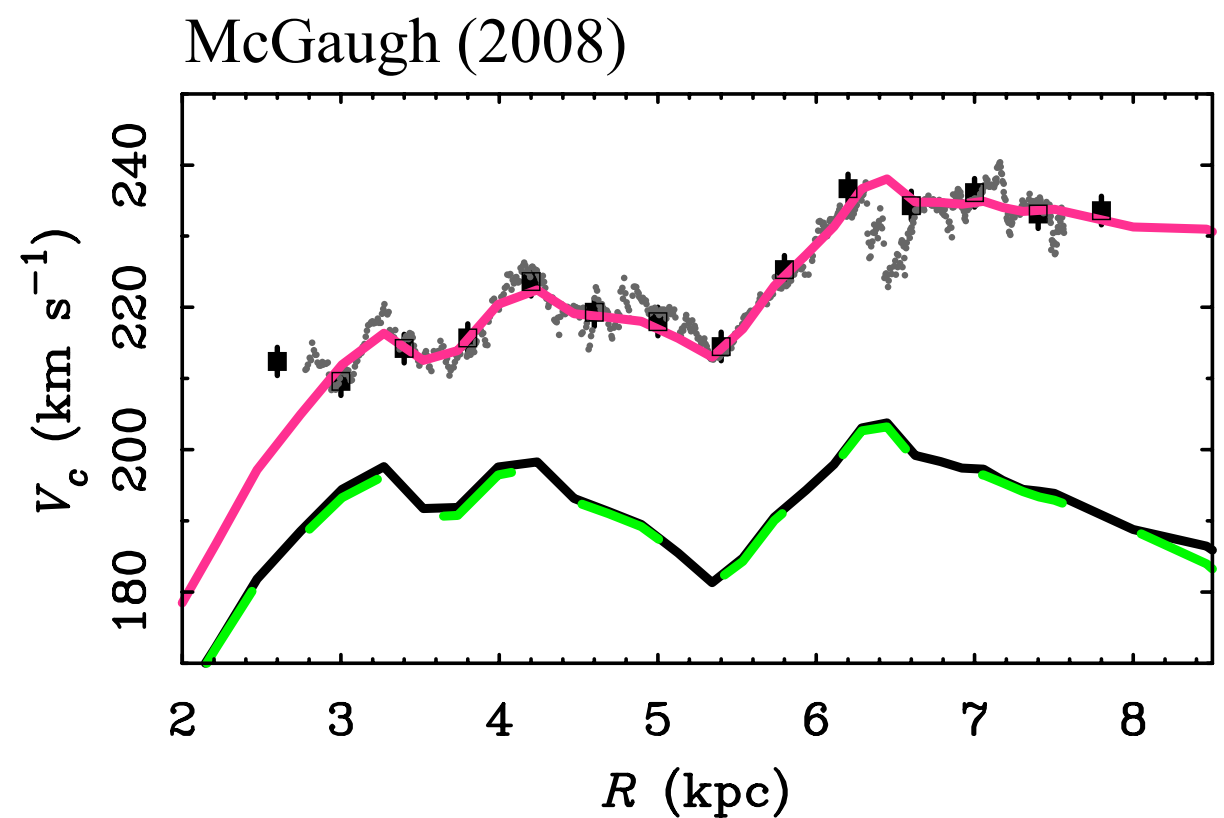
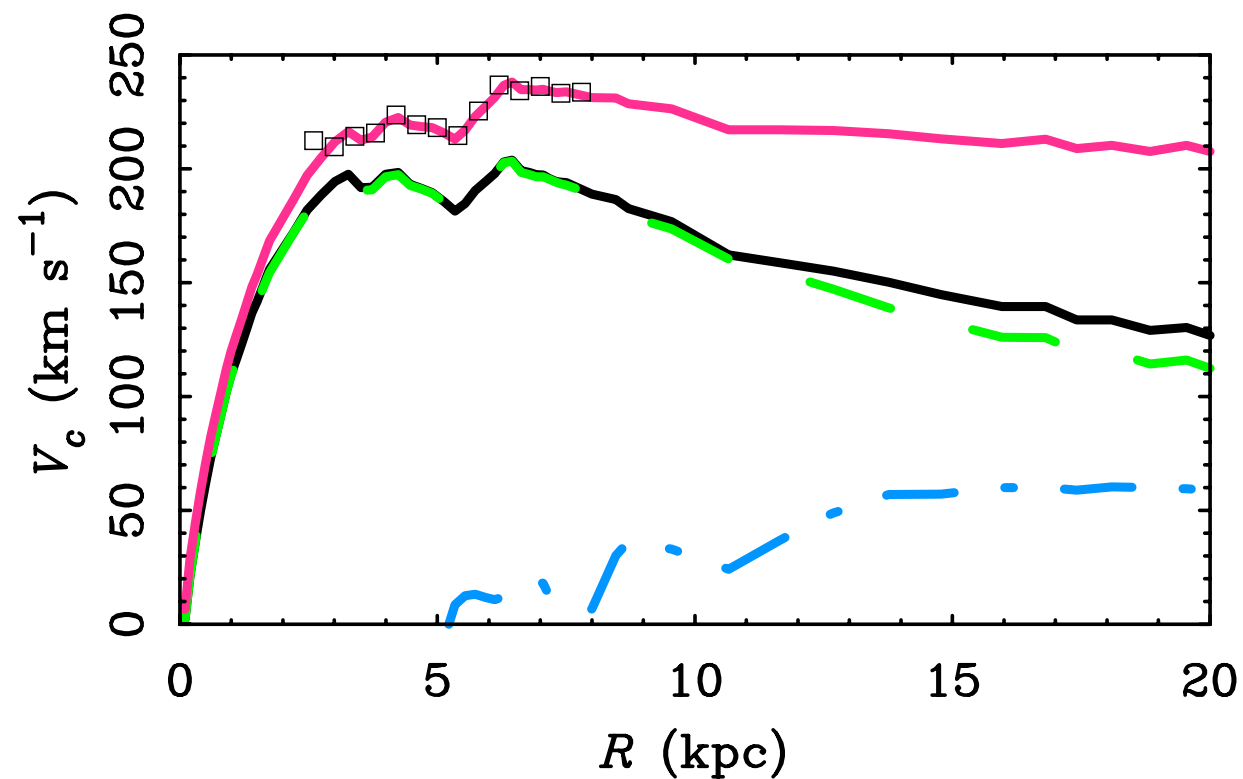
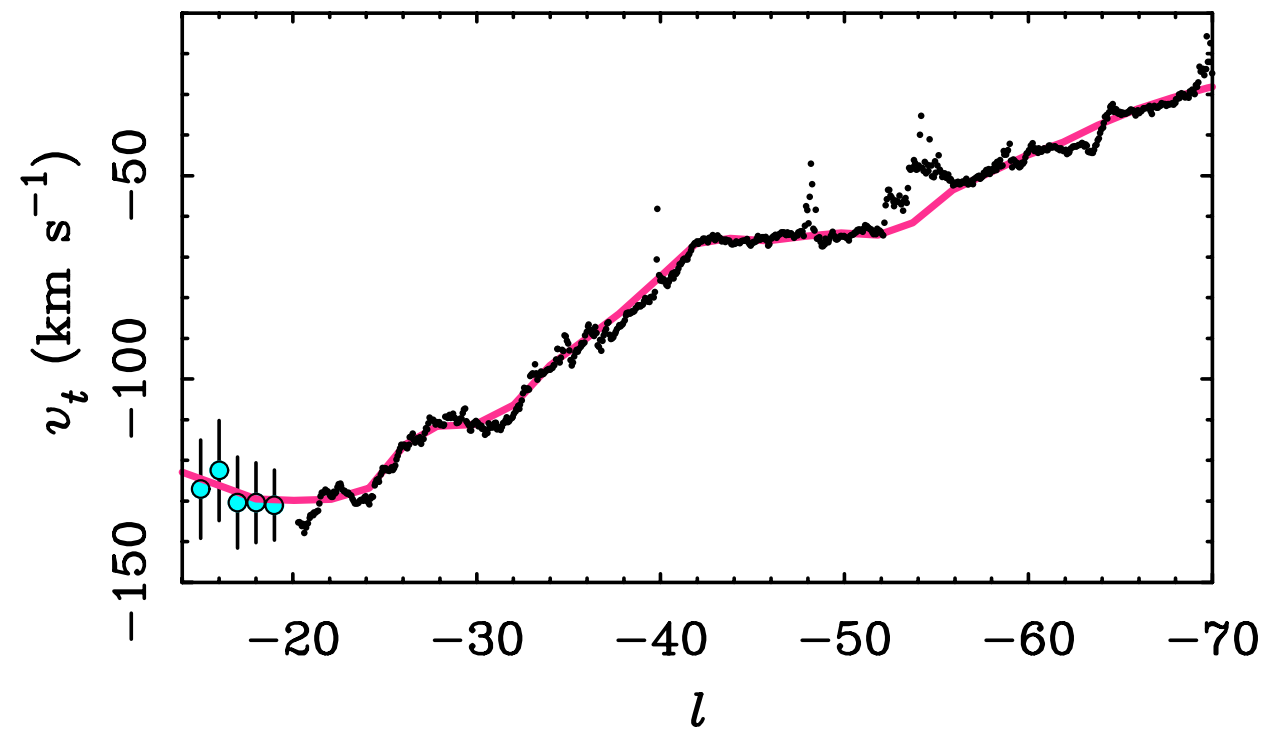
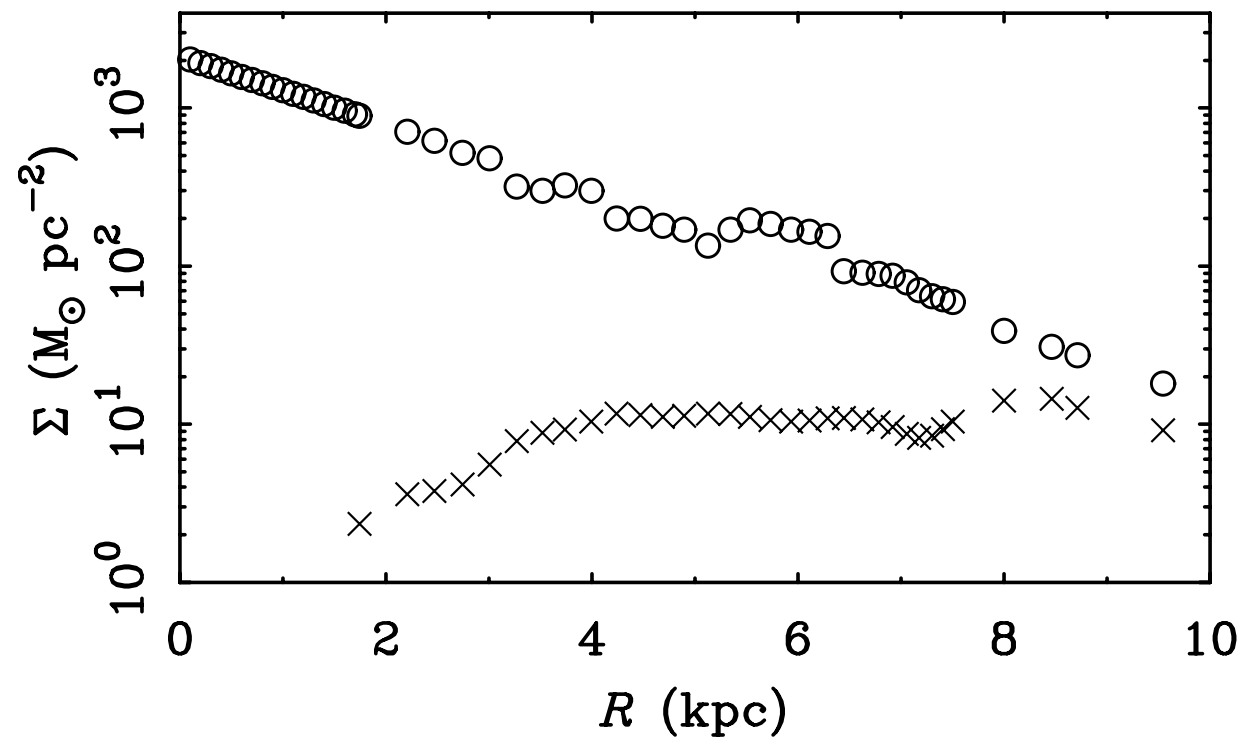
initial guess



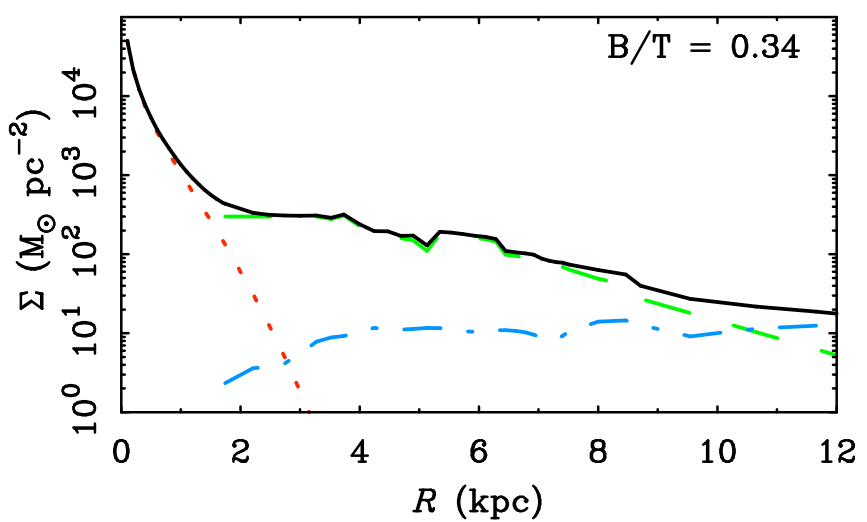
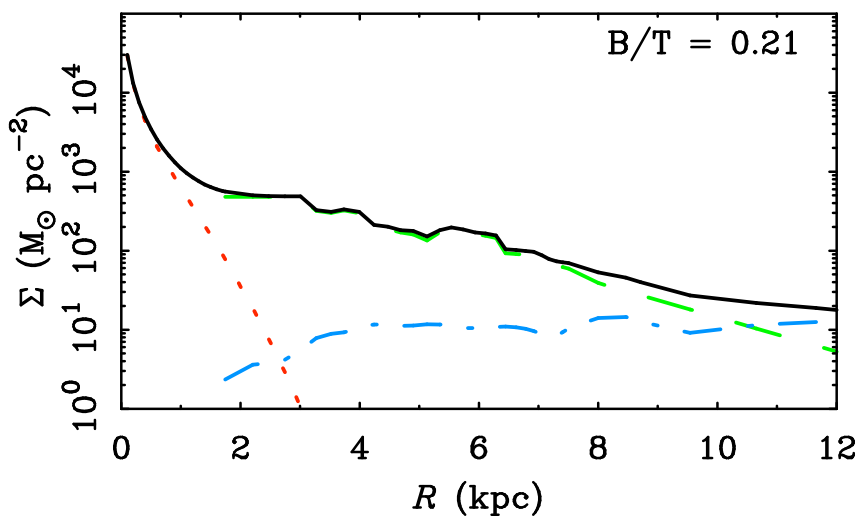
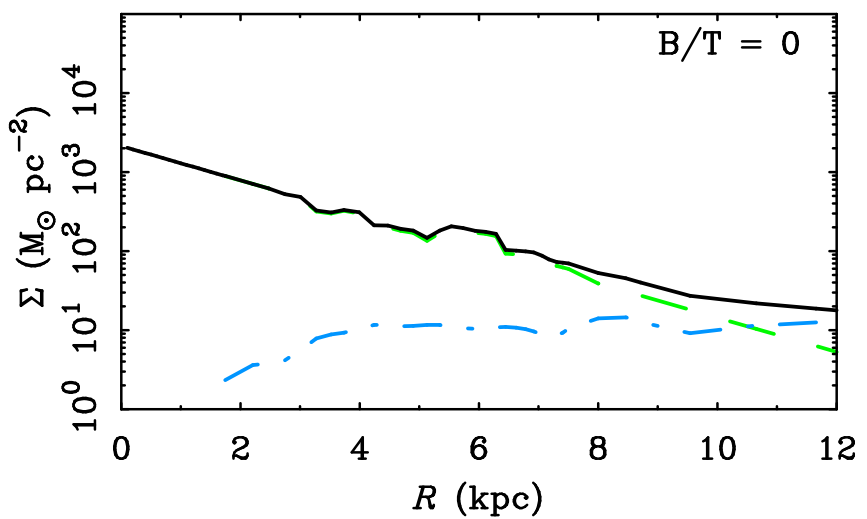
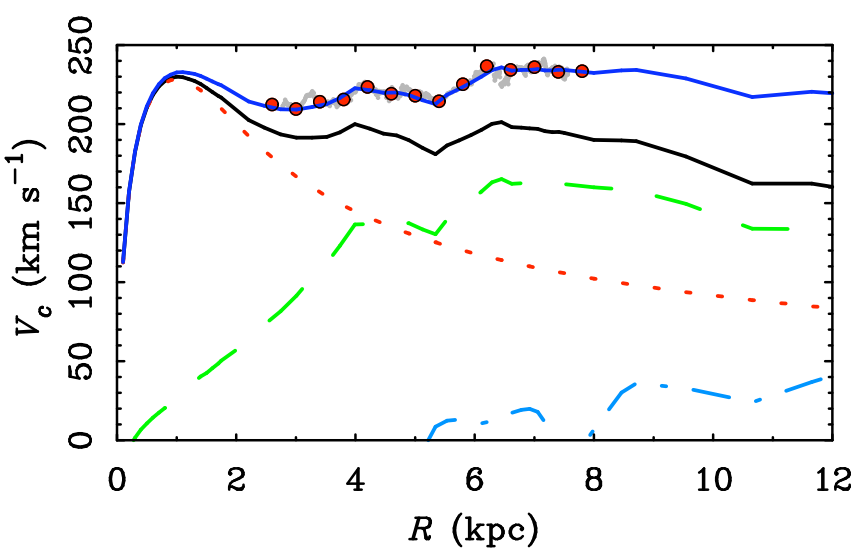
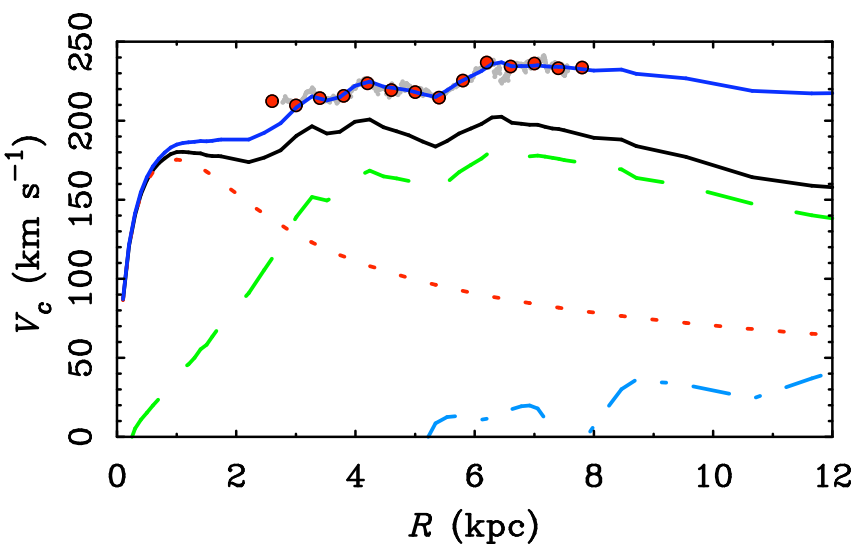
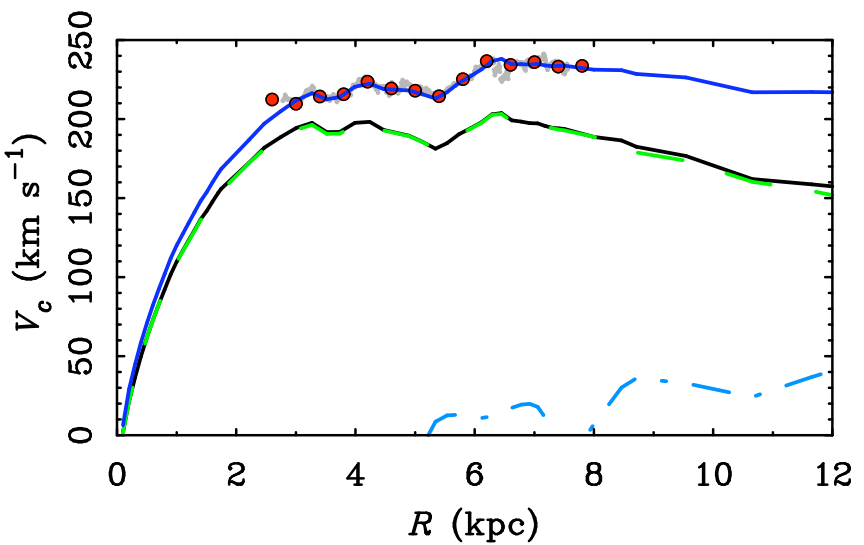
final fit



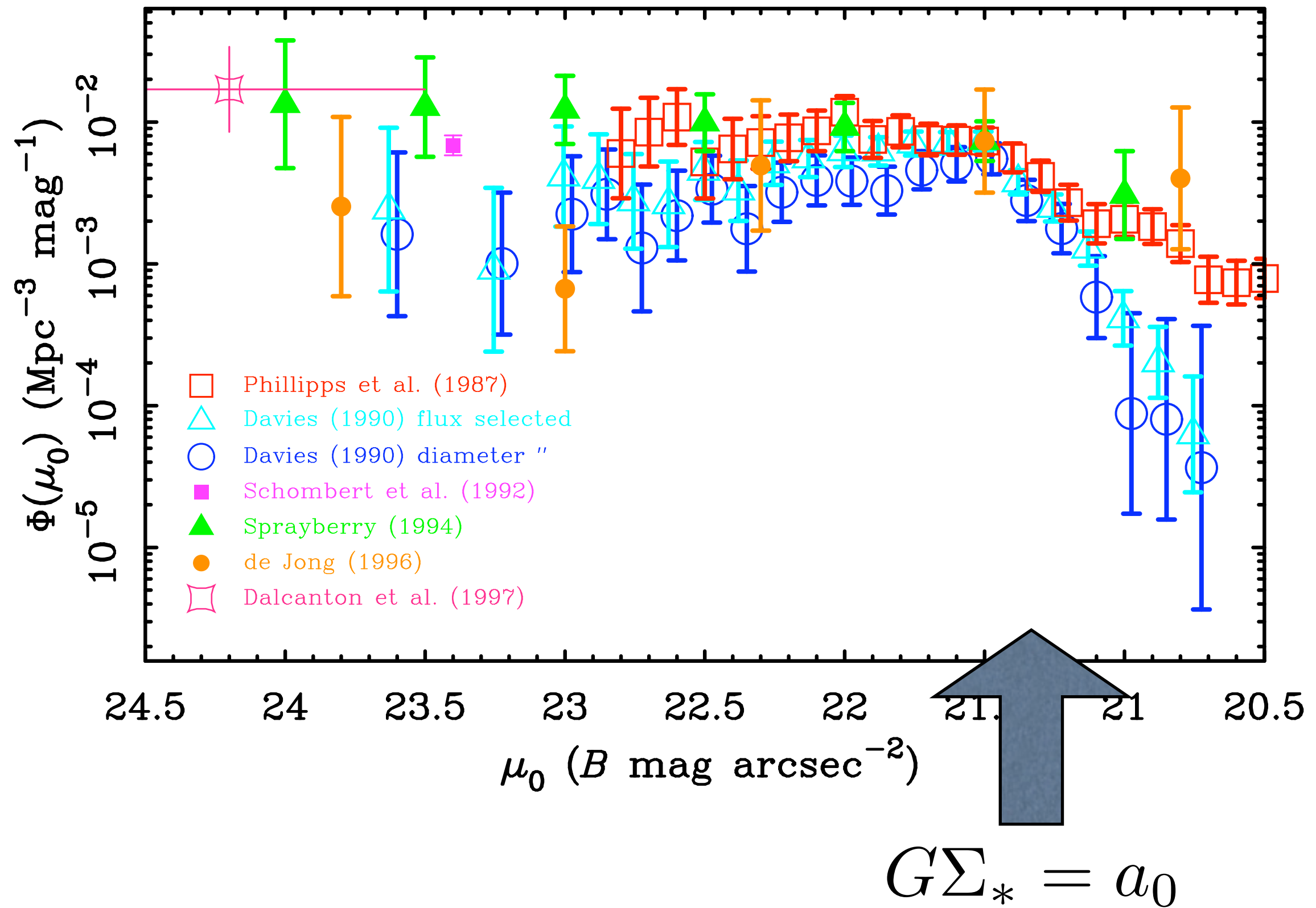
Obtain plausible mass profile; predictions testable with GAIA



Allowing for a significant bulge component
implies that the Milky Way has a Type II disk



Other tests - e.g., disk stability (Milgrom 1989)



disk stability

Brada & Milgrom (1998)

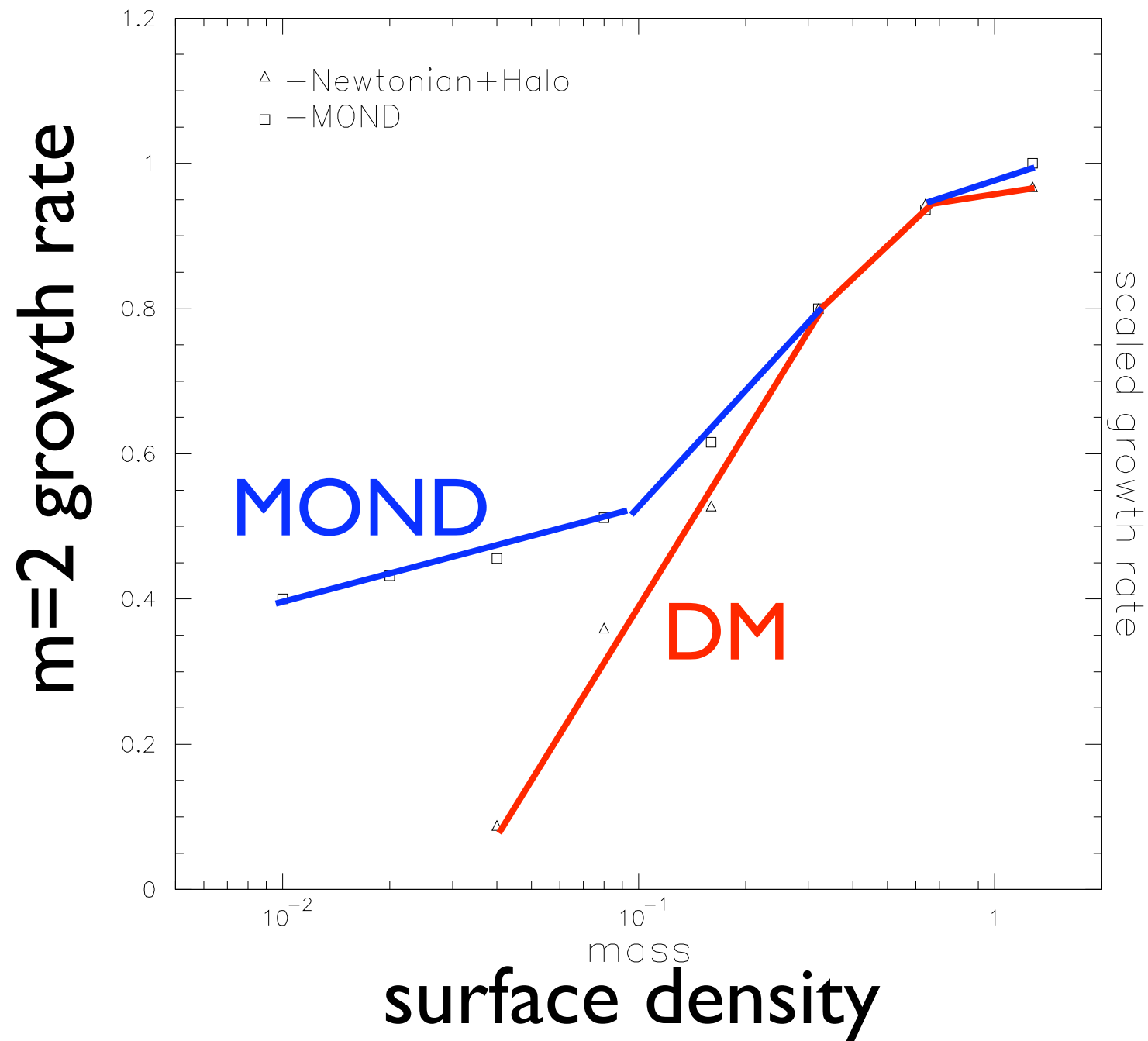
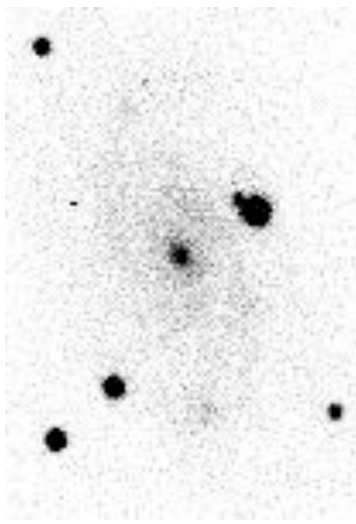


Figure 11: The growth rate, in units of the dynamical time, for the $m=2$ mode as a function of the total mass of the disk. \square MOND, \triangle Newtonian + Halo.

m	Q	time step scaling	Growth rate		halo mass at R=1
			MOND	Newt+DM	

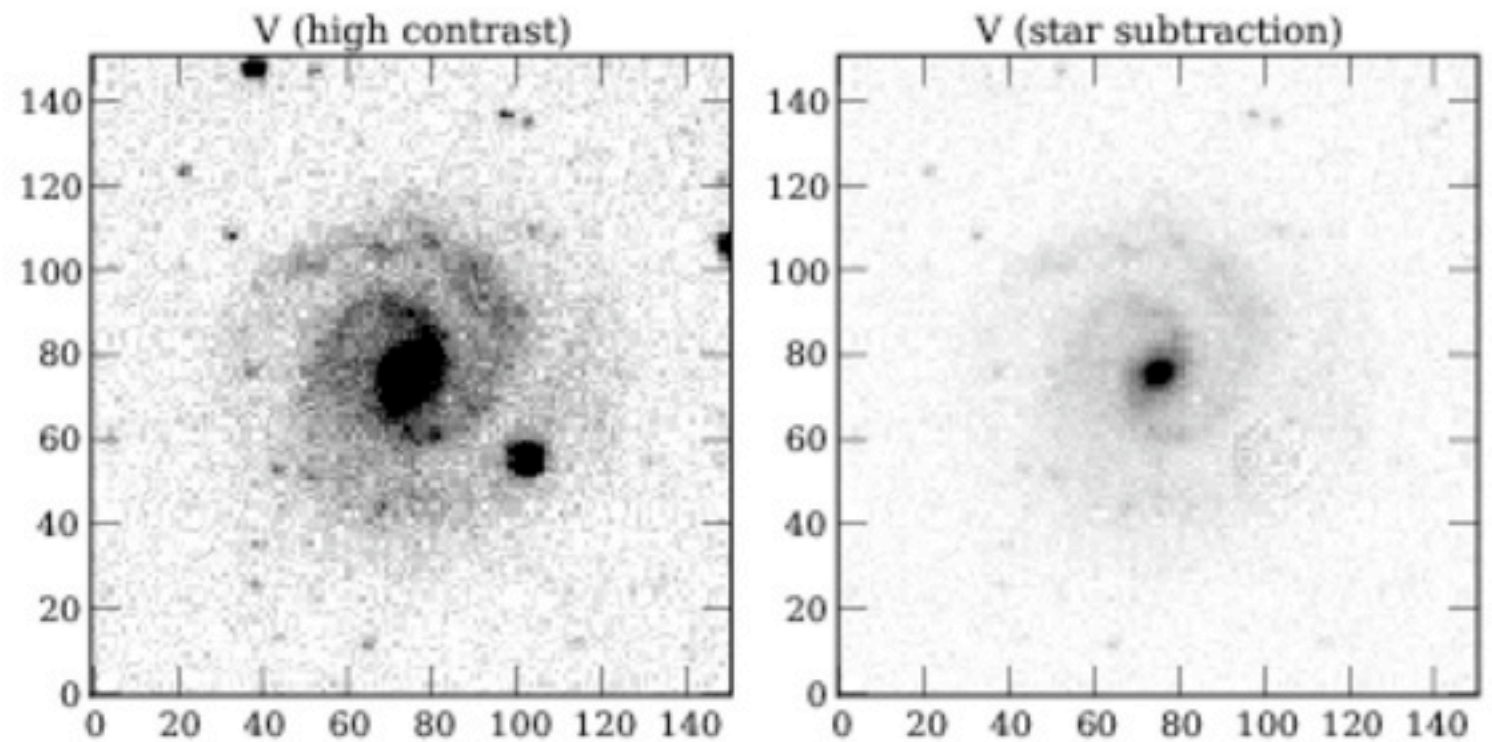


LSB galaxies
got spiral arms!

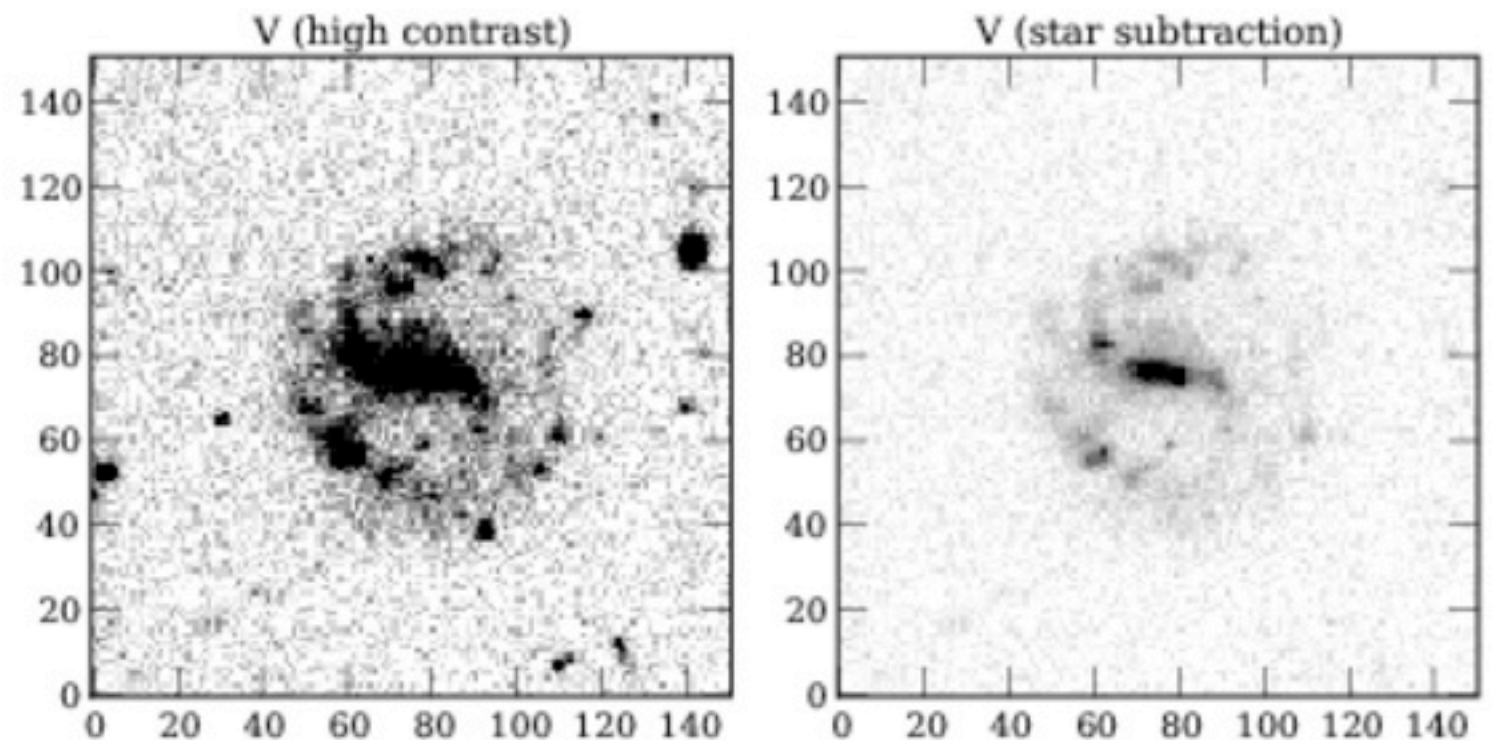
To explain this, we
anticipate the
need for very massive
disks to drive spiral
density waves in LSBs

McGaugh & de Blok
(1998), *ApJ*, 499, 66

F568-1



F577-V1



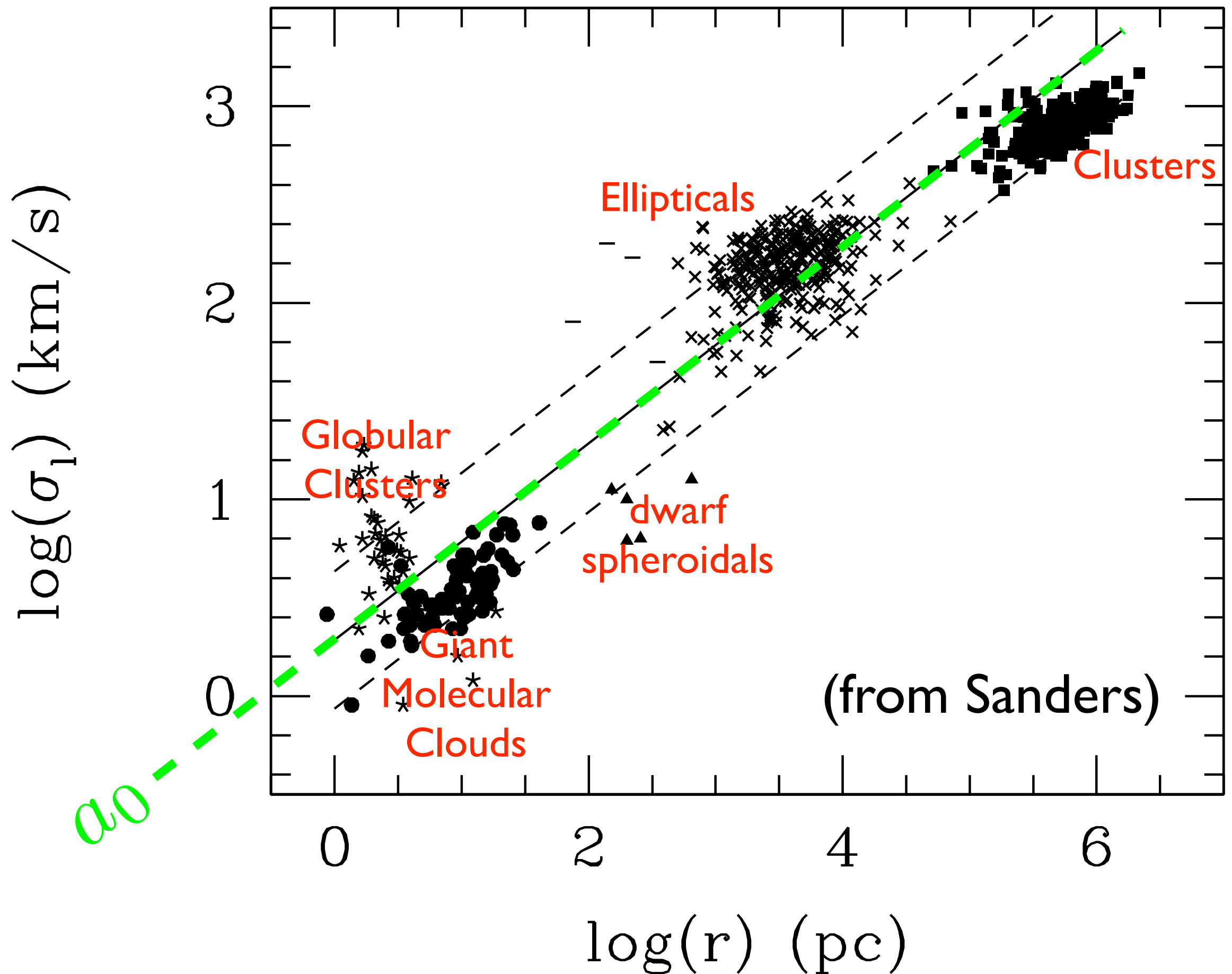
Disk Masses from Density Waves

Galaxy	$(M/L)_*$	AUTHOR
F568-1	14	FUCHS
F568-3	7	FUCHS
F568-6	11	FUCHS
F568-V1	16	FUCHS
UGC 128	4	FUCHS
UGC 1230	6	FUCHS
UGC 6614	8	FUCHS
ESO 14-40	4	FUCHS
ESO 206-140	4	FUCHS
ESO 302-120	1.7	FUCHS
ESO 425-180	2.4	FUCHS
ESO 186-550	7.5	SABUROVA
ESO 206-140	8.8	SABUROVA
ESO 234-130	5.7	SABUROVA
ESO 400-370	9	SABUROVA

Big $(M/L)_*$'s!

Conventional analysis overestimates M_*/L , as expected

Other tests - other systems

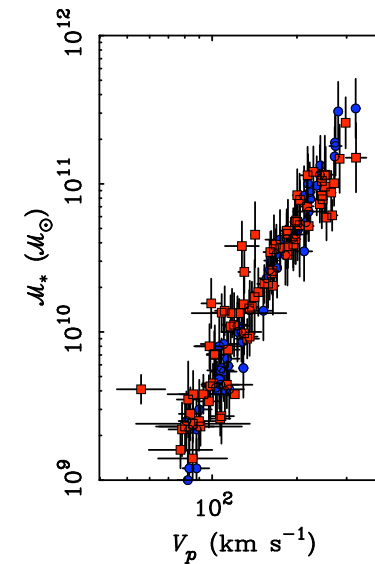
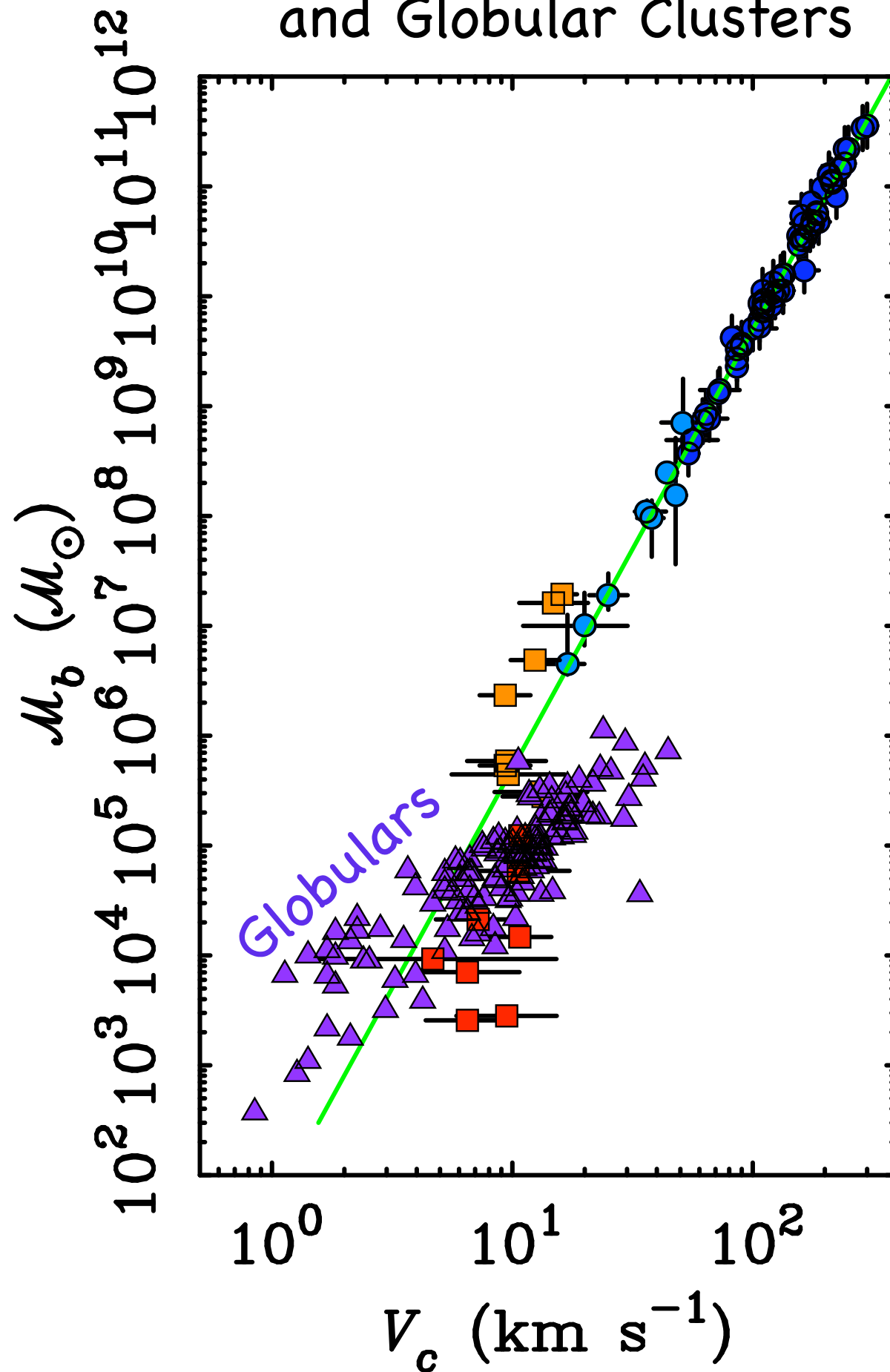


with LG dwarf Spheroidals and Globular Clusters

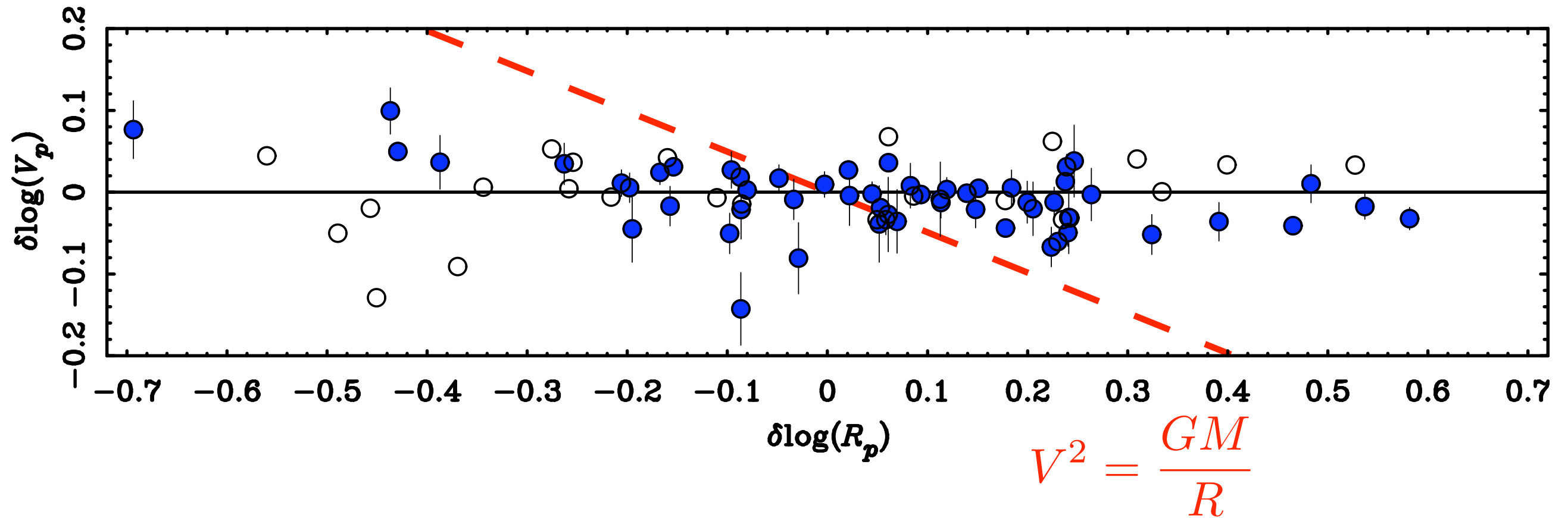
Assumes

$$\Upsilon_{\star} = 1$$

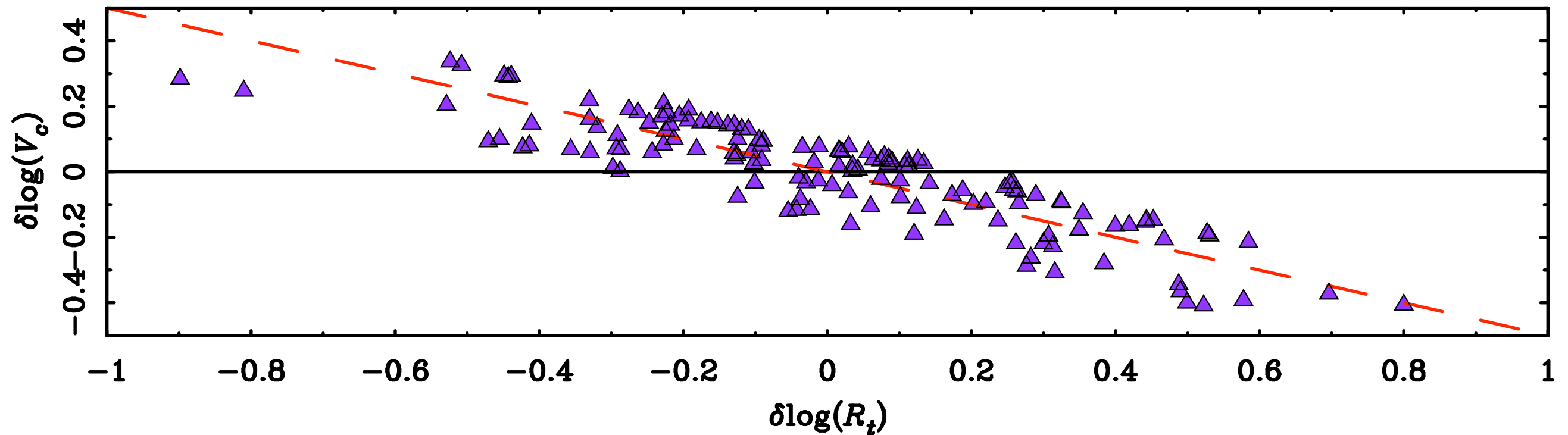
$$V_c = \sqrt{2}\sigma$$



No scale length residuals for Galaxies...



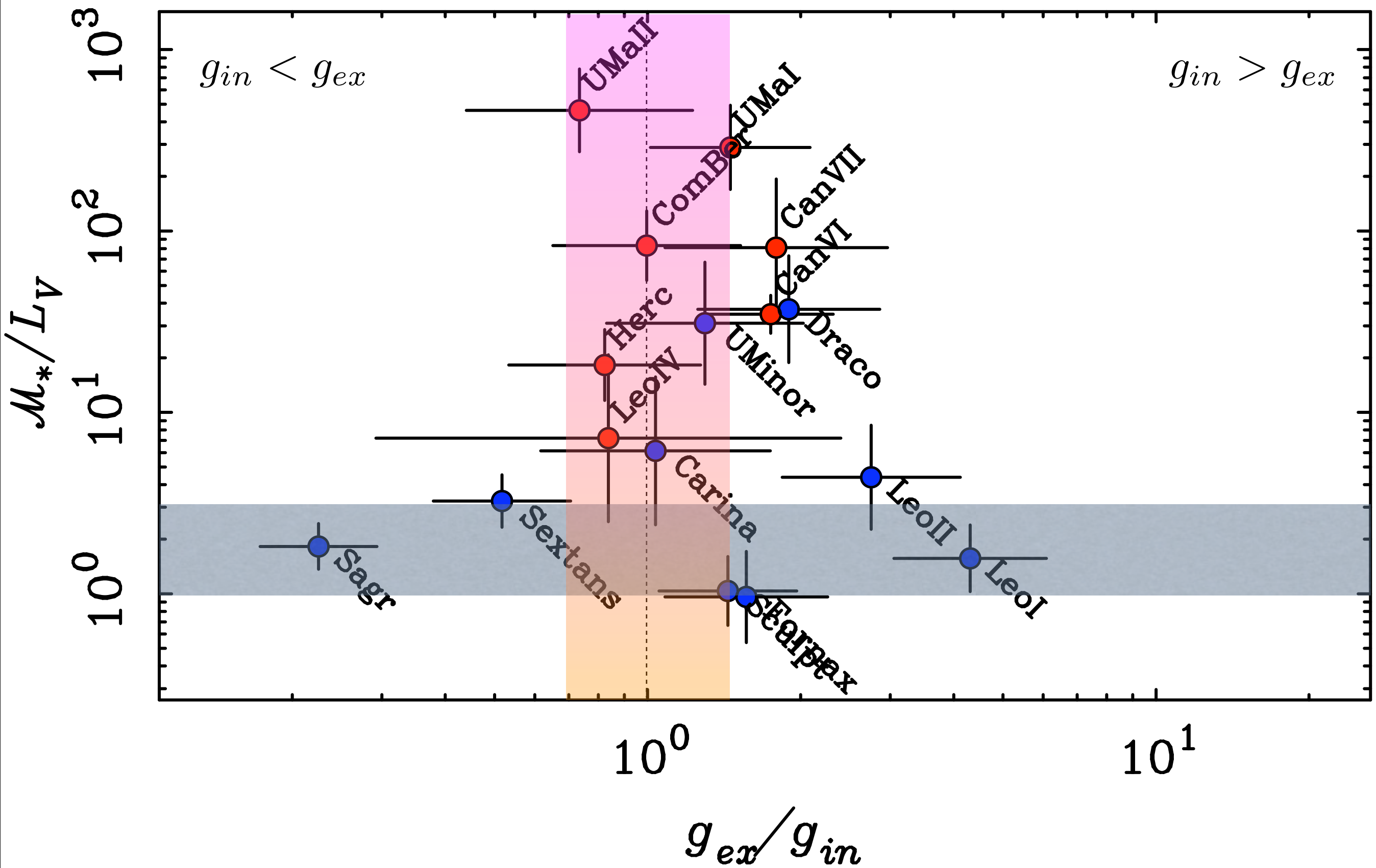
...but there are for Globular Clusters



Newton works (as he should in high density systems)

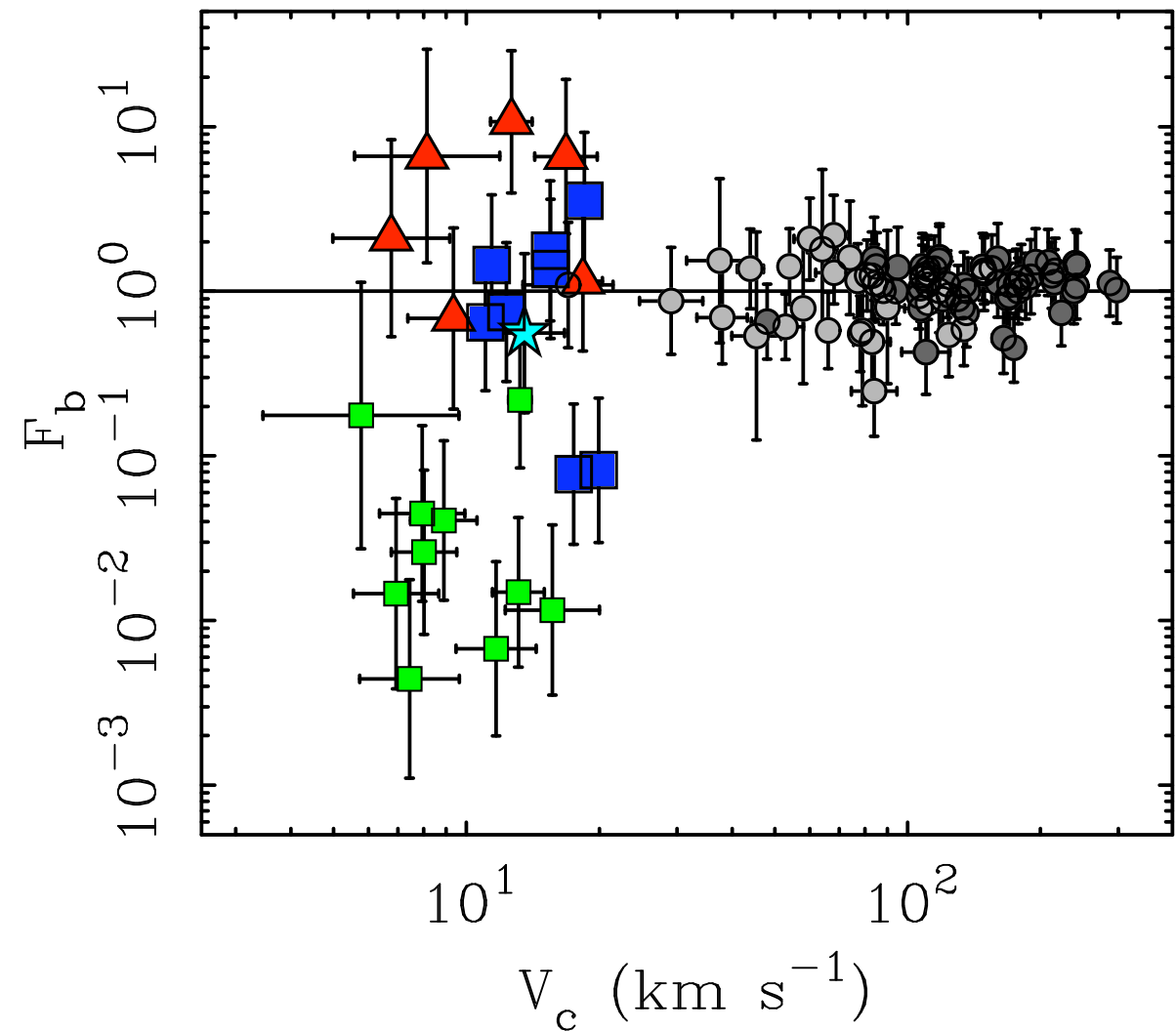
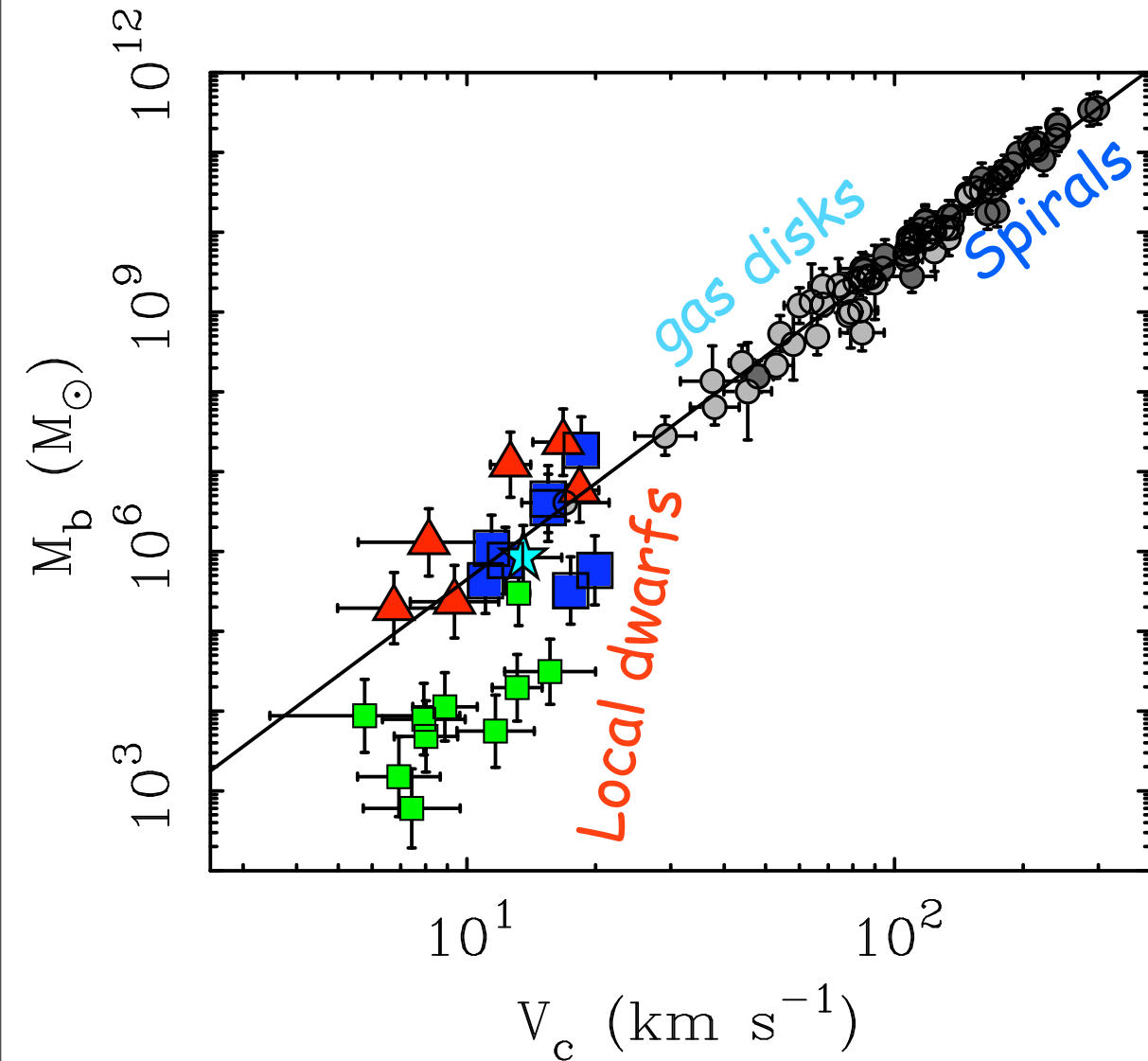
Dwarf Spheroidals

MOND M^/L OK for most classical dwarfs
but unacceptably high for ultrafaints*



Residuals of dwarf Spheroidals from Baryonic Tully-Fisher Relation

McGaugh & Wolf (2010)



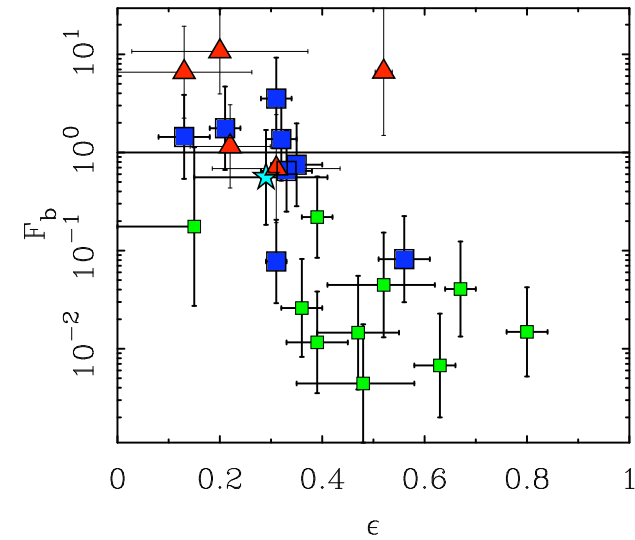
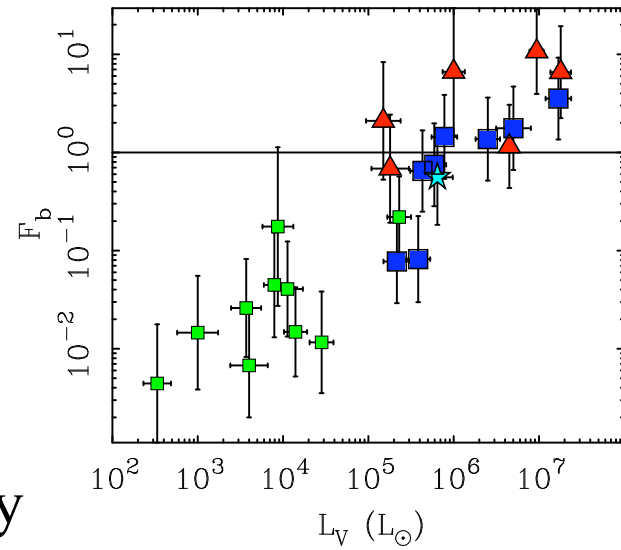
- Classical dwarfs
- Ultrafaint dwarfs
- ▲ M31 dwarfs
- ★ Leo T (contains gas)

Local dwarf data: Wolf et al. (2010)
 Kalirai et al. (2009; M31)
 M*/L as per Mateo et al. (1998)
 & Martin et al. (2008)

$$F_b = \frac{M_b}{AV_c^4}$$

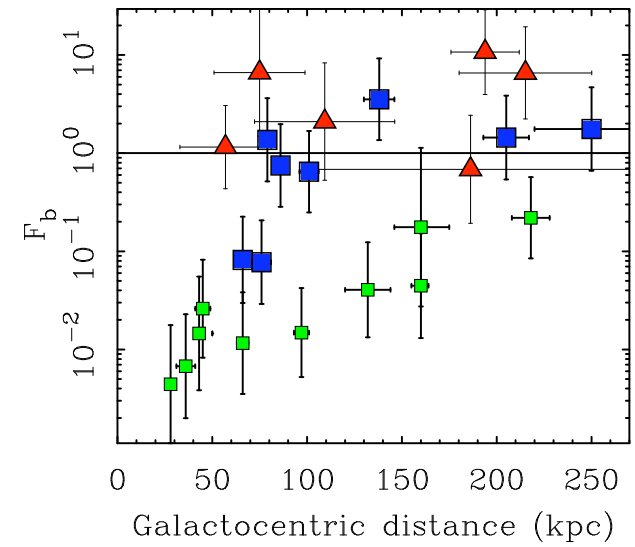
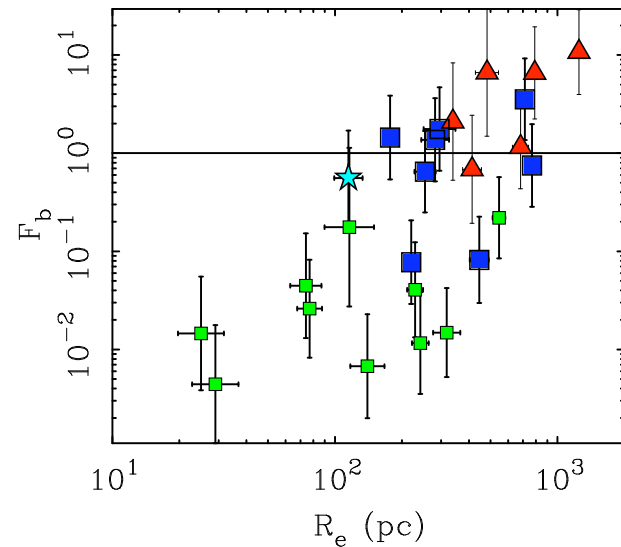
dSph BTFR residuals
correlate with

Luminosity



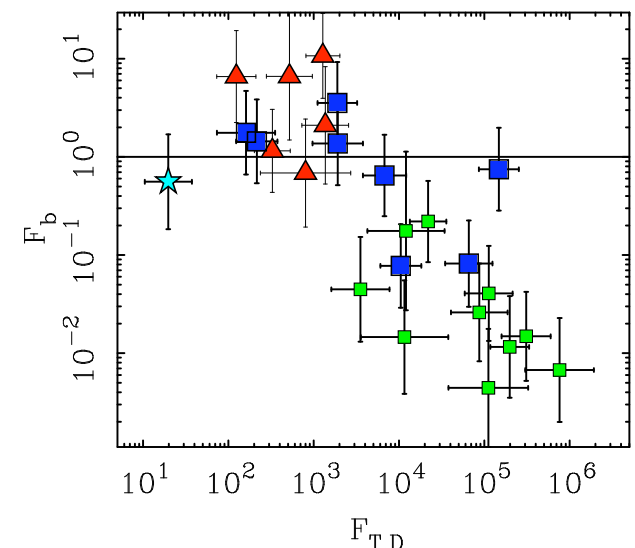
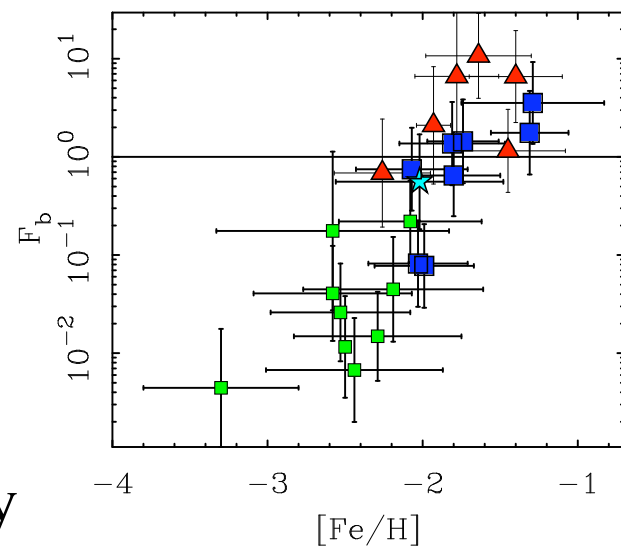
Shape

Size



Distance

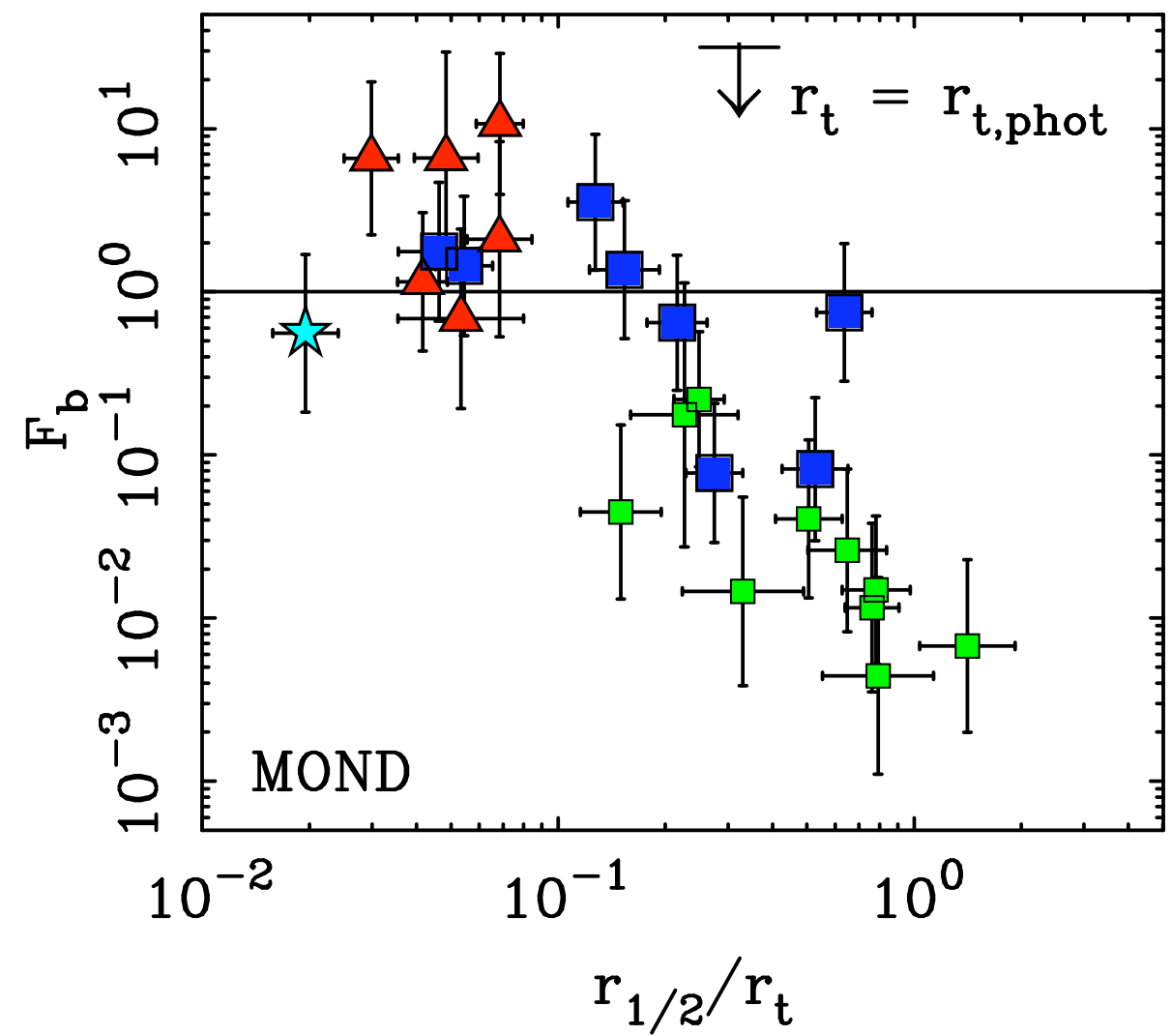
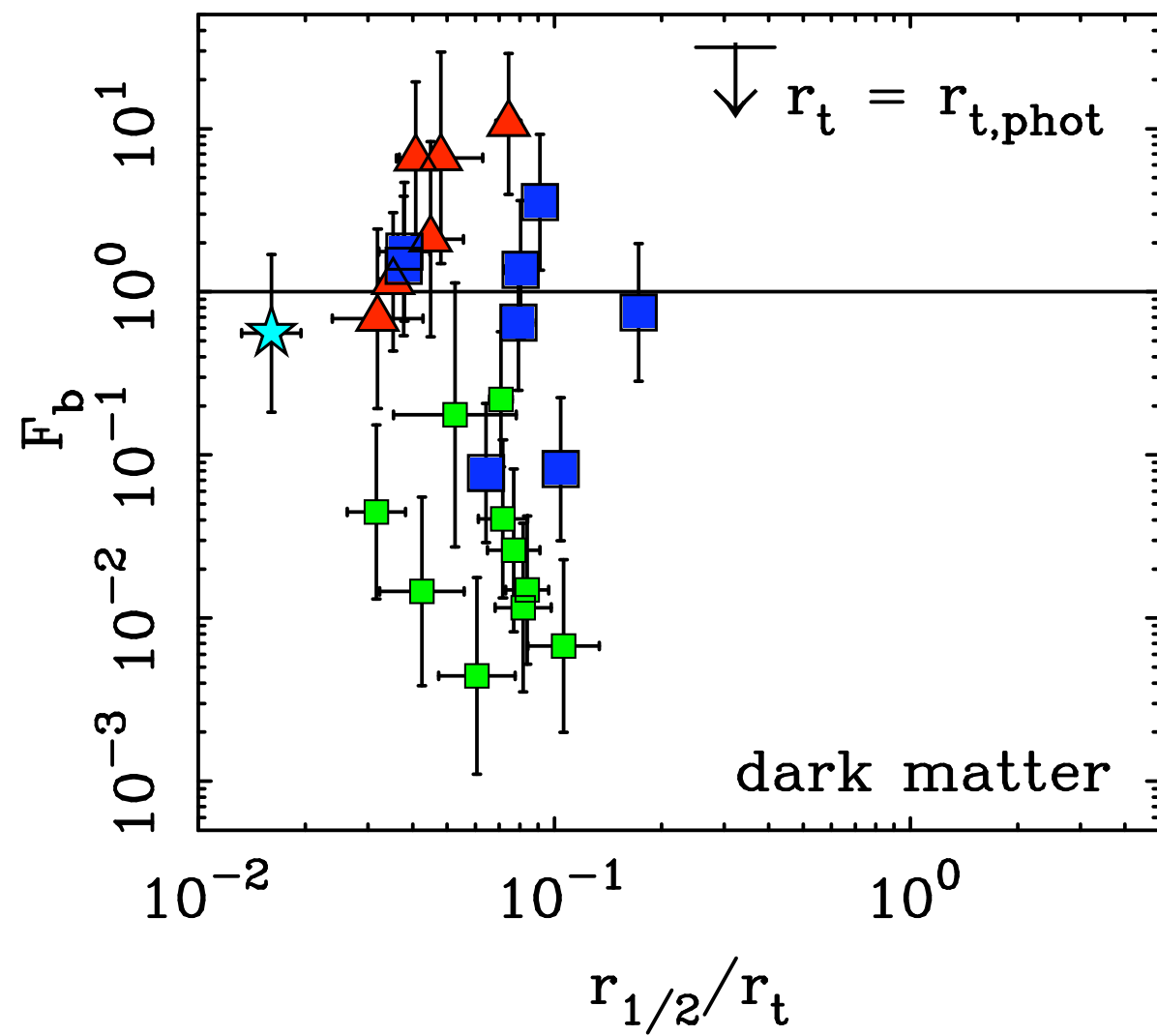
Metallicity

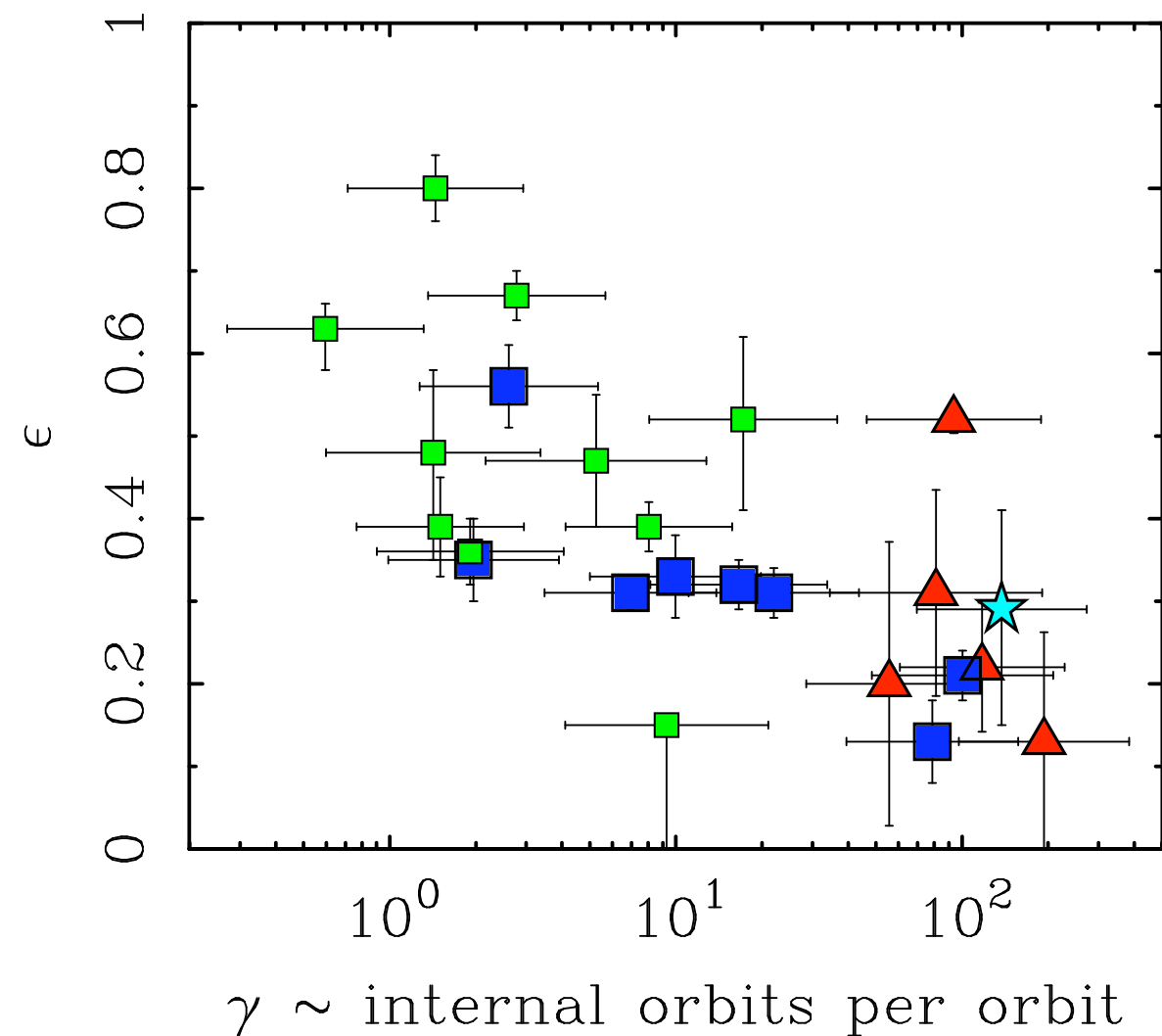
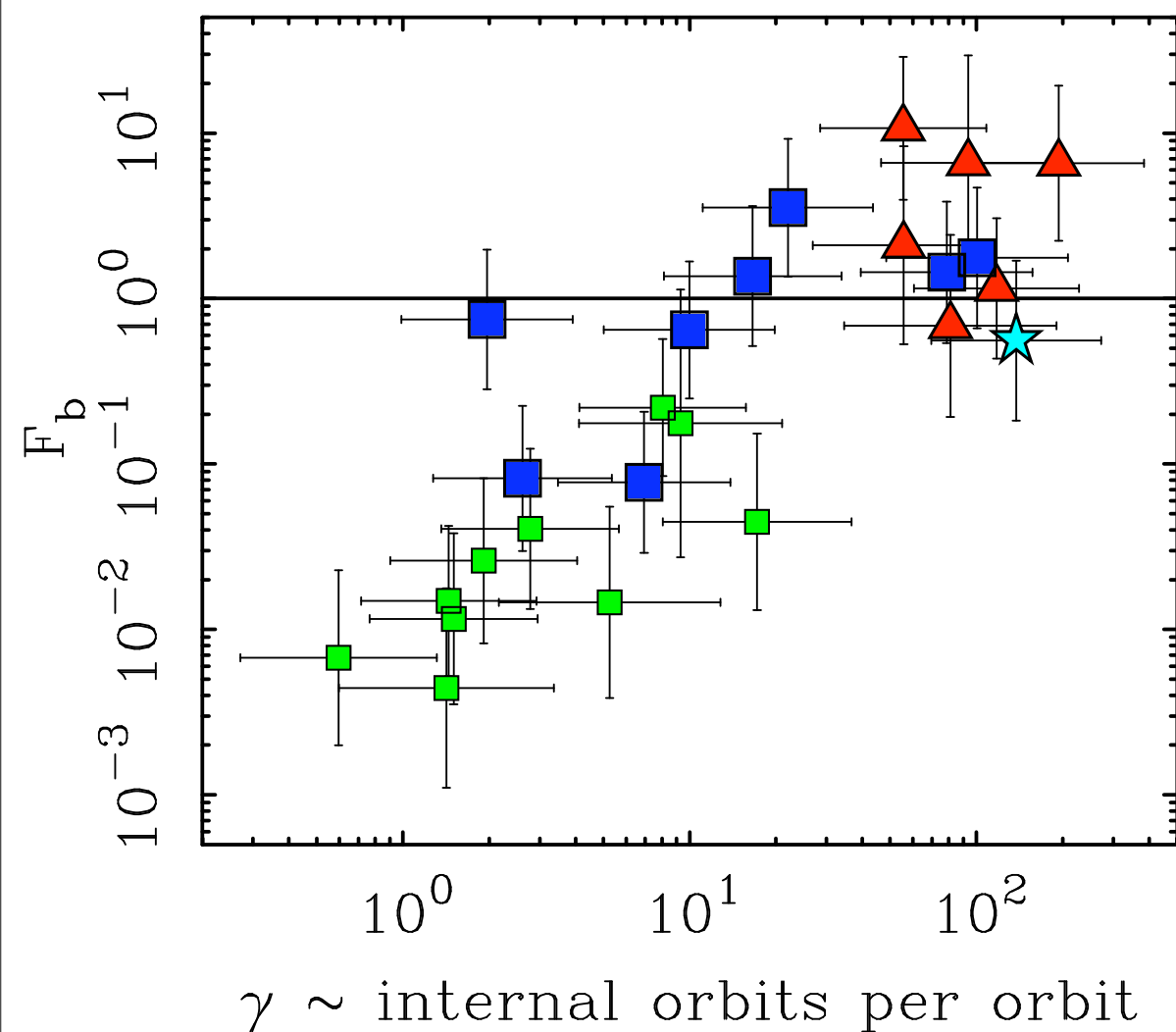


$$F_{T,D} = \frac{M}{m} \left(\frac{r}{D} \right)^3$$

Tidal Susceptibility

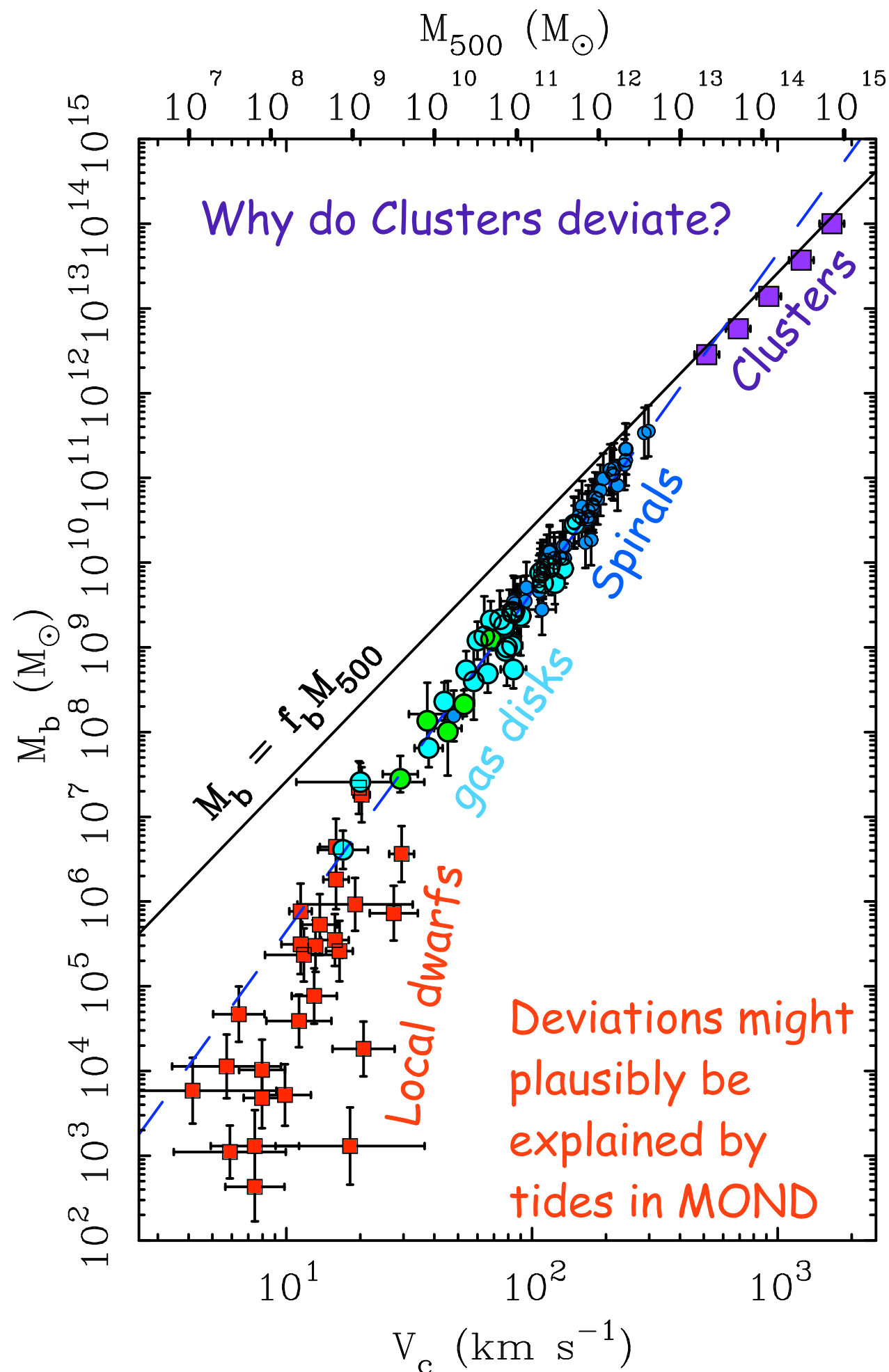
tidal radii in dark matter and MOND





Dwarfs whose stars have little time to adjust to changes in the potential suffer the largest deviations and have more elliptical shapes.

That the ultrafaints are tidally affected by the Milky Way in MOND.
Their M^*/L are overestimated by the usual equilibrium calculation.



Clusters have less baryonic mass than expected.

Cluster data: Giodini et al. (2009)

$$M_\Delta = B_\Delta V_\Delta^3$$

$$B_{500} = 1.5 \times 10^5 M_\odot \text{ km}^{-3} \text{ s}^3$$

Spiral data: McGaugh et al. (2005)

Gas dominated disks:

Stark et al. (2009)

Trachternach et al. (2009)

Local dwarf data: Walker et al. (2009)

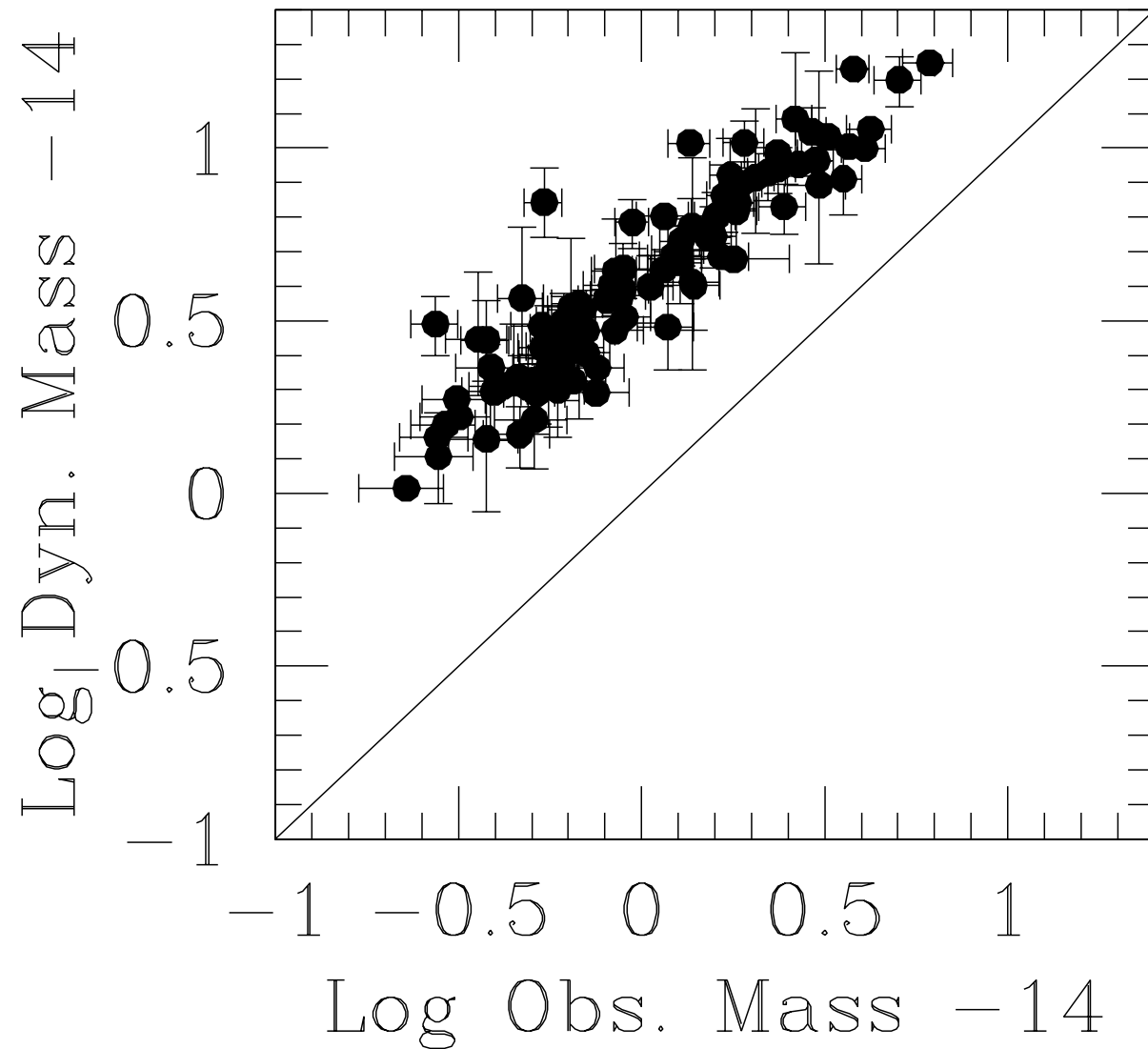
M^*/L as per Mateo et al. (1998)

$$V_c = \sqrt{3}\sigma$$

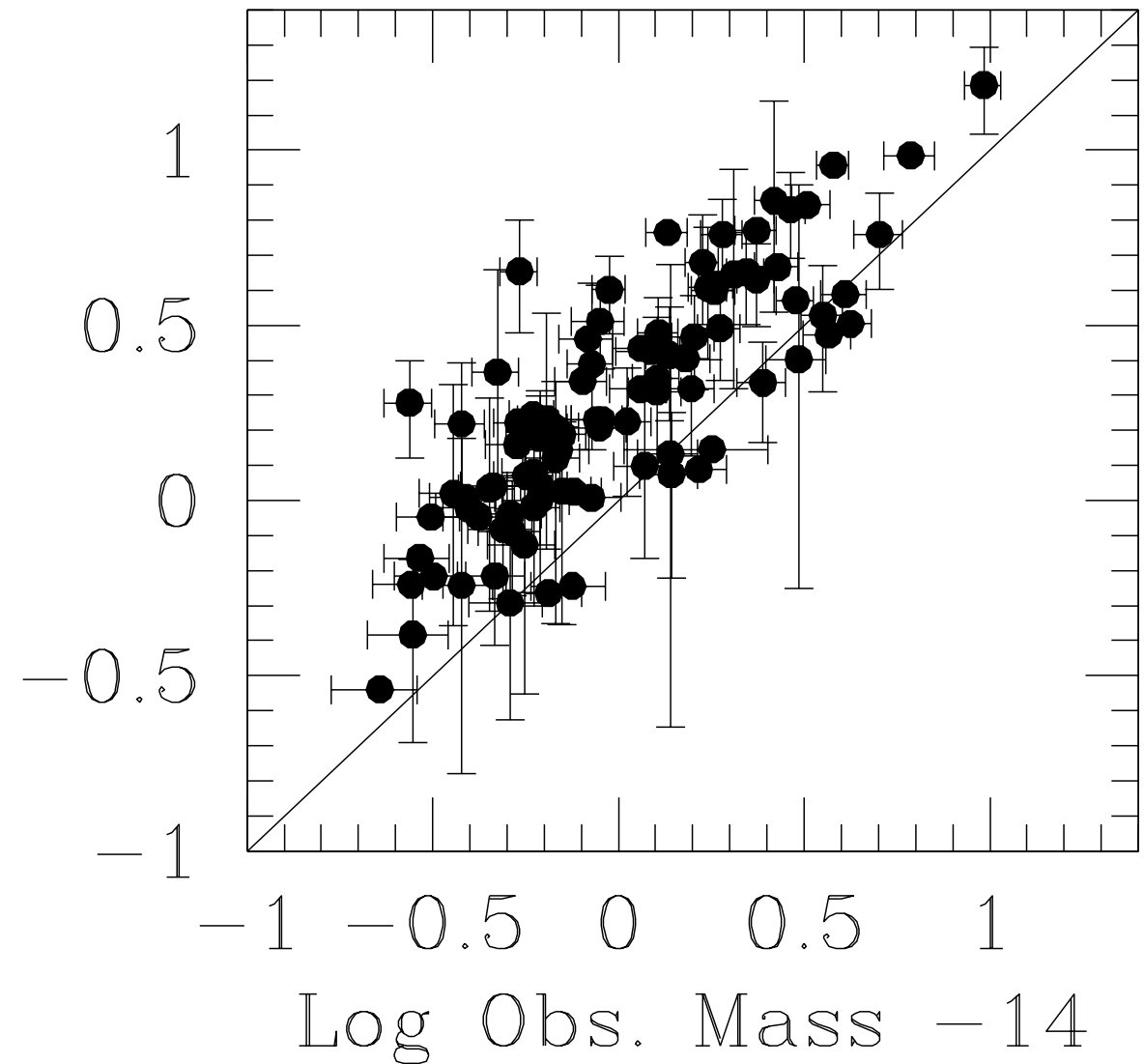
McGaugh et al. (2010)

Clusters of Galaxies

Newton

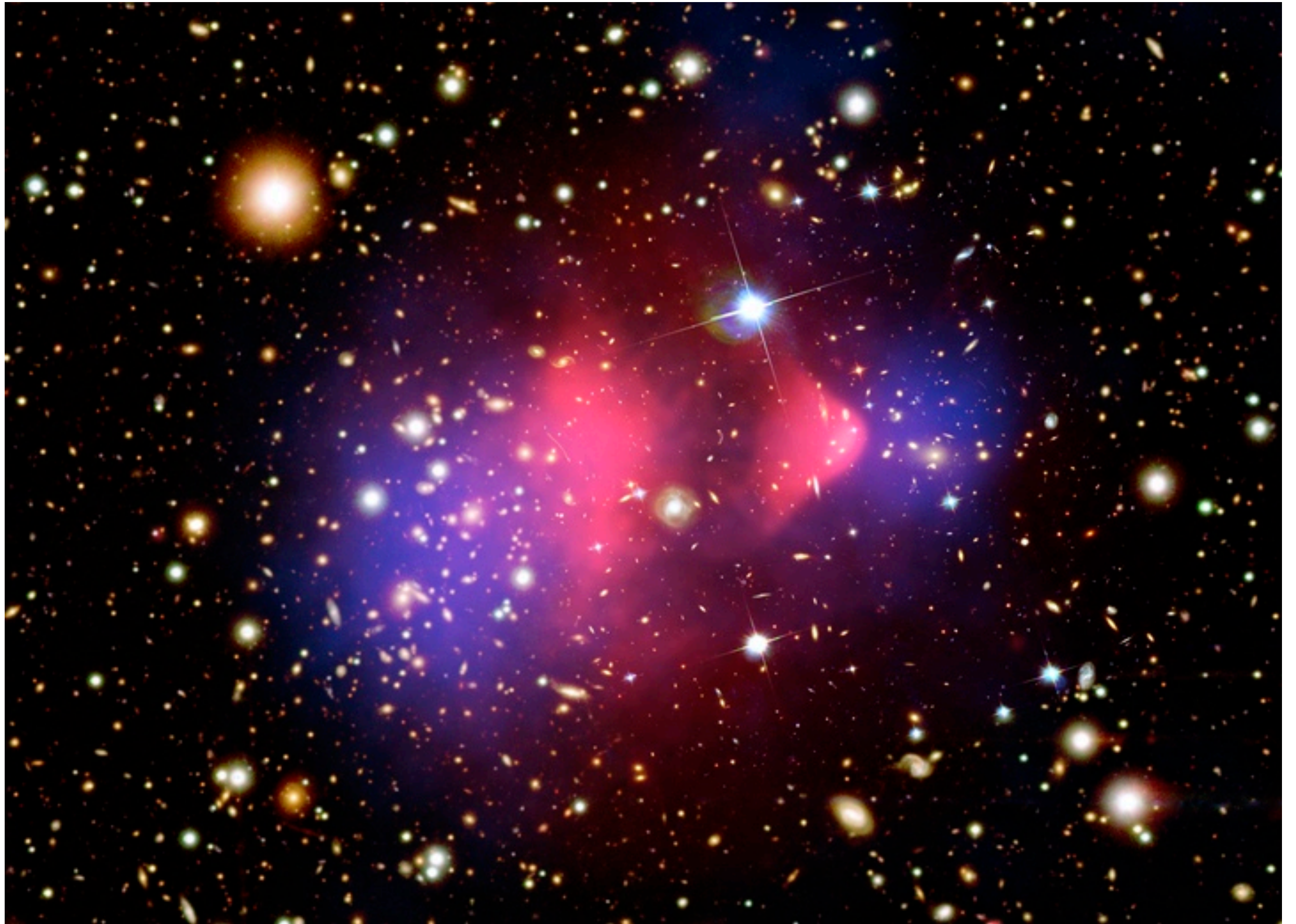


MOND

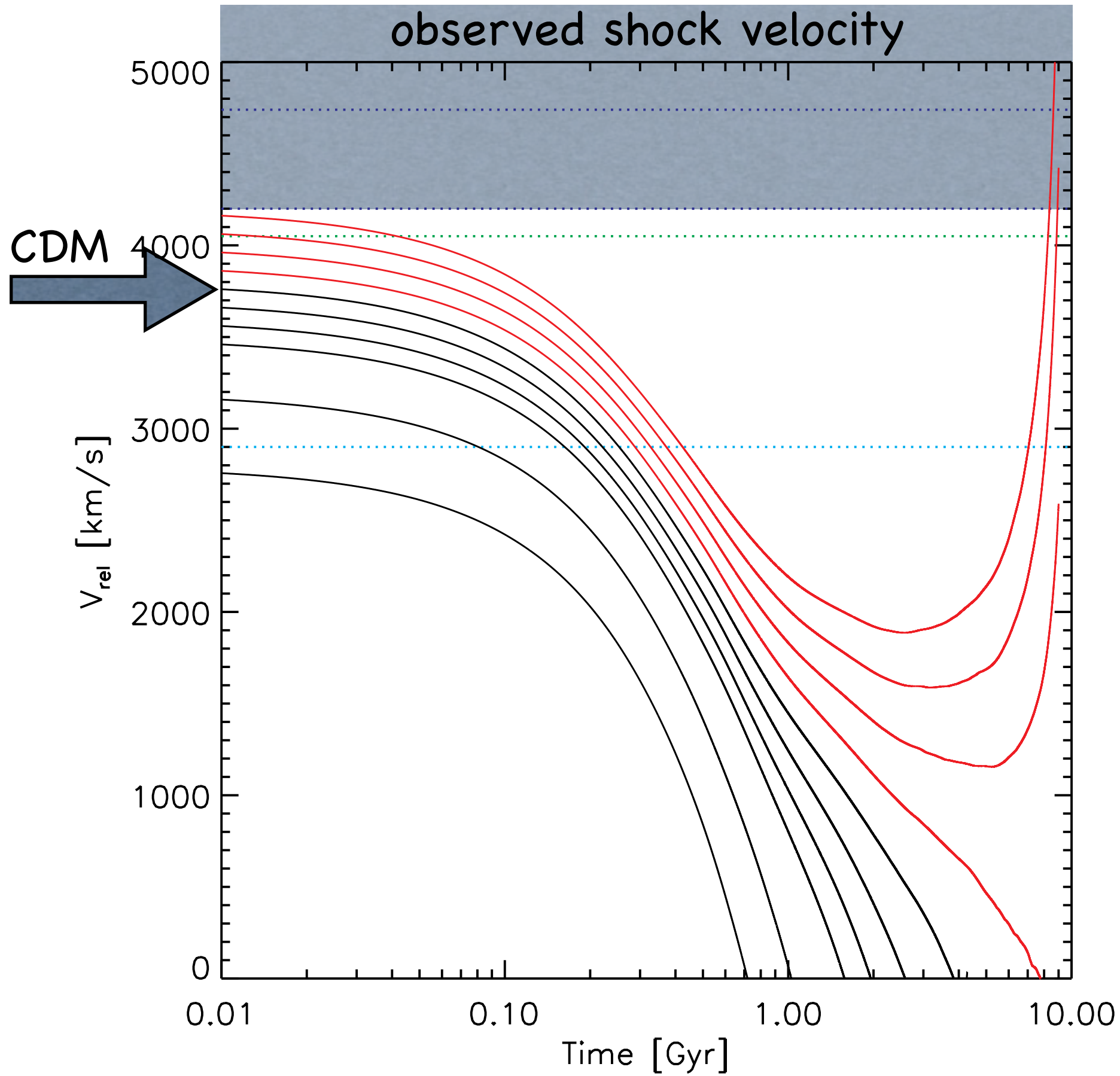


residual mass discrepancy in clusters is real...
the bullet cluster is a special case of a more general problem.

1E 0657-56 - “bullet” cluster (Clowe et al. 2006)

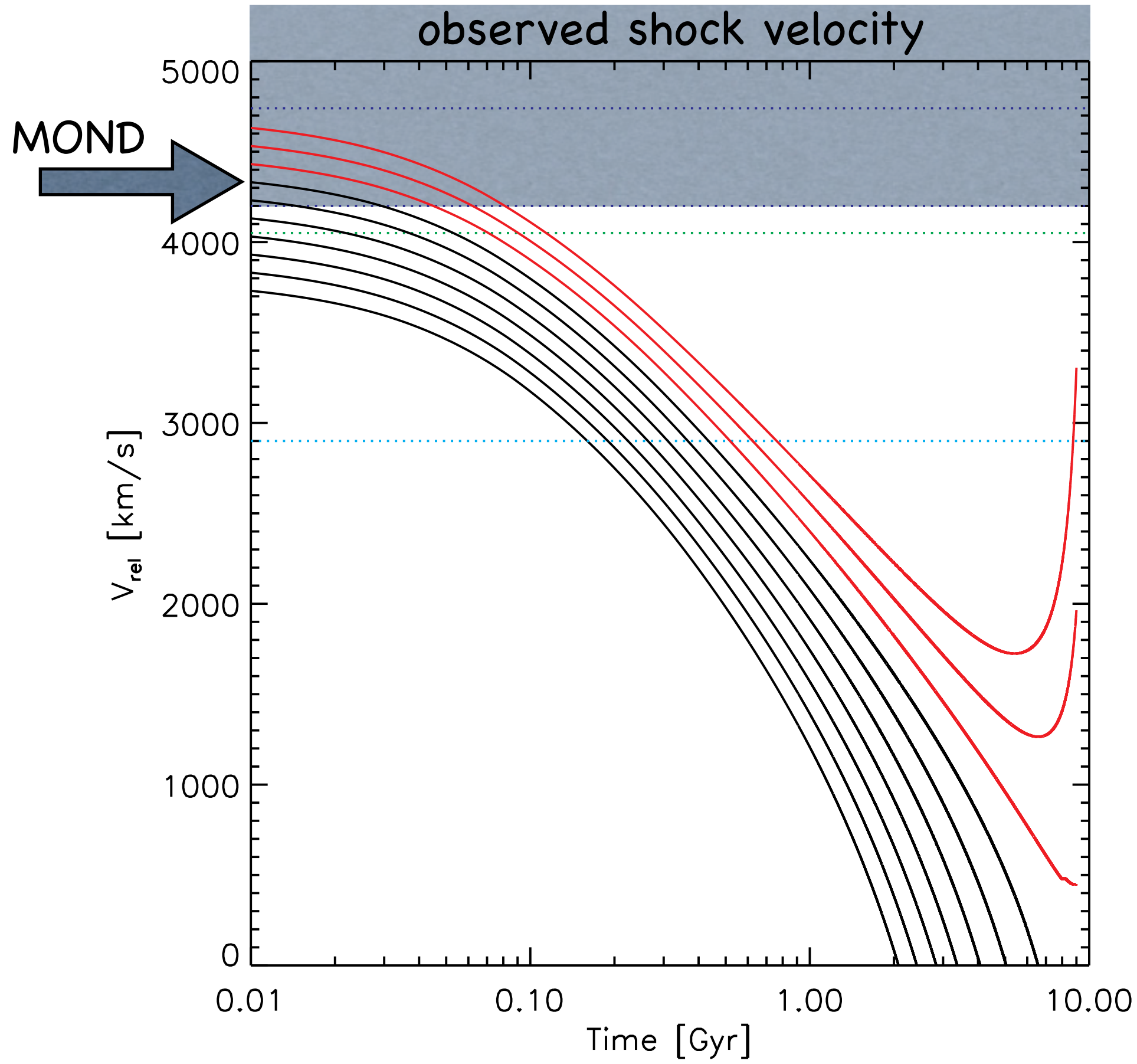


bullet cluster collision velocity



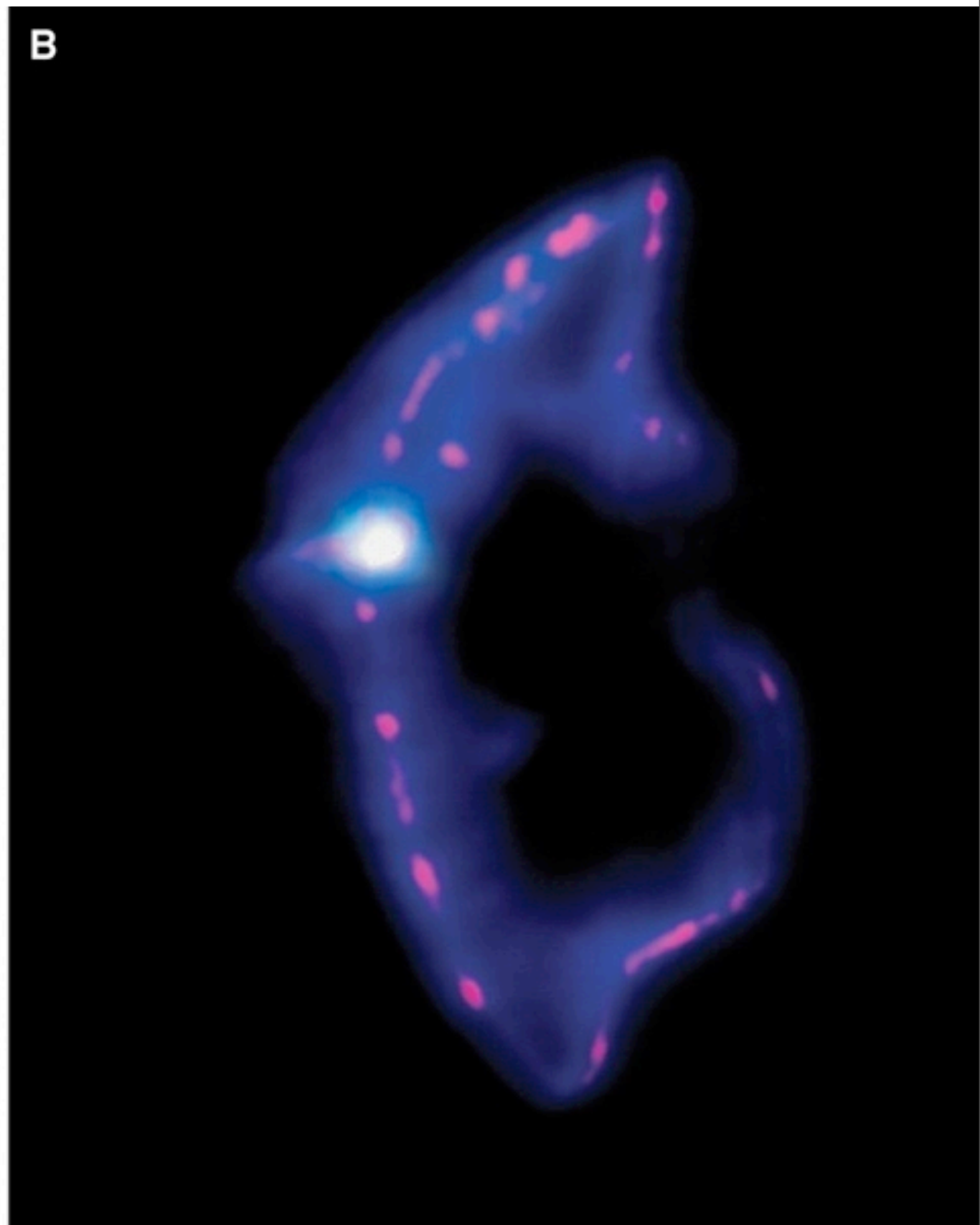
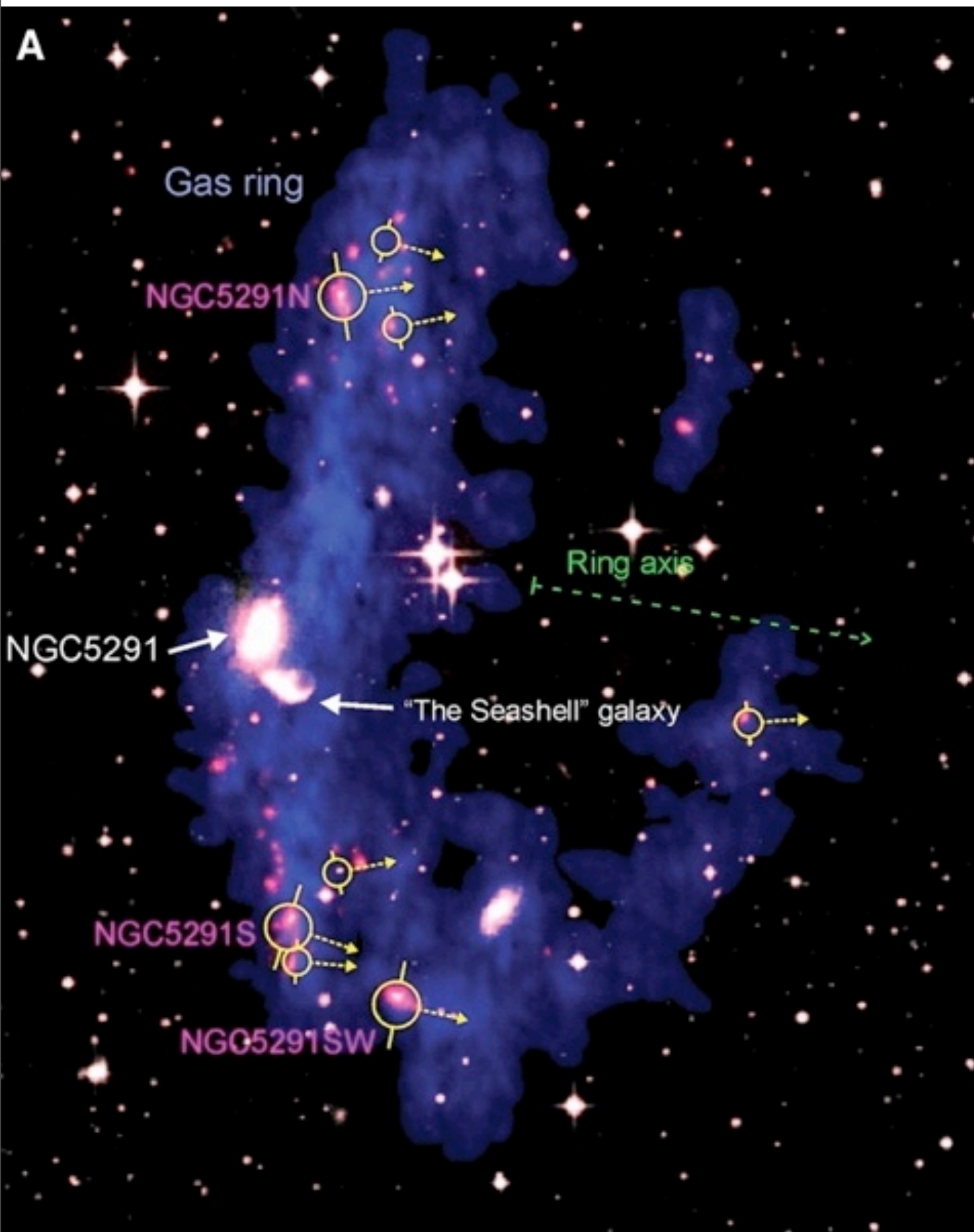
mass discrepancy more naturally explained by CDM

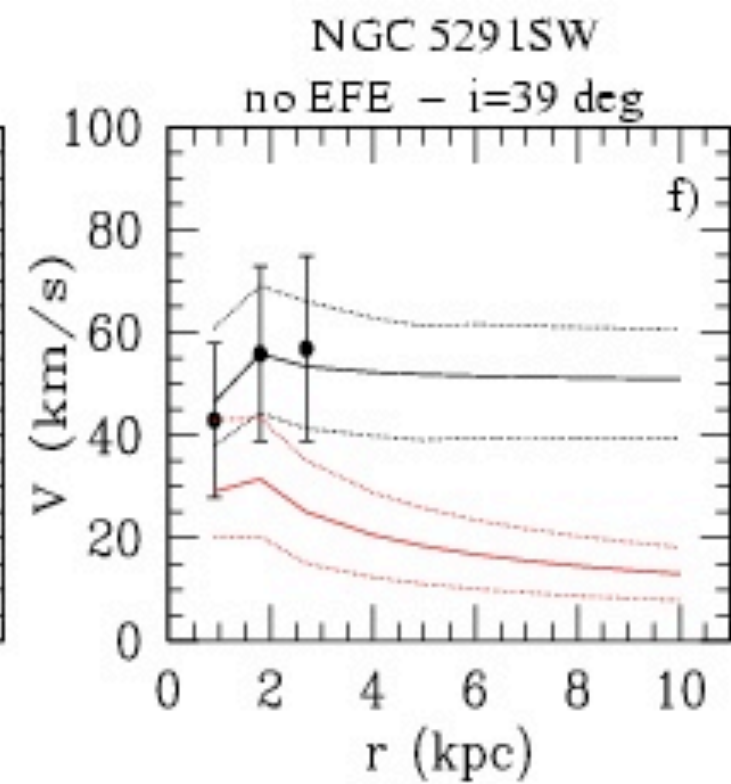
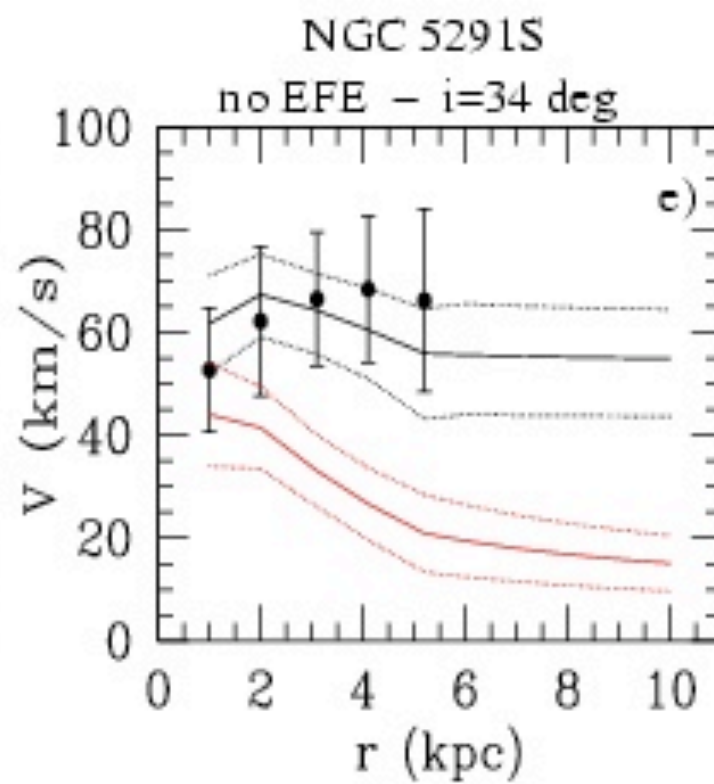
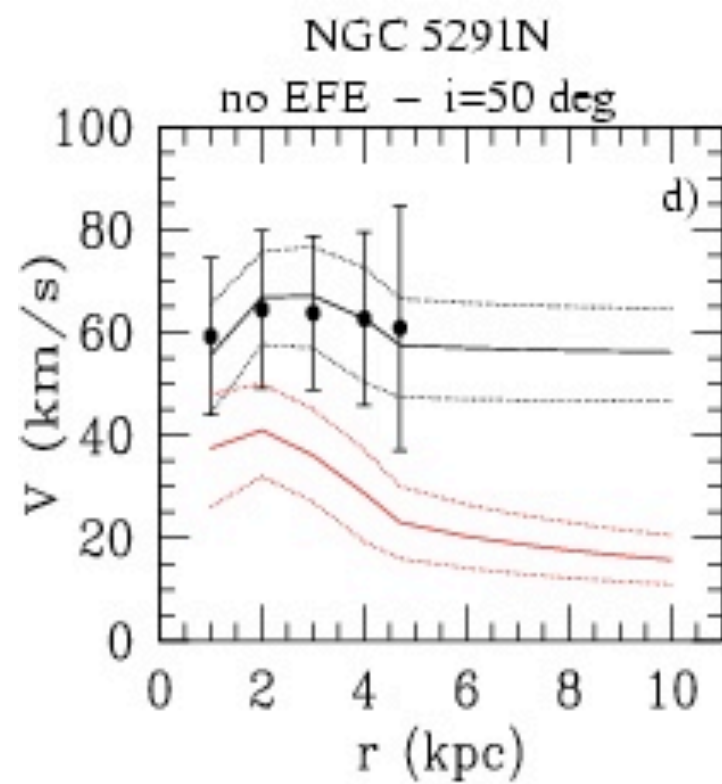
bullet cluster collision velocity



collision velocity more naturally explained by MOND

Tidal Debris Dwarfs - should be devoid of Dark Matter

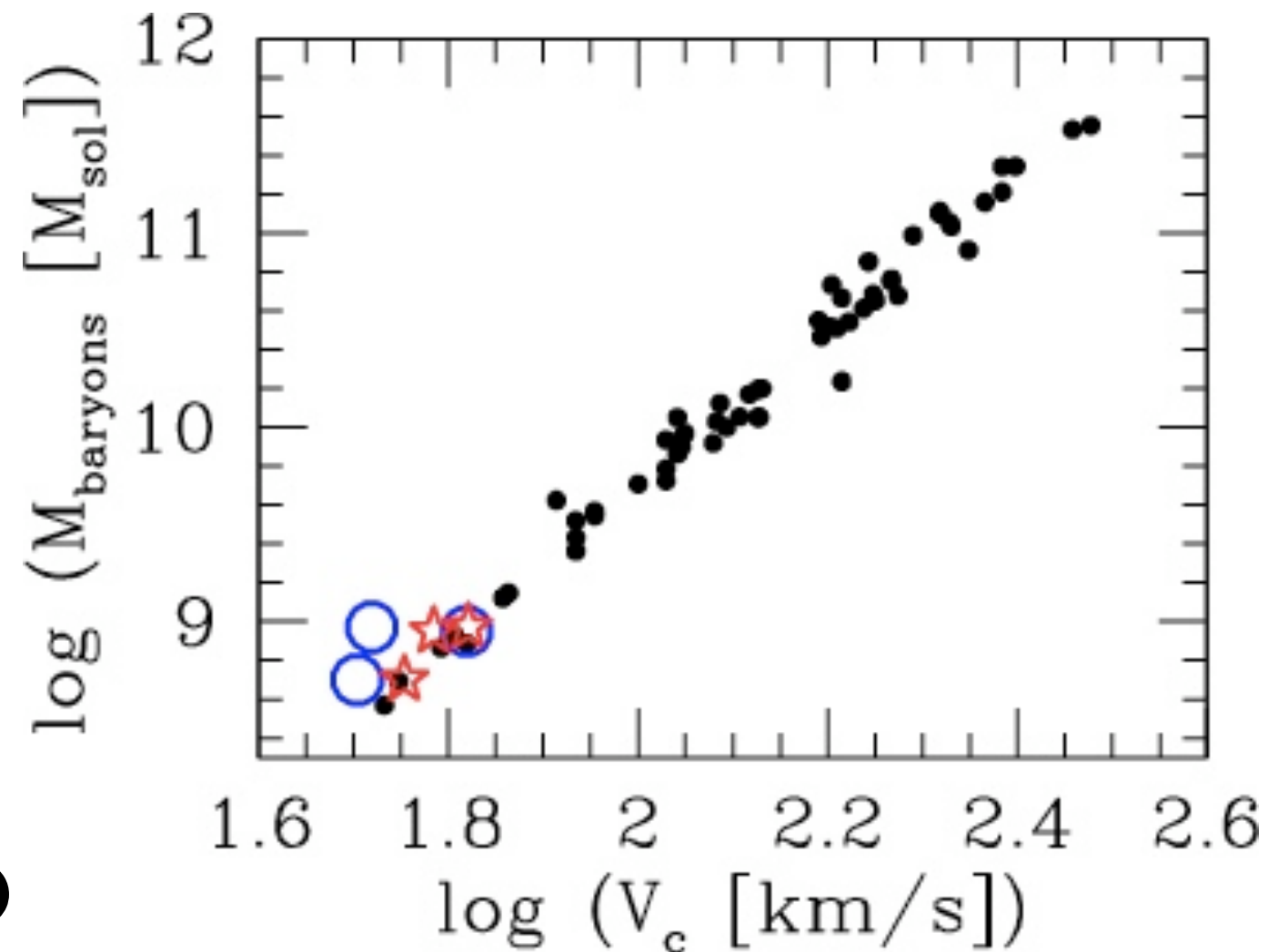




Gentile et al. (2007)

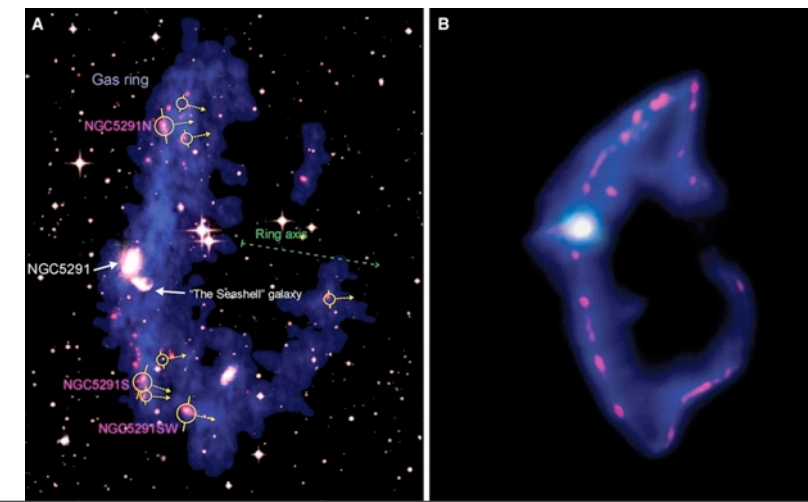
A&A, 472, L25

Tidal dwarfs
do show mass
discrepancies as
expected in MOND



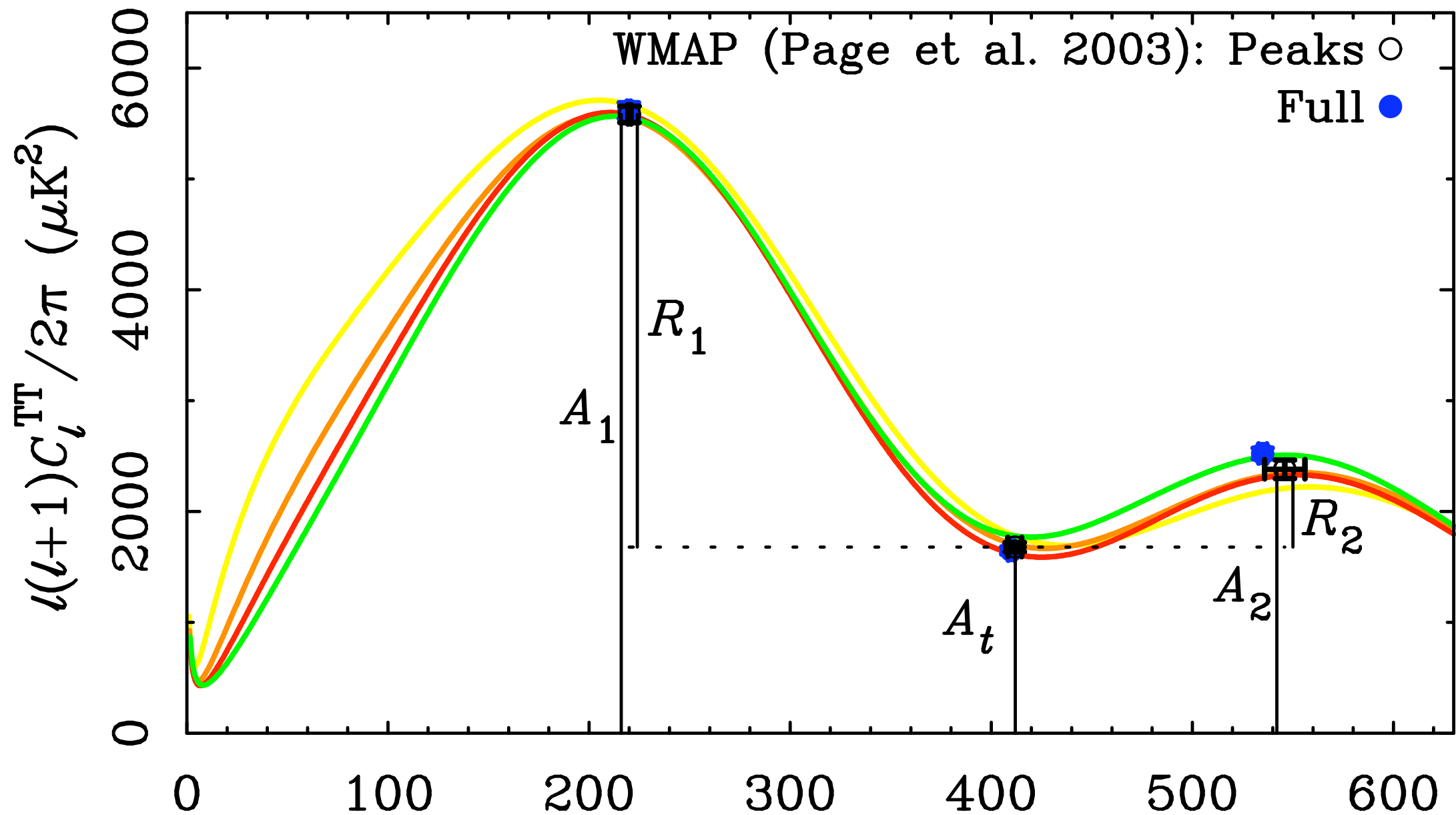
Conclusions

- MOND naturally explains a diverse array of phenomena
- Many *a priori* MOND predictions have been realized; some problems remain, especially in rich galaxy clusters
- The observed MONDian phenomenology is not naturally a part of the Λ CDM paradigm
- Can CDM be falsified???



No CDM prediction (McGaugh 1999): $A_{1:2} = 2.4$

Subsequent measurement: $A_{1:2} = 2.34 \pm 0.09$



Model	$L_{2:1}$	$A_{1:2}$	$R_{1:2}$
1999 LCDM	2.4	1.8	3.5 *
No CDM	2.6	2.4	5.4 *
WMAP Data	2.48	2.34	5.56